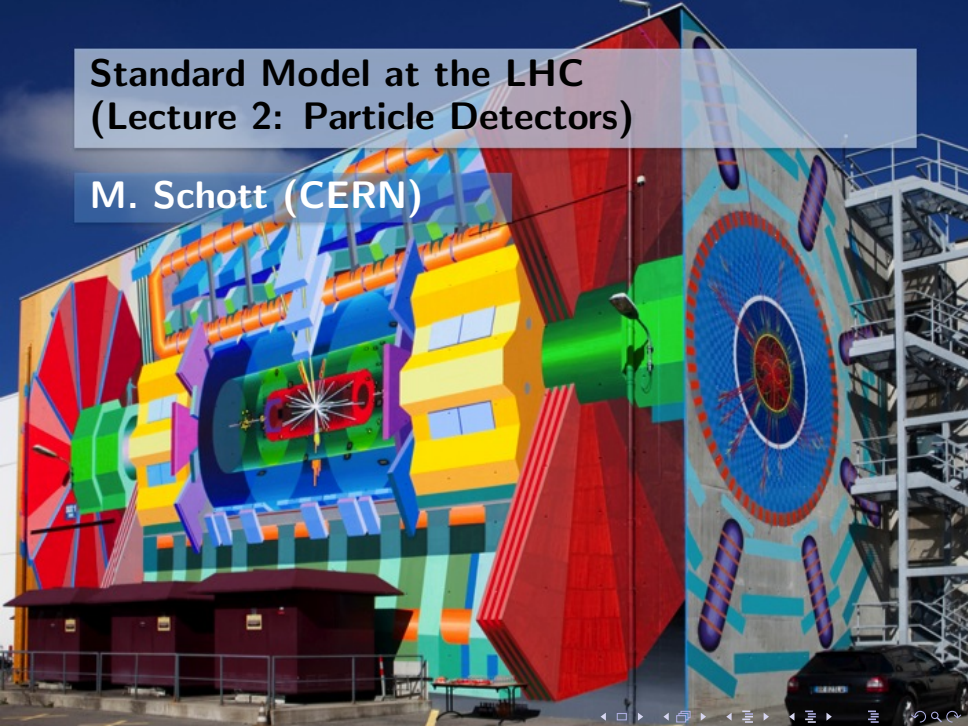


Standard Model at the LHC (Lecture 2: Particle Detectors)

M. Schott (CERN)



Content

- 1 Introduction
- 2 The LHC
- 3 Particle/Matter Interactions
- 4 Momentum Measurement
- 5 Energy Measurement
- 6 LHC Detectors

Why colliding protons

Electron-Positron Collider

- - interactions of point-like particles
- - exact energy known
- - QED processes 'easy' to calculate

Proton-Proton Collider

- - only a fraction of the beam energy 'used' in collisions
- - 'messy' QCD environment
- + negligible synchrotron radiation: $\Delta E_s \sim \frac{E^4}{m^4} \frac{1}{R} \rightarrow$ high energies
- + high luminosities
- \rightarrow good for discovering something new

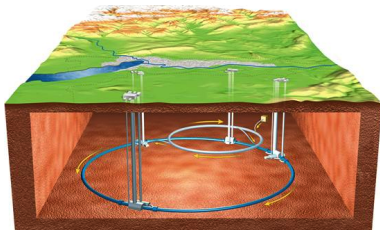
The Large Hadron Collider (1/2)

The Large Hadron Collider is

- 27km in circumference
- 100m underground (old LEP tunnel)
- a proton-proton collider with a center of mass energy of currently $\sqrt{s} = 8\text{TeV}$

Design Parameters

- 2835×2835 proton-bunches
- distance: 7.5m (25ms)
- per bunch 10^{11} proton
- per bunch-crossing: ~ 25 proton-proton collisions



The Large Hadron Collider (2/2)

The LHC consists of

- 1232 supraconducting dipole magnets
- field-strength 8.6 T with liquid helium

Design luminosity

- $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- high luminosity important as cross-sections for new physics are small
 - $\dot{N} = L \cdot \sigma$



Stable Particles

Relevant particles for the interaction with matter:

