

# QCD@LHC for beginners

## Lesson 4

Y. Kurihara  
(KEK)  
VSOP-18@Quy Nhon



H. Kawamura (KEK)

# Outline

- Lesson 4

- Around "IR divergence"

- IR divergence in QCD/KLN theorem
- Factorization
- DGLAP Equation/PDF

GOAL: Understanding

(1) IR divergence structure.

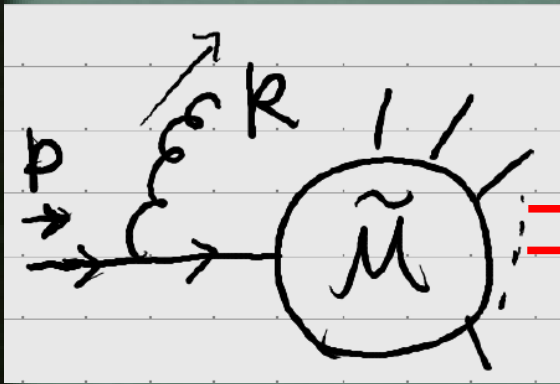
(2) what is PDF.



# IR divergence/Structure

- For QED case:

$k_0 \doteq 0$  Soft or collinear  $\vec{k} \parallel \vec{p}$



$$\tilde{M} \frac{i}{\not{p} - \not{k} - m_e} i g \gamma^\mu u(p)$$

$$\tilde{M} g \frac{\xi p^\mu}{p \cdot k} u(p)$$

$$\left. \begin{aligned} (\not{p} - m_e)u(p) &= 0 \\ \{\gamma^\mu, \gamma^\nu\} &= 2g^{\mu\nu} \end{aligned} \right\}$$

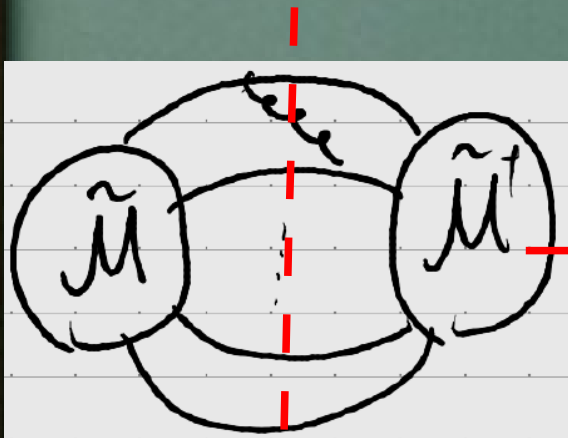
$$p - k \approx \xi p$$

$$p \cdot k = |\vec{k}| (E_p - |\vec{p}| \cos \theta)$$



# IR divergence/Structure

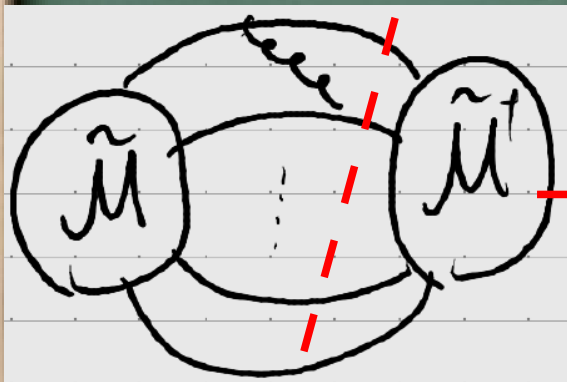
- For QED case:



$$\int d|\vec{k}|^2 d\cos\theta^{(1)} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

$$d=4+2\varepsilon_{\text{IR}}$$

same IR structure



$$-\int \frac{d^3\vec{k}}{(2\pi)^3 2|\vec{k}|} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

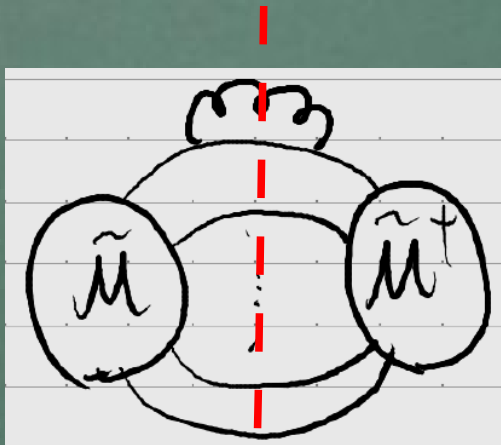
$$1/\varepsilon_{\text{IR}} - \gamma_E + \ln(4\pi), \ln(m_e^2/s)$$

