

QCD@LHC for beginners

Lesson 5

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Outline

- Lesson 5

- Examples

- Drell-Yan process
 - Higgs production/decay

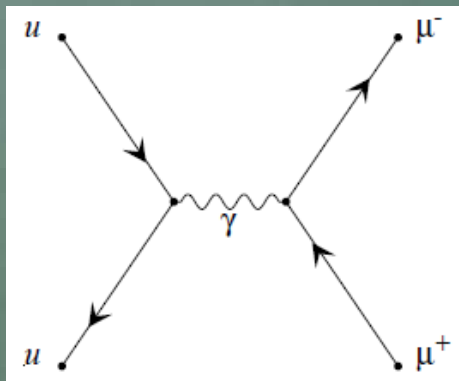
Crosssection calculations
understand a behavior

- Non-perturbative effect

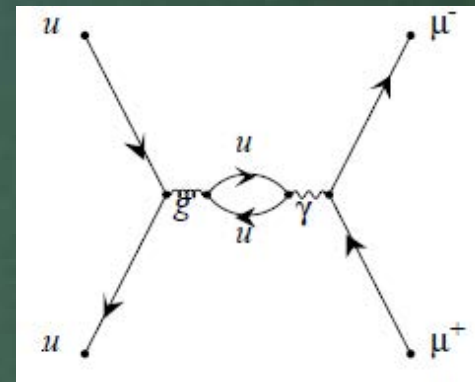
- Hadronization



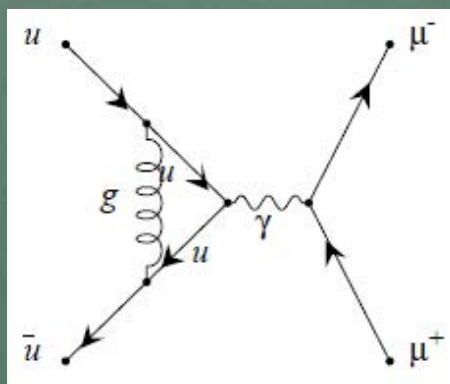
Examples/Drell-Yan



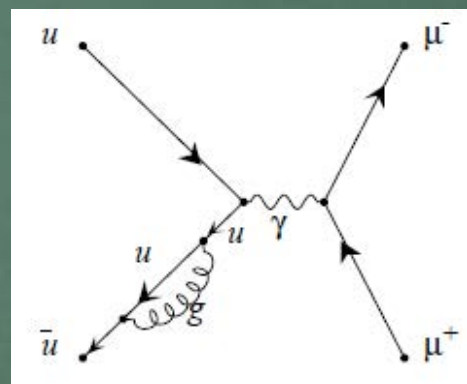
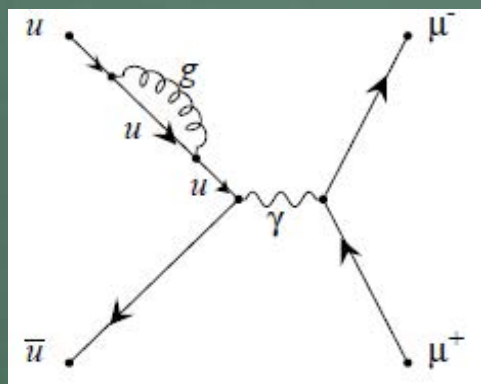
Tree diagram



$\equiv 0$

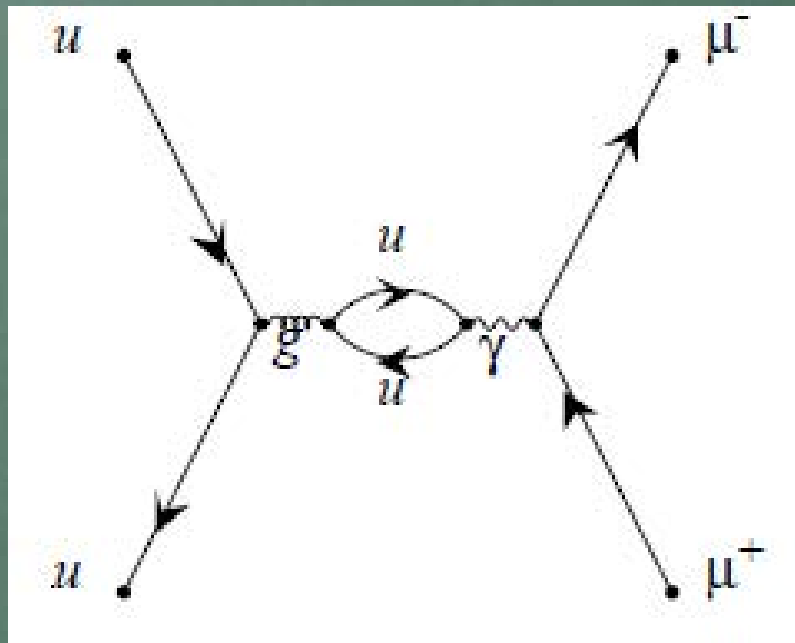


vertex correction



quark self-energy

Examples/Drell-Yan



$$= 0$$

Why?

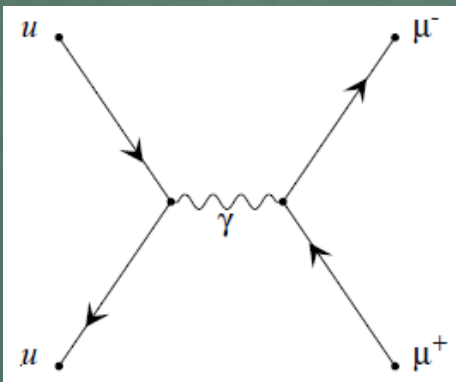


Examples/Drell-Yan

Tree cross section

$$d\sigma_0 = \frac{1}{4N_c^2} \frac{1}{2s} |\mathcal{M}_0|^2 d\Phi_2$$

$$d\Phi_2^{(d)}(P \rightarrow p_1 + p_2) = \frac{d^{d-1} \vec{p}_1}{(2\pi)^{d-1} 2p_1^0} \frac{d^{d-1} \vec{p}_2}{(2\pi)^{d-1} 2p_2^0} (2\pi)^d \delta^{(d)}(P - p_1 - p_2)$$



2

=

$$|\mathcal{M}_0|^2$$



Examples/Drell-Yan

Tree cross section

Phase space:

$$P \rightarrow p_1 + p_2,$$

$$P = (W, \vec{0}^{d-1}), \quad P^2 = W^2$$

=1

$$d\Phi_2^{(d)}(P \rightarrow p_1 + p_2) = \frac{d^{d-1} \vec{p}_1}{(2\pi)^{d-1} 2p_1^0} \frac{d^{d-1} \vec{p}_2}{(2\pi)^{d-1} 2p_2^0} (2\pi)^d \delta^{(d)}(P - p_1 - p_2)$$

$$d\Phi_2^{(d)} = \frac{1}{(2\pi)^{d-1}} \frac{d^{d-1} \vec{p}_1}{2p_1^0} \boxed{2p_2^0 dp_2^0 \delta(p_2^2 - m_2^2)} \frac{d^{d-1} \vec{p}_2}{(2\pi)^{d-1} 2p_2^0} (2\pi)^d \delta^{(d)}(P - p_1 - p_2)$$

$$= \frac{1}{(2\pi)^{d-2}} \frac{d^{d-1} \vec{p}_1}{2p_1^0} \delta(W^2 - 2Wp_1^0 + m_1^2 - m_2^2)$$

angle integration

$$d^{d-1} \vec{p}_1 = |\vec{p}_1|^{d-2} d|\vec{p}_1| d\Omega_{d-1}$$

d-dim. polar coordinate

Examples/Drell-Yan

angle integration

ϕ -symmetry

$$d\Omega_{d-1} = d\theta \sin^{d-3}\theta d\Omega_{d-2}$$

$$= d\theta \sin^{d-3}\theta \frac{2\pi^{\left(\frac{d-2}{2}\right)}}{\Gamma\left(\frac{d-2}{2}\right)}$$

$$d\Phi_2^{(d)} = \frac{1}{(2\pi)^{d-2}} \frac{\pi^{\left(\frac{d-2}{2}\right)}}{\Gamma\left(\frac{d-2}{2}\right)} (p_1^0)^{d-3} \beta_1^{d-3} dp_1^0 \delta(W^2 - 2W p_1^0 + m_1^2 - m_2^2) d\cos\theta (1 - \cos^2\theta)^{\frac{d-4}{2}}$$

p_1^0 integration

0

$$\beta_1 = \frac{|\vec{p}_1|}{p_1^0}$$

Examples/Drell-Yan

p_1^0 integration

$$p_1^0 = \frac{W^2 + m_1^2 - m_2^2}{2W}$$

$$\beta = \frac{|\vec{p}_1|}{p_1^0}$$

$$d\Phi_2^{(d)} = \frac{1}{(2\pi)^{d-2}} \frac{\pi^{\left(\frac{d-2}{2}\right)}}{\Gamma\left(\frac{d-2}{2}\right)} (p_1^0)^{d-3} \beta_1^{d-3} \frac{1}{2W} d\cos\theta (1 - \cos^2\theta)^{\frac{d-4}{2}}$$

$m_1 = m_2 = 0$

$$d\Phi_2^{(4+2\epsilon_{IR})} = \frac{1}{(16\pi)^{1+\epsilon_{IR}}} \frac{1}{\Gamma(1+\epsilon_{IR})} W^{2\epsilon_{IR}} d\cos\theta (1 - \cos^2\theta)^{\epsilon_{IR}}$$


Examples/Drell-Yan

Tree cross section

Phase space:

$$P \rightarrow p_1 + p_2,$$
$$P = (W, \vec{0}^{d-1}), \quad P^2 = W^2$$

$$d\Phi_2^{(d)}(P \rightarrow p_1 + p_2) = \frac{d^{d-1} \vec{p}_1}{(2\pi)^{d-1} 2p_1^0} \frac{d^{d-1} \vec{p}_2}{(2\pi)^{d-1} 2p_2^0} (2\pi)^d \delta^{(d)}(P - p_1 - p_2)$$


$$m_1 = m_2 = 0$$

$$d\Phi_2^{(4+2\epsilon_{IR})} = \frac{1}{(16\pi)^{1+\epsilon_{IR}}} \frac{1}{\Gamma(1+\epsilon_{IR})} W^{2\epsilon_{IR}} d\cos\theta (1 - \cos^2\theta)^{\epsilon_{IR}}$$

Exercise: Confirm above result.