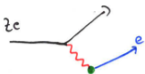
- stable particles: protons, anti-protons, electrons, anti-electrons, photons, atomic-cores, neutrinos
- long-lived particles (= do not decay within the detector): neutrons, myons, π^\pm , K^\pm

Hence we have here leptons and hadrons which interact via

- weak force (leptons): $\approx 10^{-5}$
- electromagnetic force (leptons, hadrons): $\approx 1/137$
- strong force (hadrons): ≈ 1

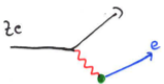
Electromagnetic Interaction of heavy particles (1/3)

The electromagnetic interaction is described by the exchange of virtual photons

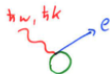


- ionization
- excitation of atoms
- polarization, cerenkov-effect, transition radiation

Simplify the interaction of particles with matter to interactions of photons with matter



\Rightarrow



Electromagnetic Interaction of heavy particles (2/3)

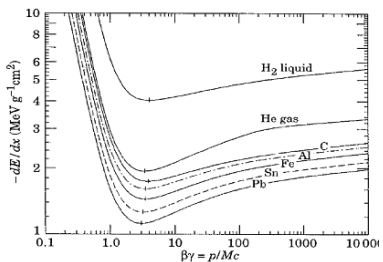
Within this model we can calculate the differential energy loss dE/dx via the famous Bethe-Bloch formulae

$$- \left\langle \frac{dE}{dx} \right\rangle = 2\pi r_e^2 m_e c^2 \rho \frac{N_0 Z z^2}{A \beta^2} \left[\ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{max}}{I} - \beta^2 \right]$$

with

- r_e : classical electron radius
- N_0 : Avogadro-number
- ρ : density of matter
- I : Ionisation potential of the atom
- T_{max} : max. kin. energy which an electron at rest can get by a central impact

Electromagnetic Interaction of heavy particles (3/3)

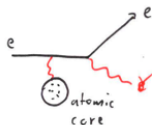


- Typical order of magnitude keV/cm
- Plateau : Saturation through energy loss via ionisation

Electrons and Positrons (1/2)

So far: heavy particles which ionize matter (atoms). We need some modifications for electrons and positrons

- Mass of impact-partners is the same
- Impact-partners cannot be distinguished
- Annihilation process ($e^+ e^- \rightarrow 2\gamma$)
- Bremsstrahlung
 - Dominating process when energies are $E \gg 10\text{MeV}$, i.e. at the LHC



Electrons and Positrons (2/2)

Cross-Section for this process is

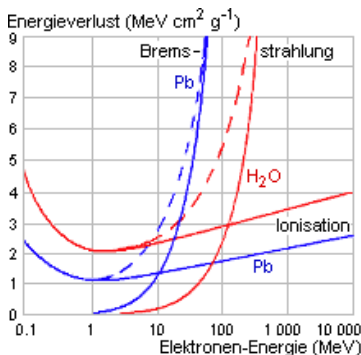
$$\sigma \sim \frac{1}{m^2},$$

- important for e^\pm but not for μ^\pm .

Note that

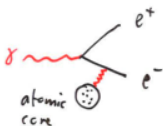
$$\langle \frac{dE}{dx} \rangle_{rad} \sim E \rightarrow -\frac{dE}{E} \sim dx$$

- Hence the energy of a particle that is transversing a detector is $E(x) = E_0 \cdot \exp(-x/X_0)$
- After a distance of X_0 (= radiation length), the initial energy is reduced to 37%.
- Important for detector dimensions ($X_0^{H_2O} = 36.1cm$, $X_0^{Pb} = 0.56cm$)



Photons

Pair-production is the dominant process for photons with an energy above $E_\gamma > 2MeV$



This is very similar to Bremsstrahlung.

As a rule of thumb: A high energetic photon converts after X_0 with a probability of $\sim 54\%$.

Summary of interactions

Summary

- electrons@LHC: Bremsstrahlung
- photons@LHC: Pair-production
- muons@LHC: Ionisation
- hadron@LHC: Ionisation, Nuclear-Interactions
- neutrinos@LHC: No signature in the detector

What do we want to measure?