# IR divergence/Structure

- For QED case:

$$\int d|\vec{k}|^2 d\cos\theta^{(l)} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

remove '  
 $p \cdot p = m_e^2$



same IR structure

$$d=4+2\varepsilon_{\text{IR}}$$

$$- \int \frac{d^3\vec{k}}{(2\pi)^3 2|\vec{k}|} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

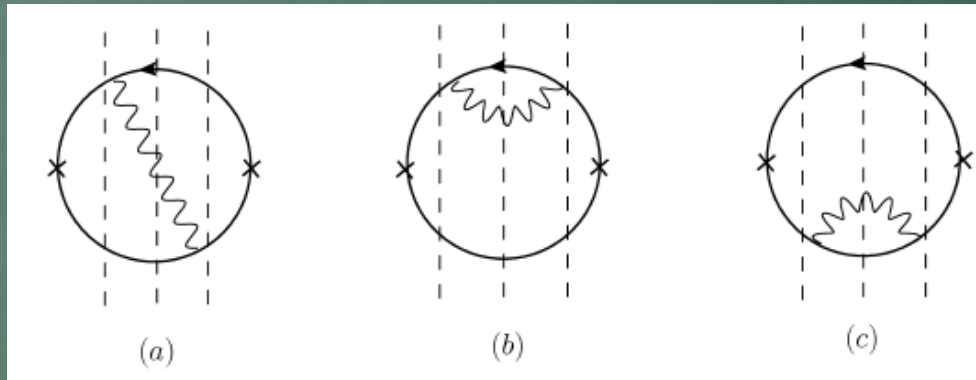
$$1/\varepsilon_{\text{IR}} - \gamma_E + \ln(4\pi), \quad \alpha/(2\pi) \ln(Q^2/m_e^2), \quad \alpha/(2\pi) \ln(Q^2/\mu_F^2)$$

# IR divergence/Structure

- For QED case: Bloch-Nordsieck Theorem

$$d\sigma^{\text{virt.}} + d\sigma^{\text{soft real}} = \text{finite}$$

Summing up all possible graphs and cuts.



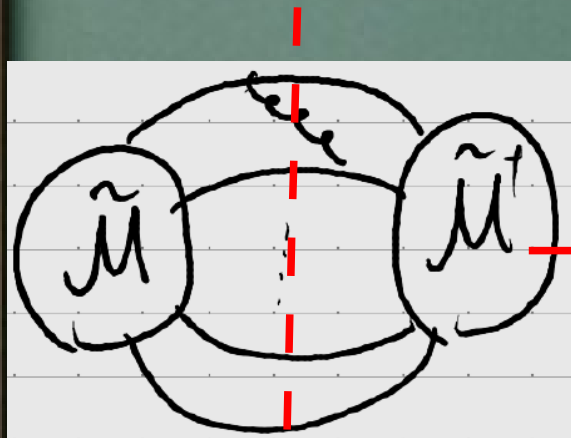
This theorem is proved for all order of QED.



# IR divergence/Structure

- For QCD case: massless quarks

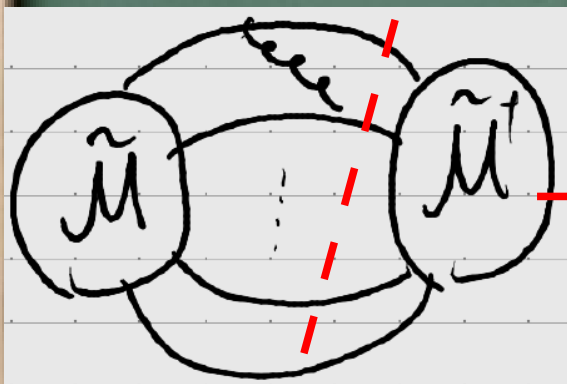
$$E_p = |\vec{p}|$$



$$\int d|\vec{k}|^2 d\cos\theta^{(1)} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