Examples/Drell-Yan

Tree cross section

Matrix element:

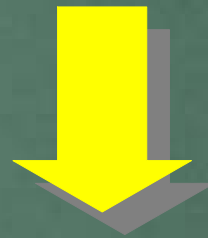
$$\left| \begin{array}{c} u \\ \searrow \\ \gamma \\ \swarrow \\ u \end{array} \begin{array}{c} \mu^- \\ \swarrow \\ \gamma \\ \searrow \\ \mu^+ \end{array} \right|^2 = |\mathcal{M}_0|^2$$

$$\begin{aligned} |\mathcal{M}_0|^2 &= \frac{e^4}{s^2} \text{Tr}[p_1 \gamma^\mu p_2 \gamma^\nu] \text{Tr}[k_1 \gamma^\mu k_2 \gamma^\nu] \\ &= 4 \frac{e^4}{s^2} \left[(d-4)(p_1 \cdot p_2)(k_1 \cdot k_2) \right. \\ &\quad \left. + 2(p_1 \cdot k_1)(p_2 \cdot k_2) + 2(p_1 \cdot k_2)(p_2 \cdot k_1) \right] \end{aligned}$$

Examples/Drell-Yan

Tree cross section

$$\int d\Phi_2^d [(d-4)(p_1 \cdot p_2)(k_1 \cdot k_2) + 2(p_1 \cdot k_1)(p_2 \cdot k_2) + 2(p_1 \cdot k_2)(p_2 \cdot k_1)]$$
$$= \frac{2 \pi^{3/2} s^2 (\epsilon + 1) \csc(\pi \epsilon)}{\Gamma(-\epsilon - 1) \Gamma(\epsilon + \frac{5}{2})}$$



$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2 \frac{(1 + \epsilon_{IR})^2}{(3 + 2\epsilon_{IR})} \left(\frac{s}{4\pi}\right)^{\epsilon_{IR}} \frac{1}{\Gamma(1 + \epsilon_{IR})}$$

Examples/Drell-Yan

Tree cross section

$$\begin{aligned} |\mathcal{M}_0|^2 &= \frac{e^4}{s^2} \text{Tr}[\not{p}_1 \gamma^\mu \not{p}_2 \gamma^\nu] \text{Tr}[\not{k}_1 \gamma^\mu \not{k}_2 \gamma^\nu] \\ &= 4 \frac{e^4}{s^2} [(d-4)(p_1 \cdot p_2)(k_1 \cdot k_2) \\ &\quad + 2(p_1 \cdot k_1)(p_2 \cdot k_2) + 2(p_1 \cdot k_2)(p_2 \cdot k_1)] \end{aligned}$$

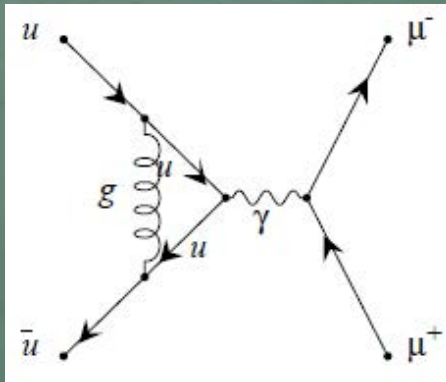
$$\begin{aligned} &\int d\Phi_2^d [(d-4)(p_1 \cdot p_2)(k_1 \cdot k_2) + 2(p_1 \cdot k_1)(p_2 \cdot k_2) + 2(p_1 \cdot k_2)(p_2 \cdot k_1)] \\ &= \frac{2 \pi^{3/2} s^2 (\epsilon + 1) \csc(\pi \epsilon)}{\Gamma(-\epsilon - 1) \Gamma(\epsilon + \frac{5}{2})} \end{aligned}$$

$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2 \frac{(1 + \epsilon_{IR})^2}{(3 + 2\epsilon_{IR})} \left(\frac{s}{4\pi}\right)^{\epsilon_{IR}} \frac{1}{\Gamma(1 + \epsilon_{IR})}$$

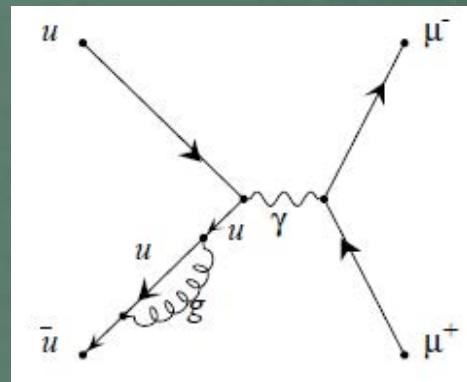
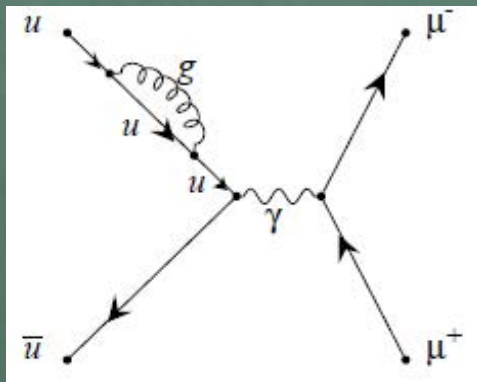
Exercise: Confirm above result.

Examples/Drell-Yan

NLO cross section



vertex correction

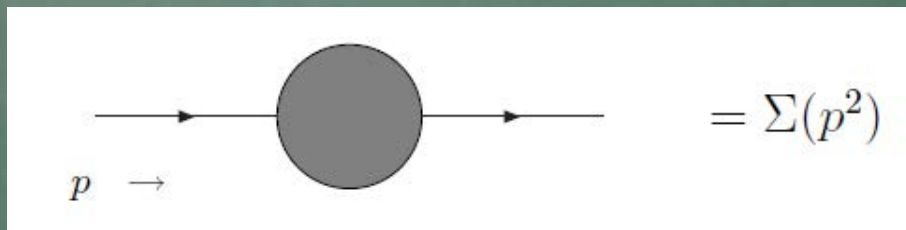


quark self-energy

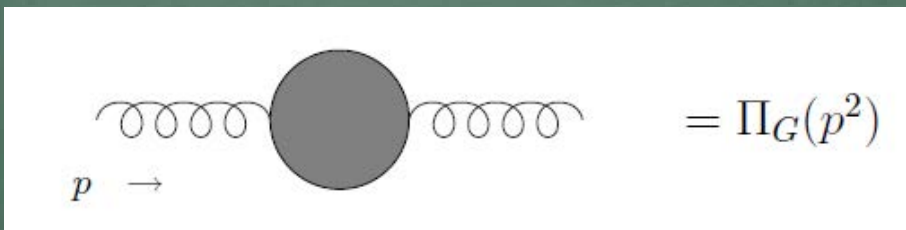


Examples/Drell-Yan

NLO cross section



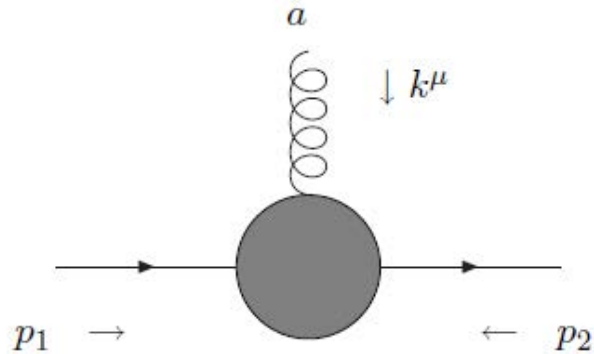
on-/off-shell	self-energy
$p^2 \neq 0$	$\Sigma(p^2) = C_F \frac{\alpha_s}{4\pi} (1 - \ln \frac{-p^2}{\mu^2})$
$p^2 = 0$	$\Sigma(0) = C_F \frac{\alpha_s}{4\pi} \frac{1}{\epsilon_{IR}}$



on-/off-shell	self-energy
$p^2 \neq 0$	$\Pi_G(p^2) = -C_F \frac{\alpha_s}{4\pi} (\frac{31}{9} - \frac{5}{3} \ln \frac{-p^2}{\mu^2}) + T_R N_f \frac{\alpha_s}{4\pi} (\frac{20}{9} - \frac{4}{3} \ln \frac{-p^2}{\mu^2})$
$p^2 = 0$	$\Pi_G(0) = -C_F \frac{\alpha_s}{4\pi} \frac{1}{\epsilon_{IR}} + T_R N_f \frac{\alpha_s}{4\pi} \frac{4}{3} \frac{1}{\epsilon_{IR}}$