- Ultimate goal is particle identity and 4-Momentum (Lorentz-Momentum Vector)
- particle identification via detector design
- 4-momentum measurement: Measure $|\vec{p}| \approx E$ and the flight direction w.r.t to collision point

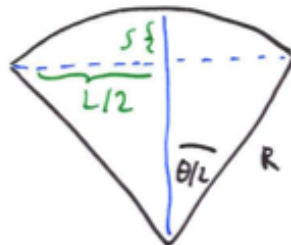
Momentum Measurement: Basic Concept (1/2)

- Idea: Use magnetic field to bend trajectory with charged particle. Momentum is given by

$$p = e \cdot E \cdot R$$

- where R is the bending radius and B is the magnetic field.
- Measurement via Sagitta

$$S = R - \sqrt{R^2 - (L/2)^2} \approx R \frac{\theta^2}{8}$$



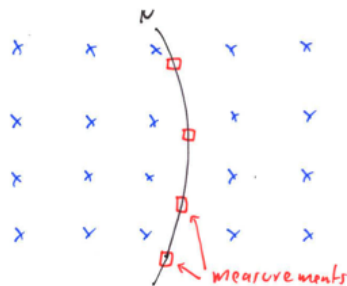
$$p = e \cdot B \cdot \frac{L}{2 \sin \frac{\theta}{2}} \rightarrow S \approx \frac{e \cdot B \cdot L^2}{8p}$$

Momentum Measurement: Basic Concept (2/2)

- Trajectory of particle in magnetic field is measured at N (equidistant) points
- Momentum resolution (for solenoid magnetic field) is given by

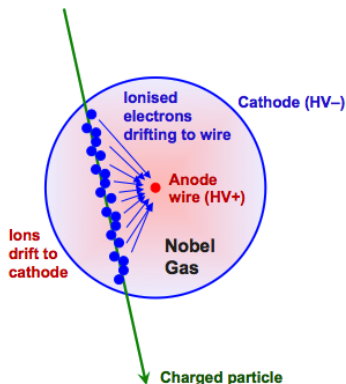
$$\frac{\sigma(p_T)}{p_T} = \frac{\sigma_{pos}}{0.3 \cdot B \cdot L^2} \sqrt{\frac{720}{N+4}} \cdot p_T$$

- Consequences
 - Larger Magnetic field \rightarrow better resolution
 - Even better to have longer measurement distance!
 - Rel. resolution worsens with increasing p_T



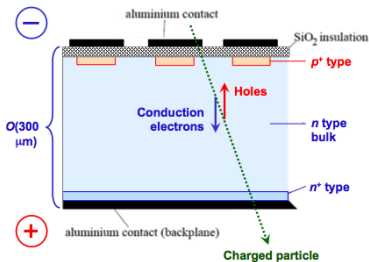
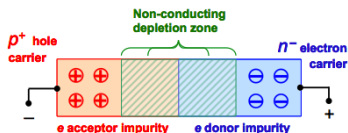
Gas Based Detectors

- Idea: Use ionisation of gas-atoms by incident particles
 - Primary Ionisation:
$$x + A \rightarrow x + A^* + \delta_e$$
 - Secondary Ionisation:
$$\delta_e + A \rightarrow \delta_e + A^* + e^-$$
 - Typical value for producing one ion-pair: $\approx 30\text{eV}$.
- Measurement: Voltage in gas
 - electrons/ions drift to cathode/anode
 - When the E-field gets strong (close to wire): Ionisation electrons can cause new ionisations
 - \rightarrow electron-avalanche \rightarrow measurable signal



Semi Conductor Detectors

- Using the p-n junction as a tracking detector
- enhance depletion zone with external voltage
- High Si density low electron-hole creation potential (3.6 eV compared to 36 eV for gaseous ionisation) allows use of very thin detectors with reasonable signal
- Very high spatial resolution ($\approx 10\mu m$)



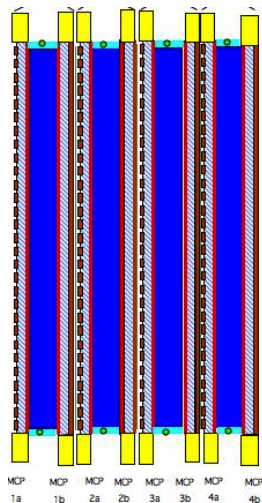
Basic Concept of Energy Measurements (1/2)