$$d=4+2\varepsilon_{\text{IR}}$$

same IR structure



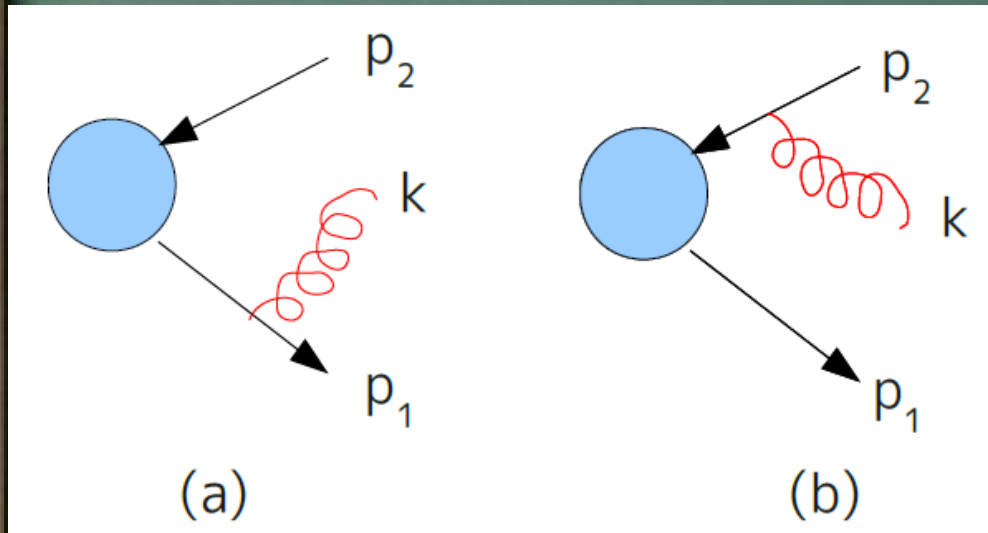
$$-\int \frac{d^3\vec{k}}{(2\pi)^3 2|\vec{k}|} \frac{\xi\xi' p \cdot p'}{|\vec{k}|^2 (E_p - |\vec{p}| \cos\theta)(E_{p'} - |\vec{p}'| \cos\theta')}$$

$$1/\varepsilon_{\text{IR}}^2, 1/\varepsilon_{\text{IR}}, \alpha_s/(4\pi) \ln(Q^2/\mu_F^2)$$