Examples/Drell-Yan

NLO cross section



$$\begin{aligned}
 = \Lambda_\mu^a(p_1, p_2, k) = & g\mathbf{T}^a(C_F - \frac{1}{2}C_G)\frac{\alpha_s}{4\pi}(\mathcal{F}_1^I\gamma^\mu + \mathcal{F}_2^I\frac{p_j^\mu k}{-p_j^2}) \\
 & + g\mathbf{T}^a\frac{1}{2}C_G\frac{\alpha_s}{4\pi}(\mathcal{F}_1^I\gamma^\mu + \mathcal{F}_2^I\frac{p_j^\mu k}{-p_j^2})
 \end{aligned}$$

$$k^2 = 0, p_i^2 = 0, p_j^2 = q^2 \neq 0, \text{ and } L = \ln \frac{-q^2}{\mu^2},$$

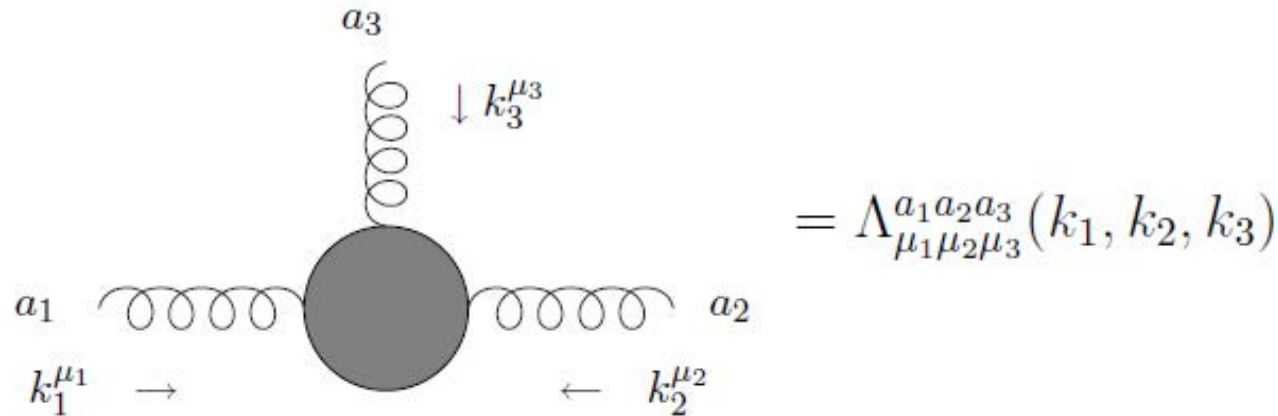
\mathcal{F}_1^I	$\frac{2}{\epsilon_{IR}} + L - 4$
\mathcal{F}_2^I	$\frac{4}{\epsilon_{IR}} + 4L - 10$
\mathcal{F}_1^{II}	$-\frac{2}{\epsilon_{IR}^2} + \frac{3-2L}{\epsilon_{IR}} + \frac{\pi^2}{6} - L$
\mathcal{F}_2^{II}	$-\frac{2}{\epsilon_{IR}^2} - \frac{2L}{\epsilon_{IR}} + \frac{12+\pi^2-6L^2}{6}$

$$k^2 = q^2, p_1^2 = p_2^2 = 0, \text{ and } L = \ln \frac{-q^2}{\mu^2}.$$

\mathcal{F}_1^I	$-\frac{2}{\epsilon_{IR}^2} - \frac{2L-4}{\epsilon_{IR}} - 8 + \frac{\pi^2}{6} + 3L - L^2$
\mathcal{F}_2^I	0
\mathcal{F}_1^{II}	$\frac{4}{\epsilon_{IR}} - 2 + L$
\mathcal{F}_2^{II}	0

Examples/Drell-Yan

NLO cross section



$$k_1^2 = k_2^2 = 0, k_3^2 = q^2 \neq 0, L = \ln \frac{-q^2}{\mu^2}$$

$$\Lambda_{\mu_1 \mu_2 \mu_3}^{a_1 a_2 a_3}(k_1, k_2, k_3) = -igf^{a_1 a_2 a_3} \frac{\alpha_s}{4\pi} \left[\frac{C_G}{2} (\mathcal{G}_1 + \mathcal{G}_2 + \mathcal{G}_3) + \frac{N_f}{2} \mathcal{G}_4 \right],$$

$$\mathcal{G}_i = \sum_{j=1}^3 c_{ij} \mathcal{P}_j$$

\mathcal{G}_1 :gluon loop, \mathcal{G}_2 :ghost loop, \mathcal{G}_3 :gluon loop (fish type), \mathcal{G}_4 :quark loop

Examples/Drell-Yan

NLO cross section

- $\mathcal{P}_1^{\mu_1\mu_2\mu_3} = (k_1 - k_2)^{\mu_3} g^{\mu_1\mu_2}$
- $\mathcal{P}_2^{\mu_1\mu_2\mu_3} = k_1^{\mu_2} g^{\mu_1\mu_3} - k_2^{\mu_2} g^{\mu_2\mu_3}$
- $\mathcal{P}_2^{\mu_1\mu_2\mu_3} = \frac{k_2^{\mu_1} k_1^{\mu_2} (k_1 - k_2)^{\mu_3}}{q^2}$

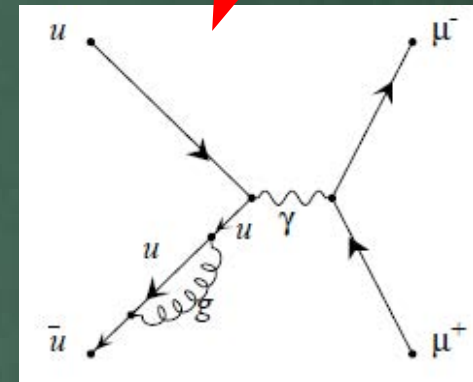
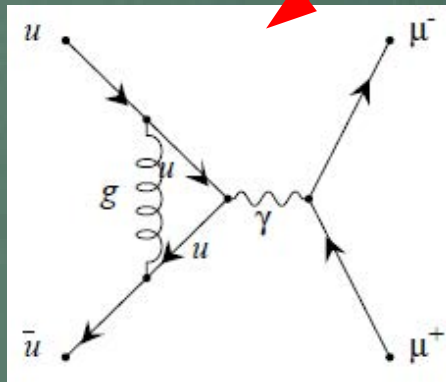
c_{ij}	$j = 1$	$j = 2$	$j = 3$
$i = 1$	$-\frac{3}{2\varepsilon_{IR}^2} - \frac{3}{2\varepsilon_{IR}}(-5 + L)$ $-\frac{1}{12}(103 - 51L + 9L^2) + \frac{\pi^2}{8}$	$\frac{2}{\varepsilon_{IR}^2} - \frac{1}{2\varepsilon_{IR}}(19 - 4L)$ $+\frac{1}{3}(19 - 9L + 3L^2) - \frac{\pi^2}{6}$	$-\frac{3}{2}$
$i = 2$	$-\frac{11-3L}{36}$	$\frac{8-3L}{18}$	$\frac{1}{6}$
$i = 3$	$-\frac{9}{2\varepsilon_{IR}}$	$\frac{2}{2\varepsilon_{IR}} + \frac{18-19L}{2}$	0
$i = 4$	$\frac{14-12L}{9}$	$-\frac{40-24}{9}$	$\frac{4}{3}$

Examples/Drell-Yan

NLO cross section

$$\sigma_{DY}^{loop} = \sigma_{DY}^0 C_F \frac{\alpha_s(\mu)}{2\pi} \left[-\frac{2}{\epsilon_{IR}^2} - \frac{2L-4}{\epsilon_{IR}} - 8 + \frac{\pi^2}{6} + 3L - L^2 - \frac{1}{\epsilon_{IR}} - \frac{1}{\epsilon_{IR}} \right]$$

$$L = \ln \left(\frac{s}{\mu_F^2} \right)$$



Examples/Drell-Yan

NLO cross section

$$\sigma_{DY}^{loop} = \sigma_{DY}^0 C_F \frac{\alpha_s(\mu)}{2\pi} \left[-\frac{2}{\epsilon_{IR}^2} - \frac{2L-4}{\epsilon_{IR}} - 8 + \frac{\pi^2}{6} + 3L - L^2 - \frac{1}{\epsilon_{IR}} - \frac{1}{\epsilon_{IR}} \right]$$

$$\sigma_{coll} = \sigma_0(s) \frac{\alpha_s}{2\pi} C_F \left[\frac{2}{\epsilon_{IR}^2} + \frac{2L-3}{\epsilon_{IR}} - \frac{\pi^2}{2} + L^2 \right] \quad L = \ln \left(\frac{s}{\mu_F^2} \right)$$

$$+ 2 \int_0^1 dx \sigma_0(xs) \phi(x, \epsilon_{IR})$$

$$+ 2 C_F \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) \left[L \frac{1+x^2}{(1-x)_+} + 2 \frac{(1+x^2) \ln(1-x)}{(1-x)_+} - \frac{1+x^2}{1-x} \ln x \right]$$

$$\phi(x, \epsilon_{IR}) = \frac{1}{\epsilon_{IR}} C_F \frac{\alpha_s}{2\pi} P(x) = \frac{1}{\epsilon_{IR}} C_F \frac{\alpha_s}{2\pi} \frac{1+x^2}{(1-x)_+}$$