Energy-measurements are destructive, i.e. the particles deposit their full energy in the detector

- cannot be used for further measurements

Detectors for the energy measurement are called calorimeters

- homogenous calorimeters: Absorber material is also detection material
- sampling calorimeter: Absorber material is passive and placed in layers. In between active detection layers for ionisation measurements



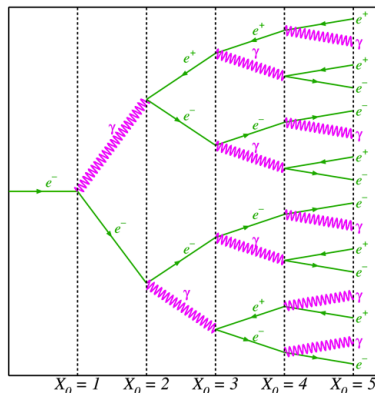
Basic Concept of Energy Measurements (2/2)

Calorimeter and Material have to be designed such that the particles leave their full energy inside \rightarrow depends on interaction

- Electromagnetic calorimeter: mainly energy of e^+ , e^- , γ
- Hadronic calorimeter for particles which interact mainly which hadronic interactions with the atomic cores of the absorber material: π^\pm , p^\pm , n , ...

Electromagnetic Calorimeters (1/3)

- High energetic electrons and photons create showers in the absorber-material:
 - pair-production, bremsstrahlung, pair-production, ...
- until critical energy E_C is reached
- The number of particles doubles in each step X_0 as long as EE_C :
 - $N(t) \approx 2^t$, where t is the number of steps in X_0
- with an average energy per particle of
- $E(t) = E_0/N(t) = E_0 \cdot 2^{-t}$



Electromagnetic Calorimeters (2/3)

The maximum of the shower is reached at $t_{max} \approx \ln(E_0/E_C)$. The total number of particles S in the shower can be calculated by

$$S = \sum_{t=0}^{t_{max}} N(t) = \sum 2^t = 2^{t_{max}+1} - 1 \approx 2^{t_{max}+1}$$

this gives

$$S = 2 \frac{E_0}{E_S} \sim E$$

Therefore we just have to measure the number of particles in the shower to determine the initial energy. Since S is a large number and its uncertainty is given by $\sqrt{(S)}$, we expect for the relative energy resolution

$$\frac{\sigma}{E} \sim \sqrt{\frac{t}{E_0}}$$

i.e. it gets more precise for higher energies!

Electromagnetic Calorimeters (3/3)

How does it work? Lets look at a sampling calorimeter:

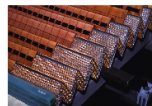
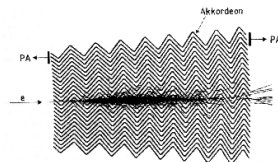
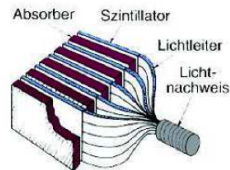
Periodic structure of passive absorber material (Keep in mind:

Bremsstrahlung cross-section goes with $\sigma \sim Z^2$)

- lead
- copper
- ...

And active materials for the detection of the shower particles

- plastic scintillators
- liquid argon (as 'ionisation chamber')
- ...

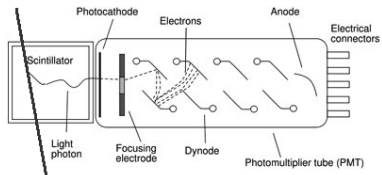


PA: Vorverstärker

Scintillators

Scintillators are particle detectors, in which ionizing particles cause light emissions in a transparent medium.