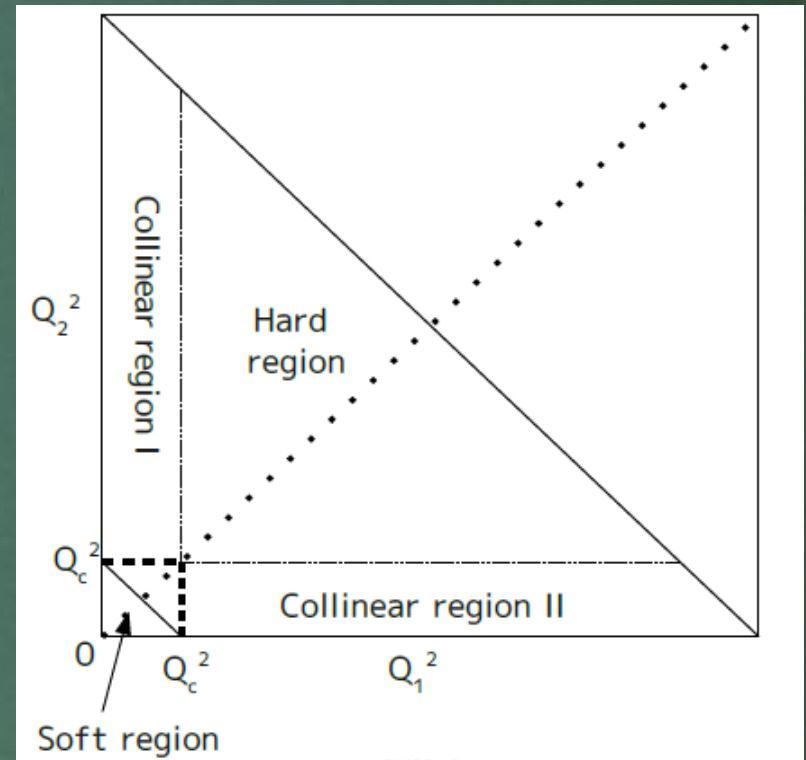
# IR divergence/Structure

- For QCD case: massless quarks

Final state radiation



Daritz Plot



$$Q_2^2 = (p_2 + k)^2$$

$$Q_1^2 = (p_1 + k)^2$$

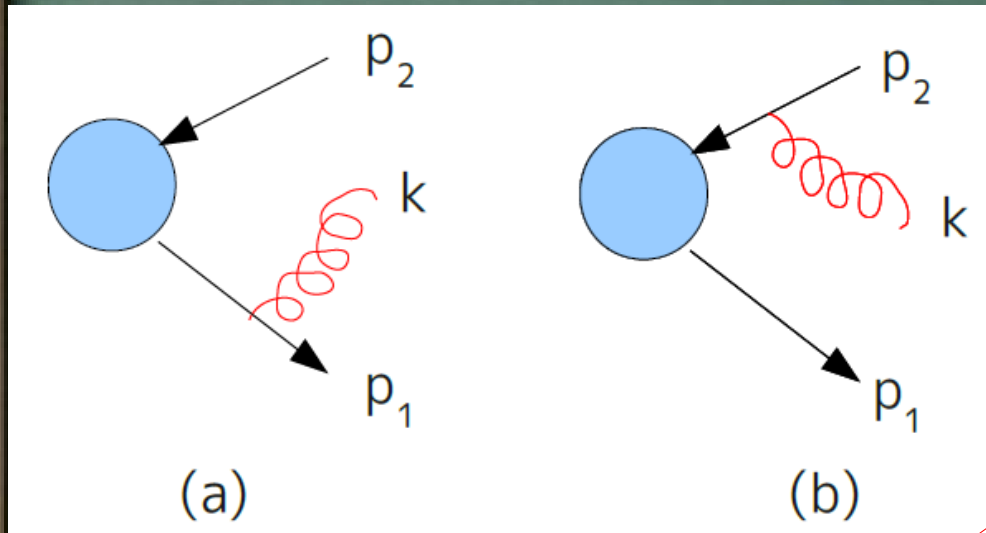


# IR divergence/Structure

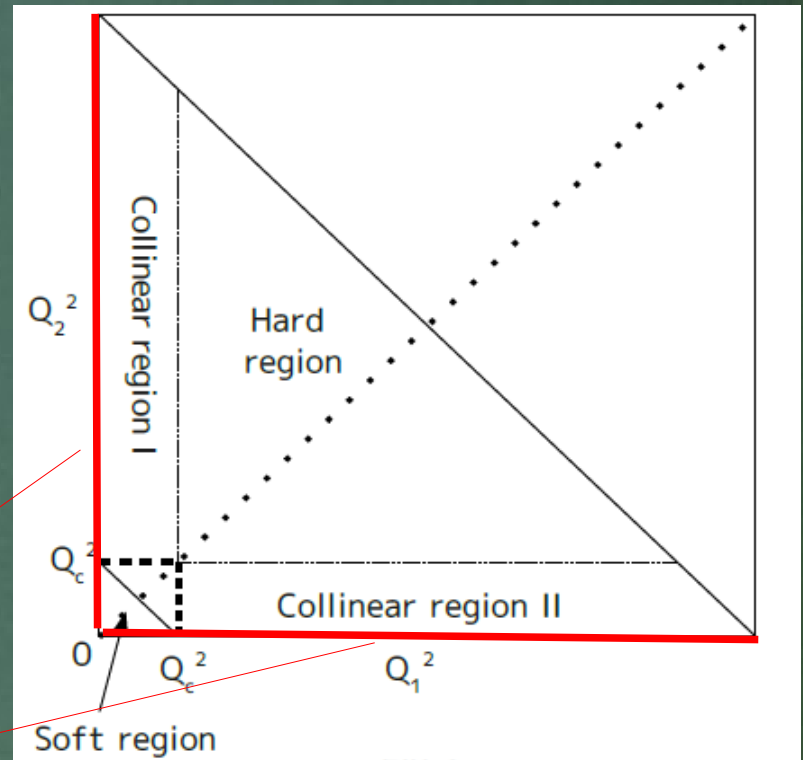
- For QCD case: massless quarks