Examples/Drell-Yan

NLO cross section

$$\sigma_{DY}^{loop} = \sigma_{DY}^0 C_F \frac{\alpha_s(\mu)}{2\pi} \left[-\frac{2}{\epsilon_{IR}^2} - \frac{2L-4}{\epsilon_{IR}} - 8 + \frac{\pi^2}{6} + 3L - L^2 - \frac{1}{\epsilon_{IR}} - \frac{1}{\epsilon_{IR}} \right]$$

$$\sigma_{coll} = \sigma_0(s) \frac{\alpha_s}{2\pi} C_F \left[\frac{2}{\epsilon_{IR}^2} + \frac{2L-3}{\epsilon_{IR}} - \frac{\pi^2}{2} + L^2 \right] \quad L = \ln \left(\frac{s}{\mu_F^2} \right)$$

$$+ 2 \int_0^1 dx \sigma_0(xs) \phi(x, \epsilon_{IR})$$

$$+ 2 C_F \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) \left[L \frac{1+x^2}{(1-x)_+} + 2 \frac{(1+x^2) \ln(1-x)}{(1-x)_+} - \frac{1+x^2}{1-x} \ln x \right]$$

$$\phi(x, \epsilon_{IR}) = \frac{1}{\epsilon_{IR}} C_F \frac{\alpha_s}{2\pi} P(x) = \frac{1}{\epsilon_{IR}} C_F \frac{\alpha_s}{2\pi} \frac{1+x^2}{(1-x)_+}$$

Examples/Drell-Yan

NLO cross section

$$\sigma_{DY}^{NLO} = \sigma_{DY}^0 \left[1 + C_F \frac{\alpha_s(\mu)}{2\pi} \left(3 \ln \left(\frac{s}{\mu_F^2} \right) - \frac{\pi^2}{3} - 8 \right) \right]$$

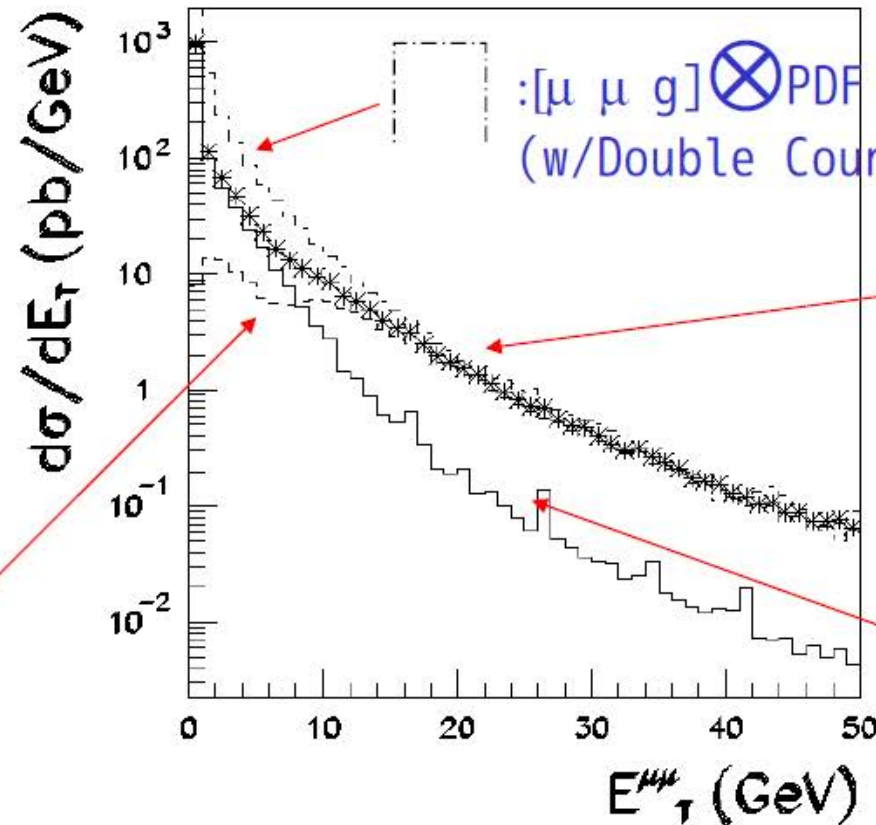
• Process :

$uu \rightarrow \mu^+ \mu^- (+ \text{gluon})$
in pp collision

Cuts:

$\sqrt{s_{\mu\mu}} > 40 \text{ GeV}$

$k_T^g > 1 \text{ GeV}$



$:[\mu \mu g\text{-LL}] \otimes \text{PS}$

$:[\mu \mu g] \otimes \text{PDF}$
(w/Double Counting)

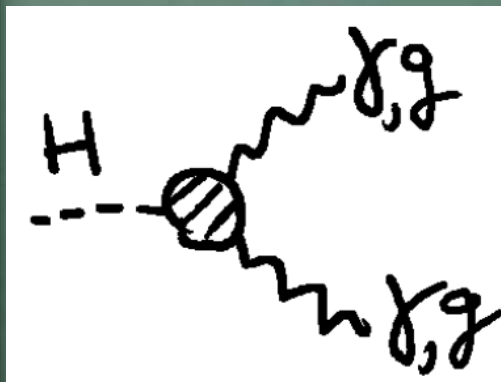
$:[\mu \mu g(\text{NLO})] \otimes \text{PS}$

$[\mu \mu g] + [\mu \mu g(\text{NLO})]$

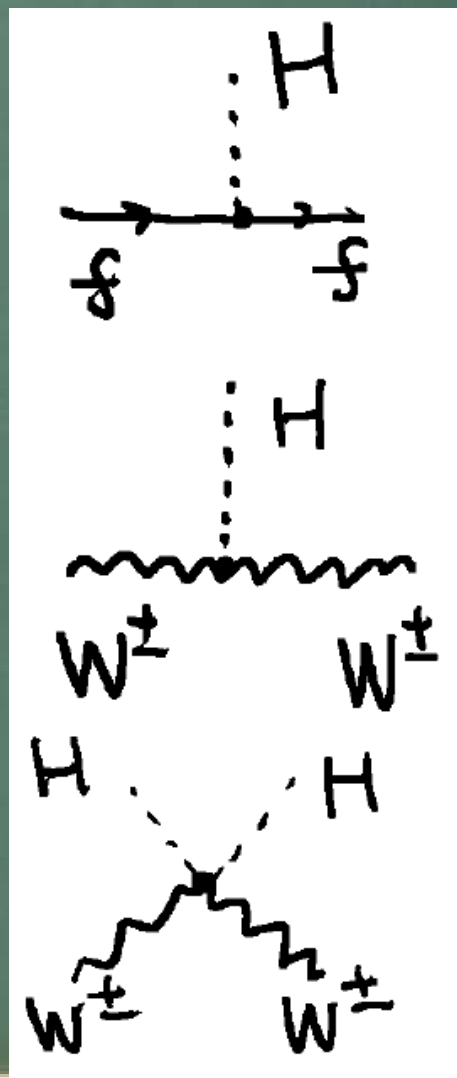
$:[\mu \mu (t+v+c)]/\text{PS}$

Examples/Higgs $\rightarrow \tau \tau$

Feynman Rule



$$g_W = \frac{e}{\sin \theta_W}$$



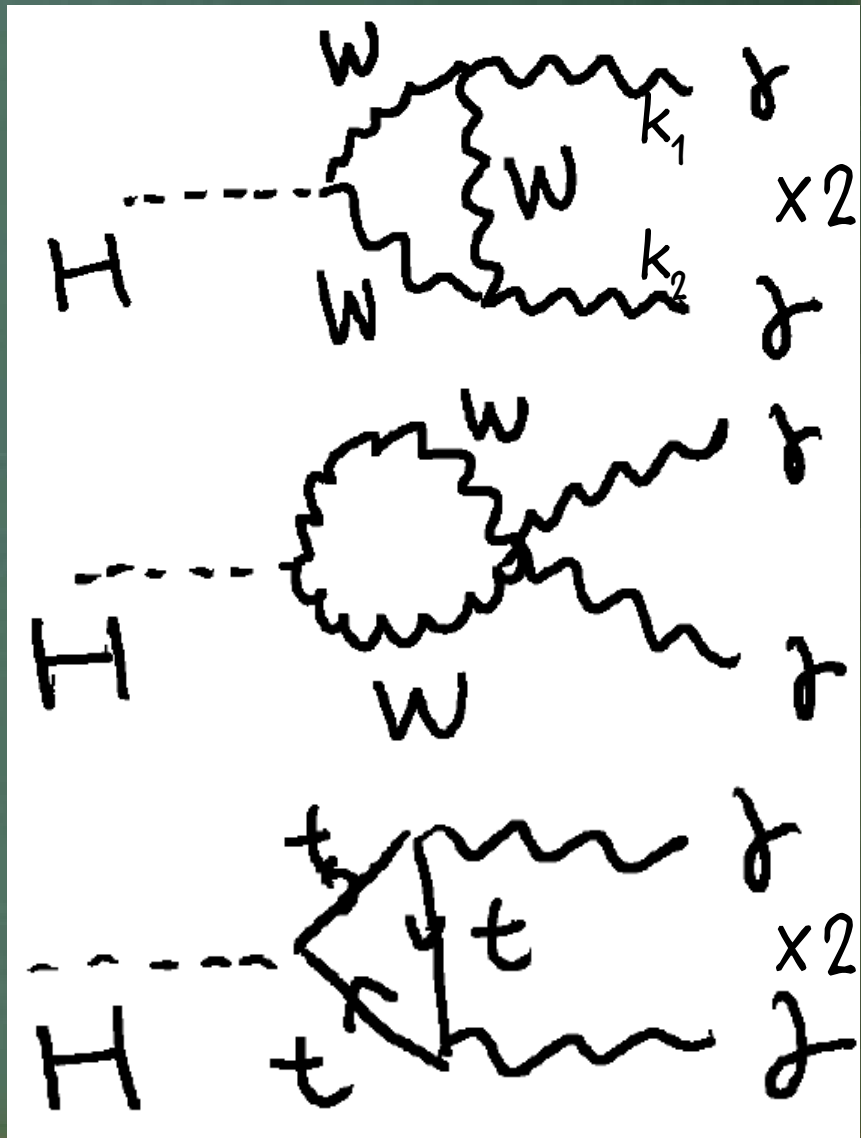
$$= -i \frac{1}{2} \frac{g_W}{m_W} m_f$$

$$= i g_W m_W g^{\mu\nu}$$

$$= i \frac{g_W^2}{2} g^{\mu\nu}$$

Examples/Higgs $\rightarrow \gamma \gamma$

$$\Gamma(H \rightarrow \gamma\gamma) = |F|^2 \left(\frac{\alpha}{4\pi}\right)^2 \frac{G_F m_H^3}{8\sqrt{2}\pi}$$