- Anorganic Scintillators
 - cristal structure, doped with activators
 - Ionizing particles create free electron-hole pairs
 - excitations travel in the lattice to an activator which gets excited and emitts light to get back into its ground state
 - → just count photons with the help of a photo-multiplier
- Organic Scintillators



Hadronic Calorimeters (1/2)

Hadronic calorimeters are designed for the energy measurement of hadrons

- Similar principle as electromagnetic calorimeters, but showers are produced via strong interaction
 - incoming hadrons interact with atomic nuclei $\rightarrow n, p, \pi^0, \pi^\pm$
 - $\pi^0 \rightarrow \gamma\gamma$: also electromagnetic shower
 - $\approx 20\%$ of energy needed to break up nuclear binding energy
- nuclear absorption lengths λ instead of radiation length X_0
 - Typical values for λ : $\lambda_{Fe} \approx 17.1cm$, $\lambda_{Pb} \approx 18.5cm$
 - Typical values for X_0 : $X_{0,Fe} \approx 1.76cm$, $X_{0,Pb} \approx 0.56cm$
 - \rightarrow hadronic calorimeters must be much larger

Hadronic Calorimeters (2/2)

Energy Resolution

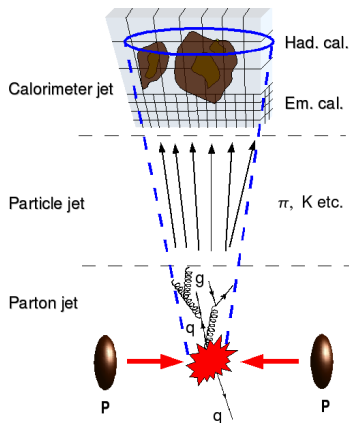
- EC: $\frac{\sigma}{E} = \frac{3-5\%}{\sqrt{E[\text{GeV}]}}$
- HC: $\frac{\sigma}{E} = \frac{50-90\%}{\sqrt{E[\text{GeV}]}}$

Worse resolution of HC due to:

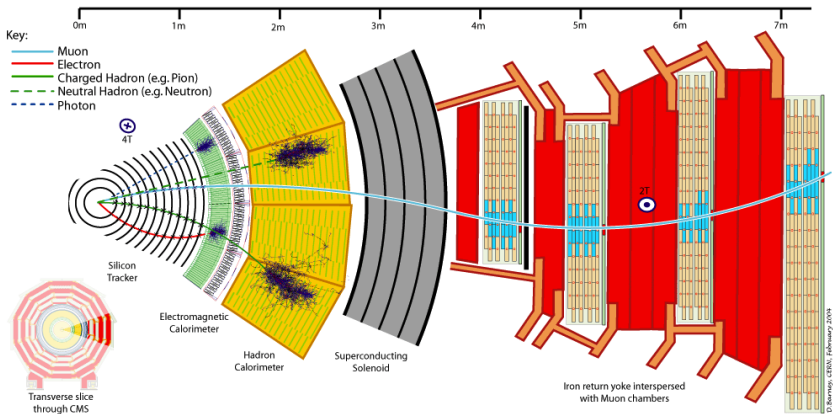
- leakage
- some shower-particles (ν, μ) escape the calorimeter
- breaking up binding energies

Important

HC do not measure directly quarks/gluons but **Particle Jets**



Basic Principle of an LHC Detector



What do we measured?

We measure

- energy / momentum of particles
- the 'initial' position (separate primary vertex and secondary vertex)
- trajectory / direction
- \rightarrow 4-momentum: $P = (E, -\vec{p})$
- Note: $P^2 = E^2 - |\vec{p}|^2 = m^2$. @LHC: $E^2 \gg m^2$

We distinguish

- electrons
- photons
- muons
- jets (from quarks, gluons,): uds-jet? b-jet? c-jet?

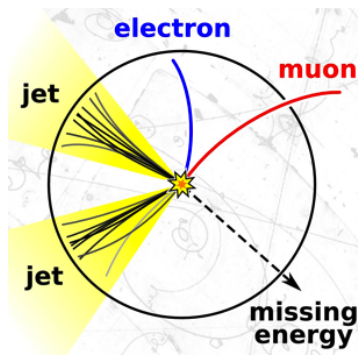
How to detect neutrinos

Neutrinos leave detector 'unseen',
but carry away energy

- total momentum in z-direction of interacting partons not known
 - Different to e^+/e^- collisions
- But $\sum E_T^{initial} = 0 = \sum_{all} E_T^{final}$
- Hence we define

$$\vec{E}_T^{miss} = - \sum_i \vec{E}_T^i$$

- Use \vec{E}_T^{miss} as estimate for neutrinos
 - No z-information!