Examples/Higgs $\rightarrow \tau \tau$

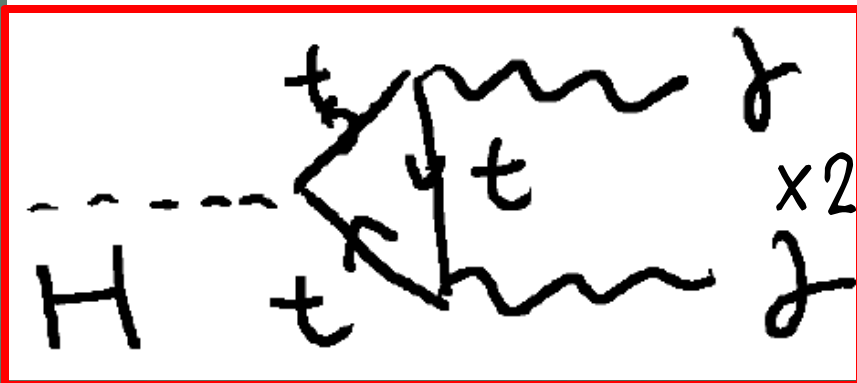
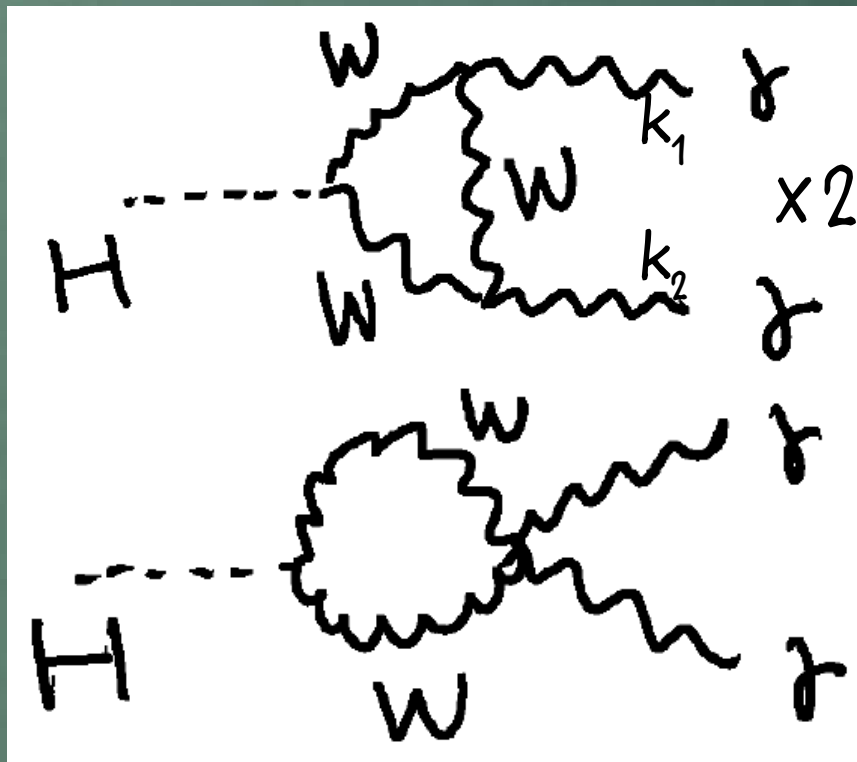
$$F = F_W(\beta_W) + \sum_f N_c Q_f^2 F_f(\beta_f)$$

$$F_W(\beta) = 2 + 3\beta + 3\beta(2 - \beta)f(\beta)$$
$$F_f(\beta) = -2\beta [1 + (1 - \beta)f(\beta)]$$

$$\beta_W = \frac{4m_W^2}{m_H^2}, \quad \beta_f = \frac{4m_f^2}{m_H^2}$$

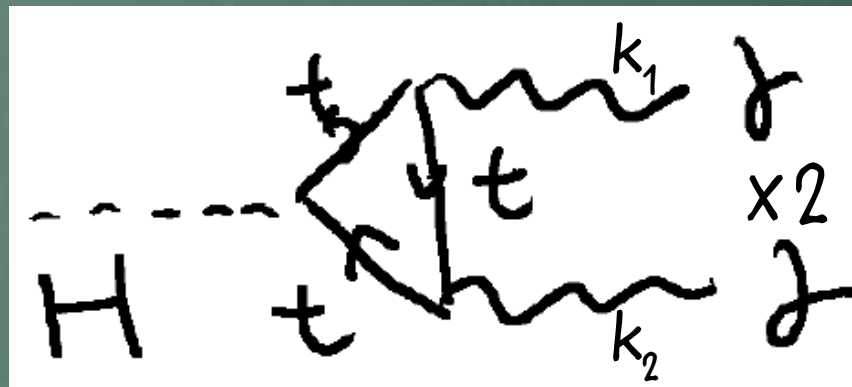
$$f(\beta) = \begin{cases} \arcsin^2(\beta^{-\frac{1}{2}}) & \text{for } \beta \geq 1 \\ -\frac{1}{4} \left[\ln \frac{1+\sqrt{1-\beta}}{1-\sqrt{1-\beta}} - i\pi \right]^2 & \text{for } \beta < 1 \end{cases}$$


Examples/Higgs $\rightarrow \tau \tau$



Examples/Higgs $\rightarrow \tau \tau$

$$d\Gamma = \frac{(2\pi)^4}{2M} |\mathcal{M}|^2 d\Phi_n (P; p_1, \dots, p_n)$$

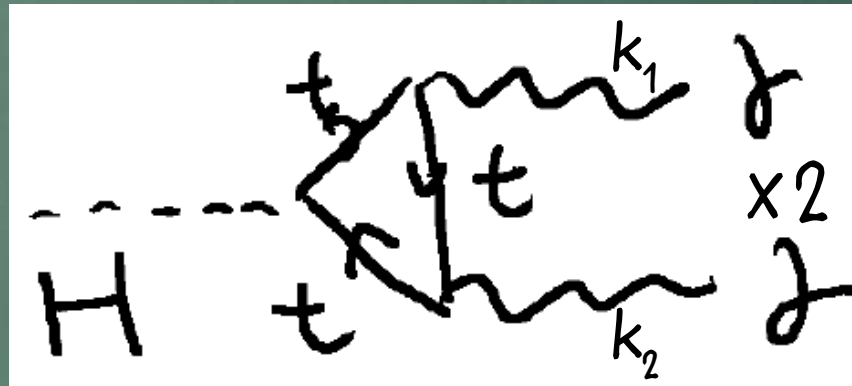


$$2\mathcal{M}_t = -ie^3 \frac{m_t}{m_W \sin \theta_W} \int \frac{d^4 k}{(2\pi)^4 i} \text{Tr} \left[\frac{i}{\not{k} + \not{k}_1 - m_t} i\not{\epsilon}(k_1) \frac{i}{\not{k} - m_t} i\not{\epsilon}(k_2) \frac{i}{\not{k} + \not{k}_2 - m_t} \right]$$

$$\sim e^3 \frac{m_t}{m_W \sin \theta_W} \int \frac{d^4 k}{(2\pi)^4 i} \text{Tr} \left[\frac{(\not{k} + m_t) \not{\epsilon}(k_1) (\not{k} + m_t) \not{\epsilon}(k_2) (\not{k} + m_t)}{(k^2 - m_t^2)^3} \right]$$

Examples/Higgs $\rightarrow \tau \tau$

$$d\Gamma = \frac{(2\pi)^4}{2M} |\mathcal{M}|^2 d\Phi_n (P; p_1, \dots, p_n)$$



$$2\mathcal{M}_t = -ie^3 \frac{m_t}{m_W \sin \theta_W} \int \frac{d^4 k}{(2\pi)^4 i} \text{Tr} \left[\frac{i}{\not{k} + \not{k}_1 - m_t} i\not{\epsilon}(k_1) \frac{i}{\not{k} - m_t} i\not{\epsilon}(k_2) \frac{i}{\not{k} + \not{k}_2 - m_t} \right]$$

$$\sim e^3 \frac{m_t}{m_W \sin \theta_W} \int \frac{d^4 k}{(2\pi)^4 i} \text{Tr} \left[\frac{(\not{k} + m_t) \not{\epsilon}(k_1) (\not{k} + m_t) \not{\epsilon}(k_2) (\not{k} + m_t)}{(k^2 - m_t^2)^3} \right]$$

$$\int \frac{d^n k}{(2\pi)^n i} \frac{(k^2)^r}{(k^2 - \mathcal{D})^m} = \frac{(-1)^{r-m}}{(16\pi)^{n/4}} \mathcal{D}^{r-m+n/2} \frac{\Gamma(r+n/2) \Gamma(m-r-n/2)}{\Gamma(n/2) \Gamma(m)}$$