The ATLAS Detector (1/2)

ATLAS and CMS are 4π detectors

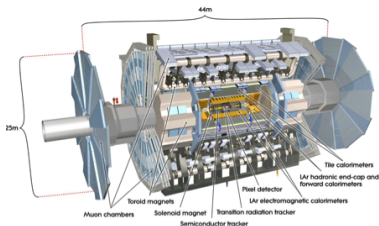
- Barrel- and Endcap-Regions

Inner Detector

- $|\eta| < 2.5$, solenoid $B = 2T$
- Si Pixels, Si strips, TRT
- Tracking and vertexing
- Resolution: $\sigma/p_T \sim 3.8 \times 10^{-4} p_T [GeV] \oplus 0.015$

EM calorimeter

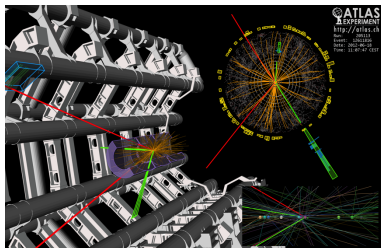
- $|\eta| < 3.2$
- LAr/Pb accordion structure e/γ
- E-resolution: $\sigma/E \sim 10\%/\sqrt{E}$



The ATLAS Detector (2/2)

HAD calorimeter

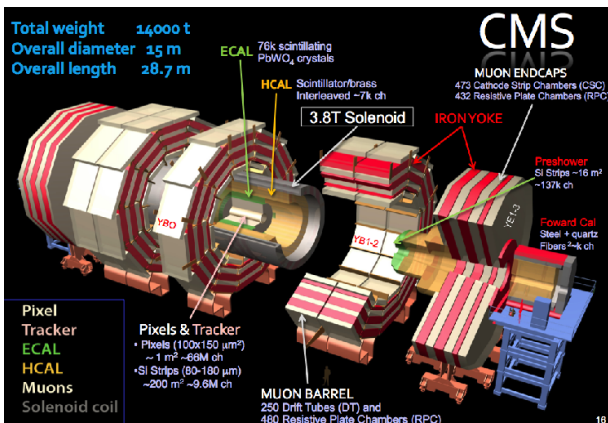
- $|\eta| < 3.2$
- Forward Calo. $|\eta| < 4.8$
- Scint./Fe tiles in the central, W(Cu)/LAr in fwd region
- Trigger, jets + missing E_t
- E-resolution:
 $\sigma/E \sim 50\%/sqrt{E} \oplus 0.03$



Muon Spectrometer

- $|\eta| < 2.7$
- Toroid B-Field (0.5 T)
- Muon Momentum resolution
< 10% up to $\sim 1\text{ TeV}$

The CMS Detector



(stolen from a fantastic CMS talk)

ALICE and LHCb Detector

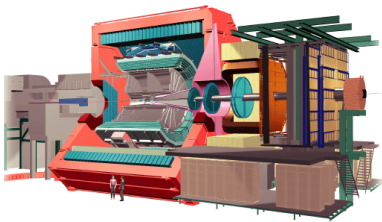
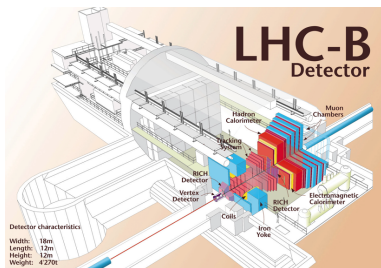
ALICE Detector

- designed to study the quark-gluon plasma during the ion-ion collisions at the LHC

LHCb Detector

- designed to study the CP-Violation in b-hadrons
- understand the Matter-Antimatter asymmetry of the universe?
- study and measure the

Both highly interesting physics programs - maybe a lecture in 2013?



Summary of Lecture 2

- ATLAS and CMS measure the 4-momentum of particles
- o) electrons, photons
 - o) muons
 - o) jets (+ some additional information)
 - o) indirectly neutrinos (or weakly interacting particles)