Examples/Higgs $\rightarrow \tau \tau$

$$d\Gamma = \frac{(2\pi)^4}{2M} |\mathcal{M}|^2 d\Phi_n (P; p_1, \dots, p_n)$$

$$\left\{ e^3 \frac{m_t}{m_W \sin \theta_W} m_t^2 m_H^2 \right\}^2$$



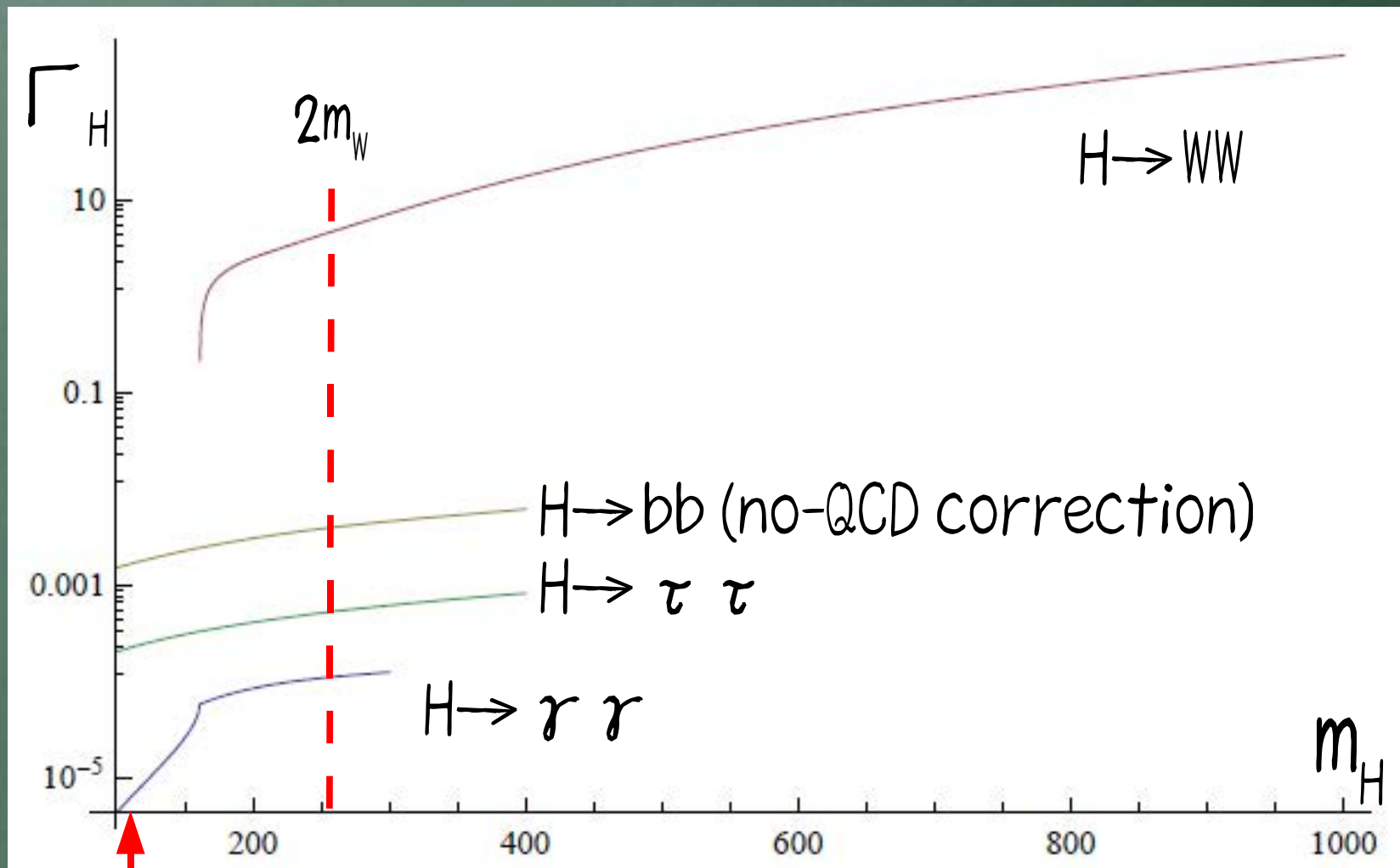
Examples/Higgs $\rightarrow \gamma \gamma$

$$d\Gamma = \frac{(2\pi)^4}{2M} |\mathcal{M}|^2 d\Phi_n (P; p_1, \dots, p_n)$$

$$\left\{ e^3 \frac{m_t}{m_W \sin \theta_W} m_t^2 m_H^2 \right\}^2$$

$$\Gamma(H \rightarrow \gamma\gamma) = |F|^2 \left(\frac{\alpha}{4\pi} \right)^2 \frac{G_F m_H^3}{8\sqrt{2}\pi}$$

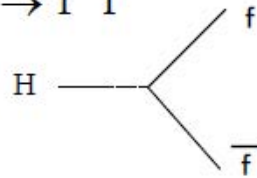
Examples/Higgs $\rightarrow \tau \tau$



Higgs @ 126 GeV?

Examples/Higgs $\rightarrow \tau \tau$

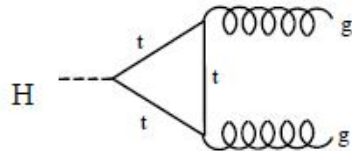
$H \rightarrow f \bar{f}$



$$\Gamma_{f\bar{f}} = C_f \frac{G_F m_f^2 m_H}{4\pi\sqrt{2}} \beta_f^n, \quad \beta_f = \left(1 - \frac{4m_f^2}{m_H^2}\right)^{1/2}$$

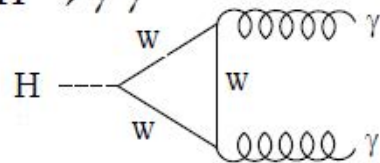
$n = 3$ (scalar), $n = 1$ (pseudoscalar)

$H \rightarrow \text{gluon} + \text{gluon}$



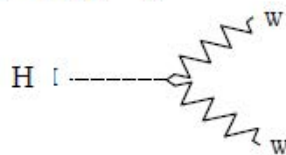
$$\Gamma_{gg} = \frac{G_F m_H^3}{36\pi\sqrt{2}} \left[\frac{\alpha_S(m_H^2)}{\pi} \right]^2 |I|^2, \quad I = 1 \sim 3$$

$H \rightarrow \gamma \gamma$



$$\Gamma_{\gamma\gamma} = \frac{G_F m_H^3}{8\pi\sqrt{2}} \left[\frac{\alpha}{\pi} \right]^2 |I|^2, \quad I \sim -\frac{1}{2}$$

$H \rightarrow W W$

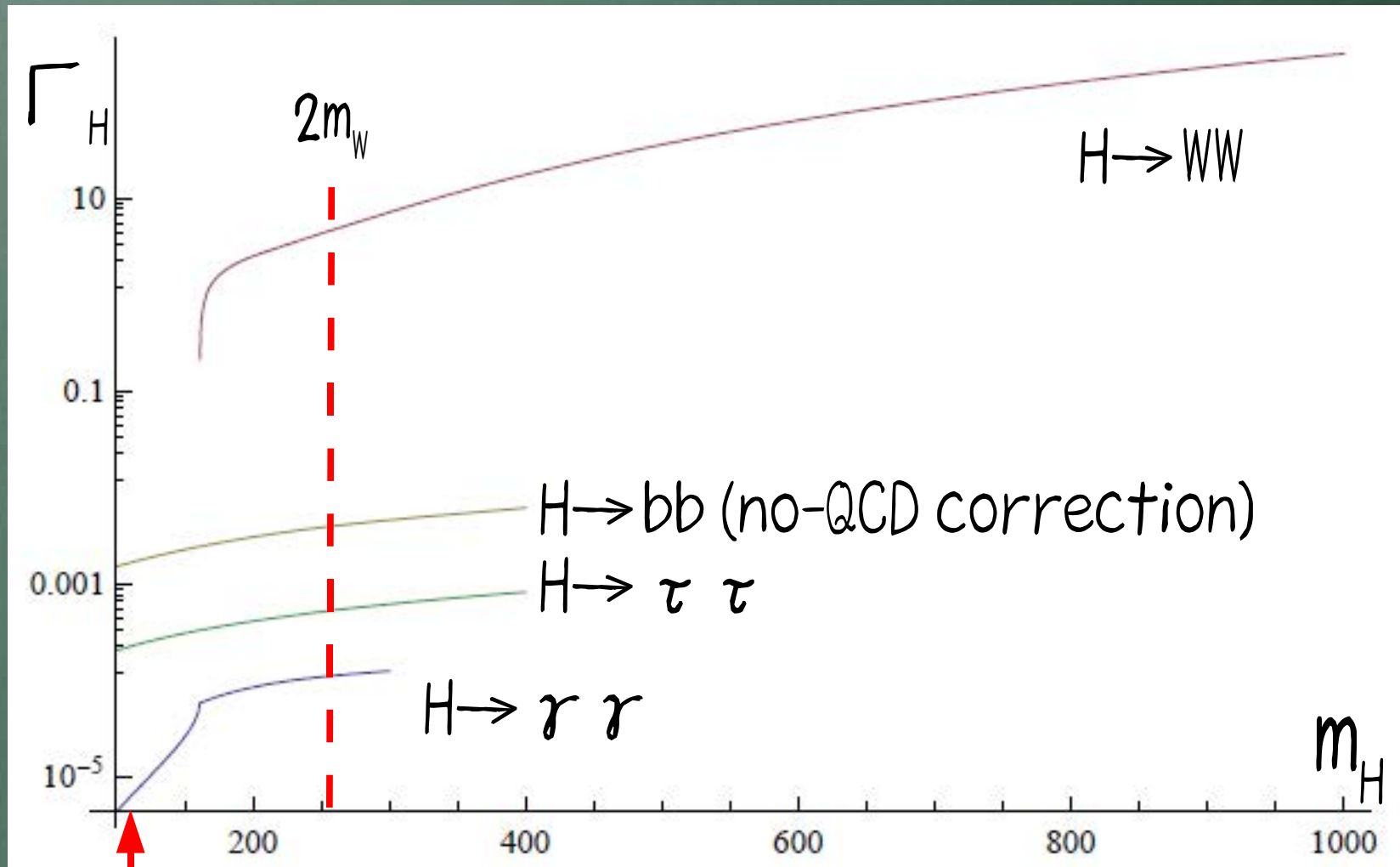


$$\Gamma_{WW} = \frac{G_F m_H^3}{16\pi\sqrt{2}} \delta_V (12x^2 - 4x + 1) \sqrt{1 - 4x}$$

$x = (m_V/m_H)^2, \quad \delta_V = 2 (W) - 1 (Z)$

Examples/Higgs $\rightarrow \tau \tau$

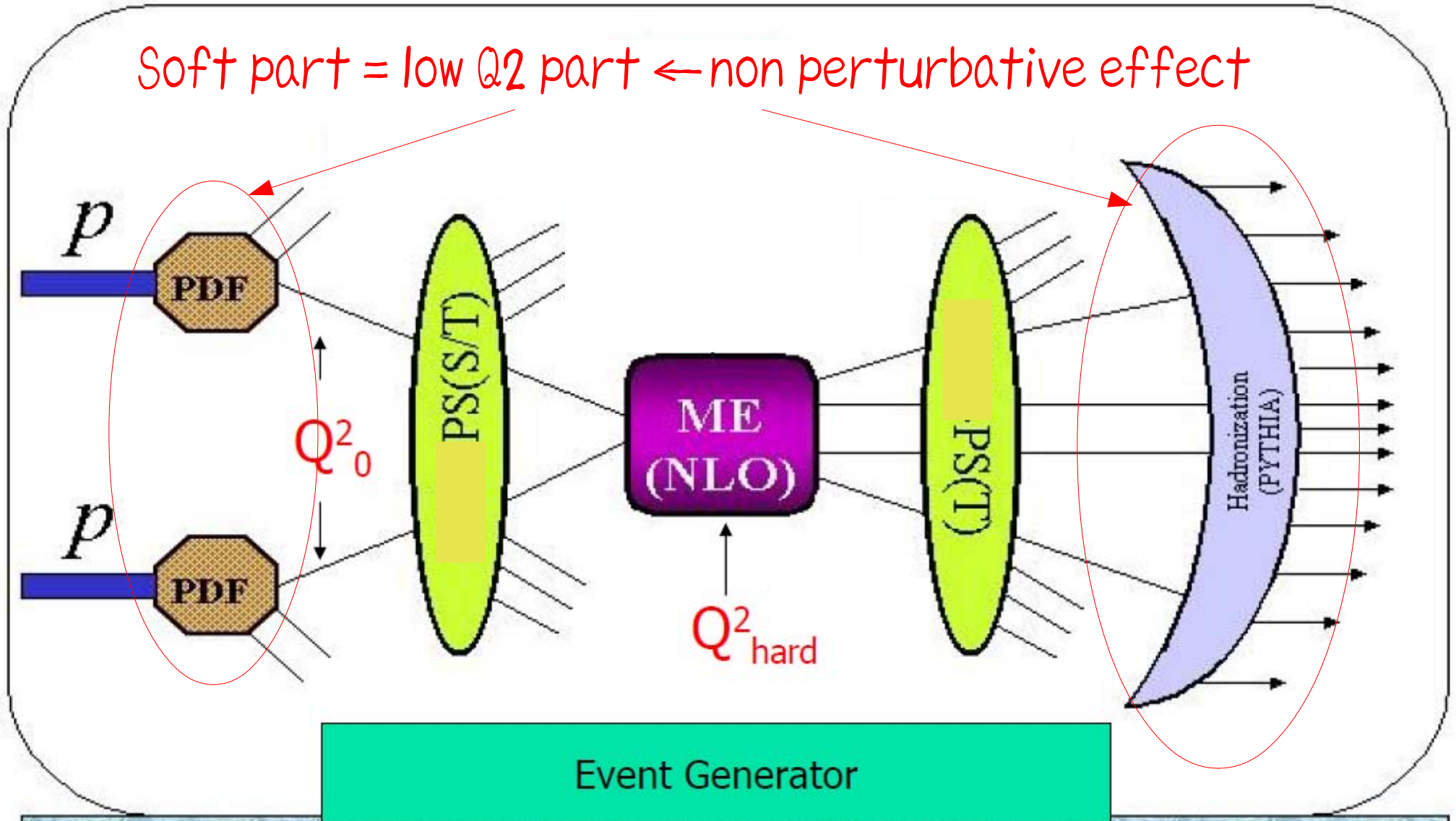
Exercise: Draw this plot.



Higgs @ 126 GeV?

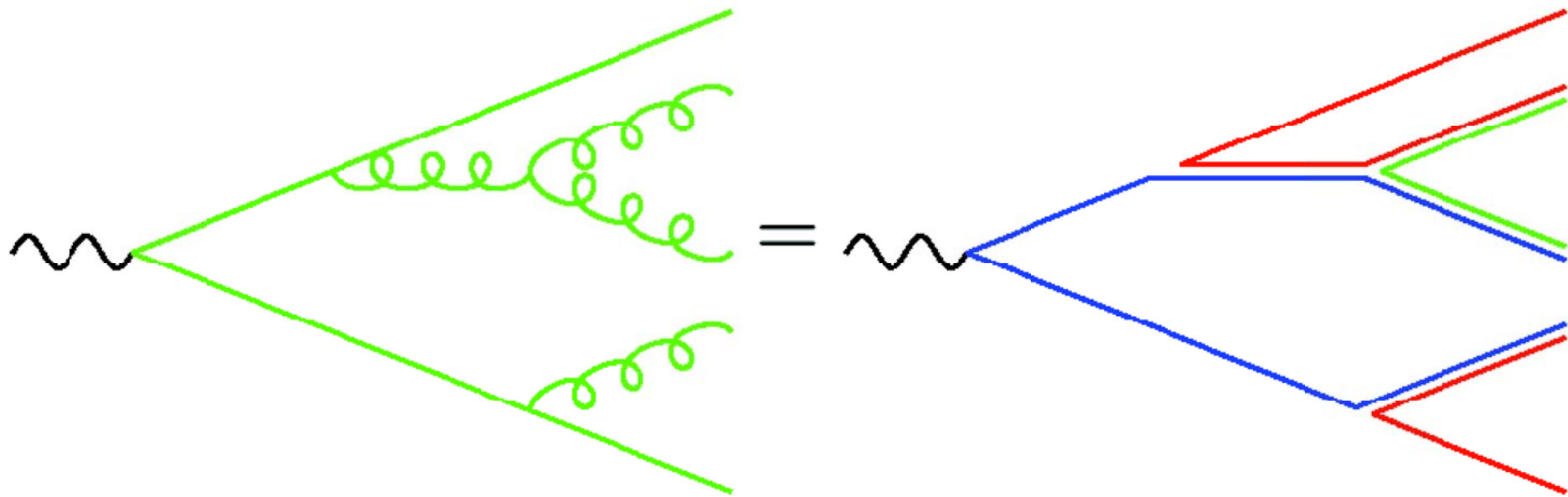
non-perturbative effects

Soft part = low Q^2 part \leftarrow non perturbative effect



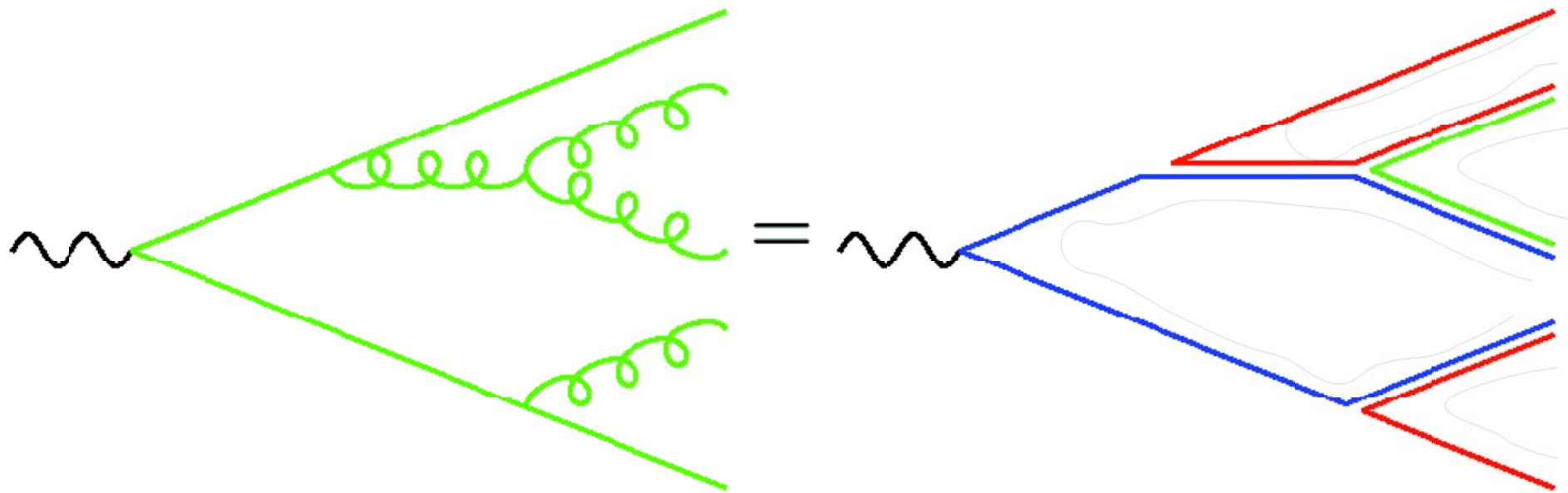
non-perturbative effects/Hadronization String Model

quark, gluon \rightarrow color line



non-perturbative effects/Hadronization String Model

quark, gluon \rightarrow color line



non-perturbative effects/Hadronization String Model

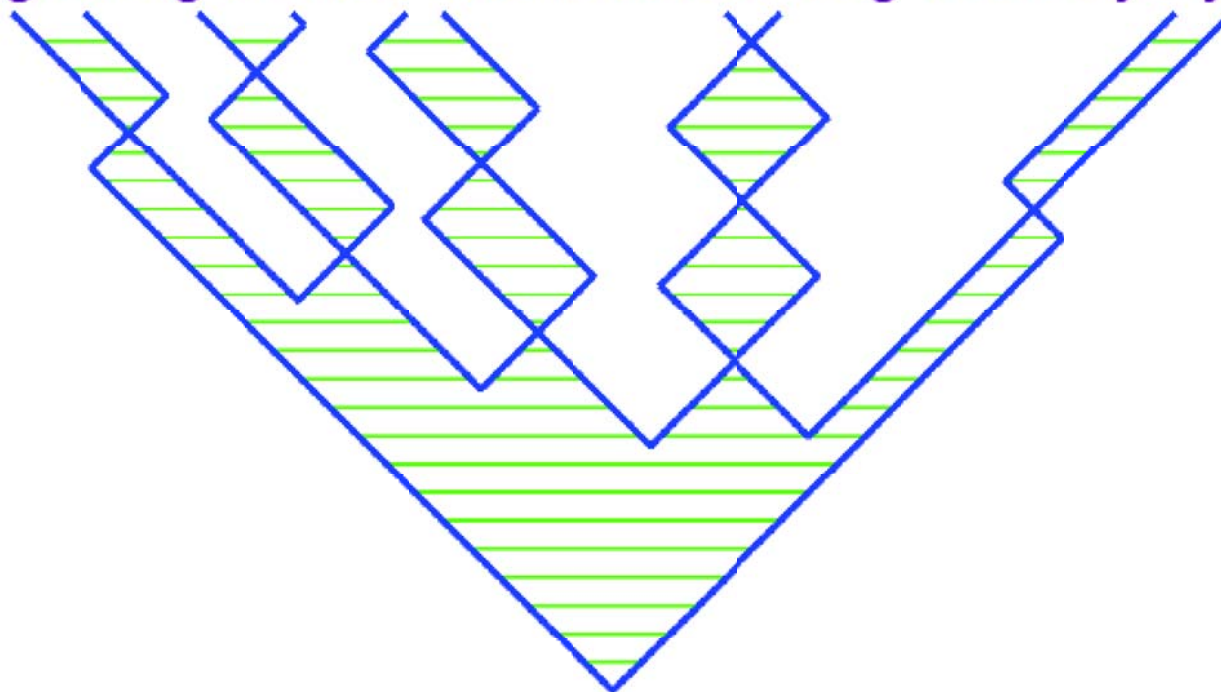
Start by ignoring gluon radiation:

e^+e^- annihilation = pointlike source of $q\bar{q}$ pairs

Intense chromomagnetic field within string \rightarrow $q\bar{q}$ pairs created by tunnelling. Analogy with QED:

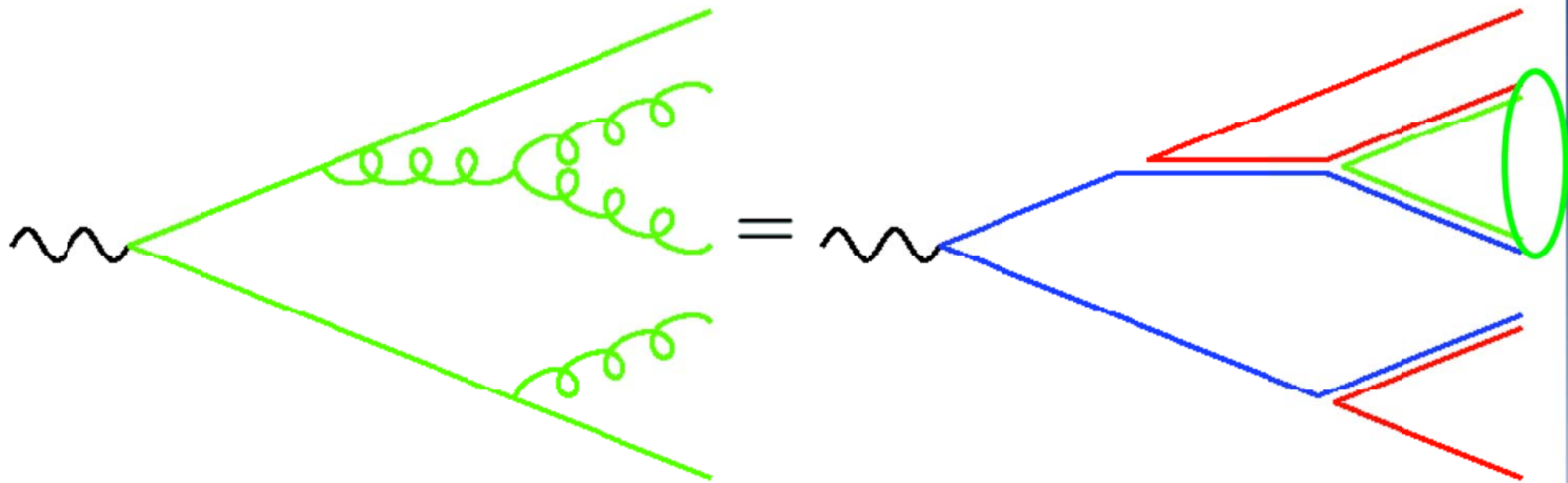
$$\frac{d(\text{Probability})}{dx dt} \propto \exp(-\pi m_q^2 / \kappa)$$

Expanding string breaks into mesons long before yo-yo point.



non-perturbative effects/Hadronization Cluster Model

quark, gluon \rightarrow color line

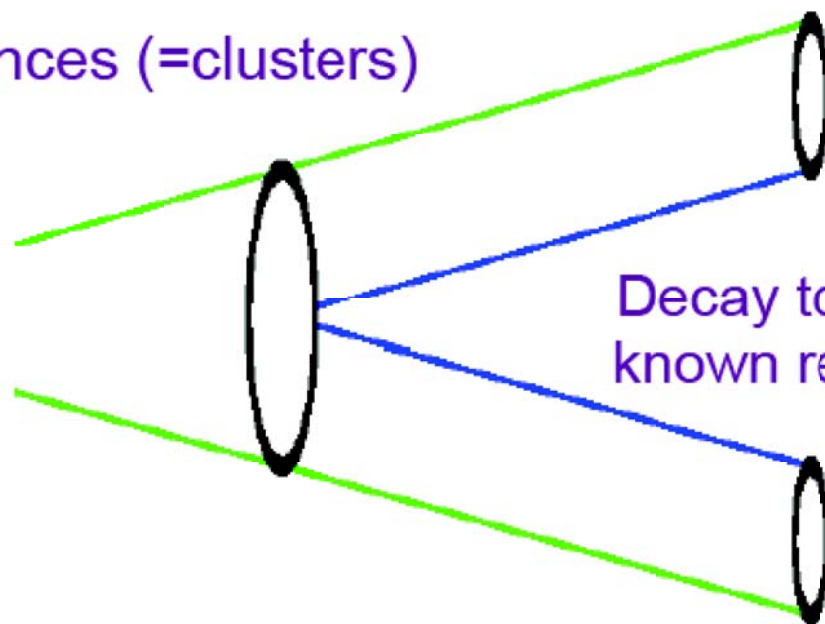


non-perturbative effects/Hadronization Cluster Model

Although cluster mass spectrum peaked at small m , broad tail at high m .

“Small fraction of clusters too heavy for isotropic two-body decay to be a good approximation” \rightarrow Longitudinal cluster fission:

mesonic resonances (=clusters)



Decay to lighter well-known resonances and stable hadrons.

Rather string-like.

Fission threshold becomes crucial parameter.

$\sim 15\%$ of primary clusters get split but $\sim 50\%$ of hadrons come from them.