Final state radiation

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$$Q_2^2 = (p_2 + k)^2$$



$$Q_1^2 = (p_1 + k)^2$$

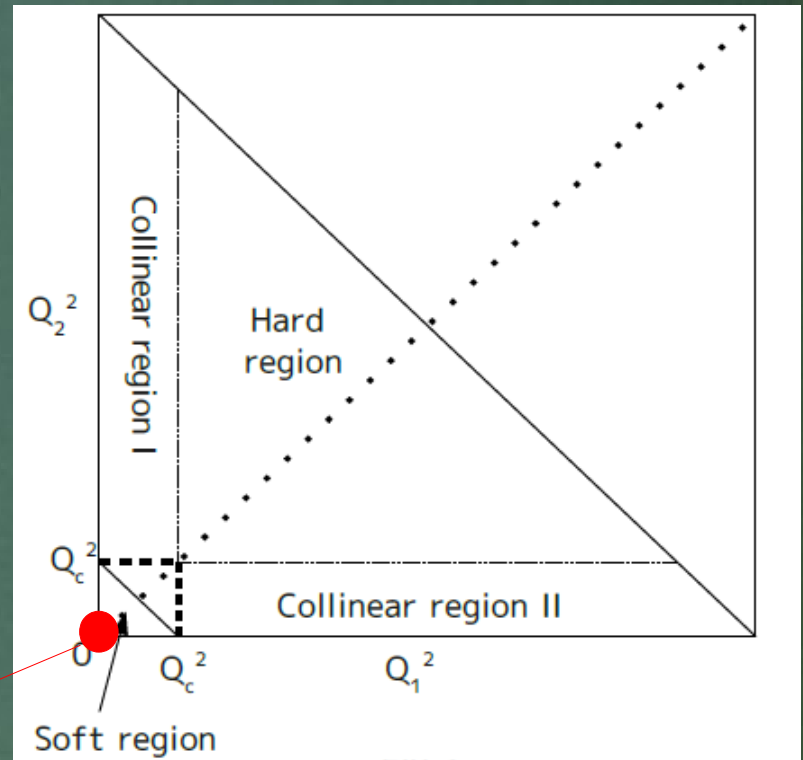
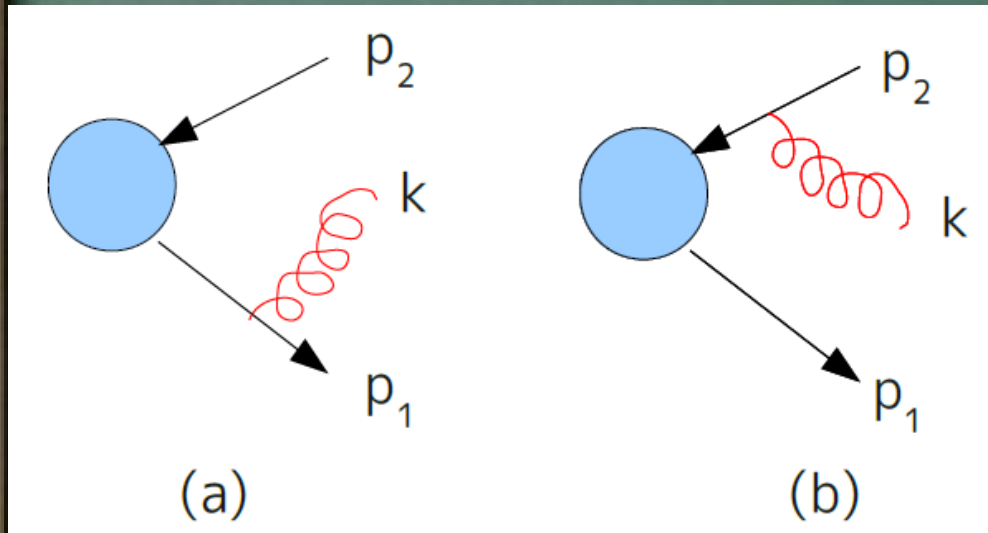
Collinear divergence:  $1/\epsilon_{IR}$

# IR divergence/Structure

- For QCD case: massless quarks

Final state radiation

Daritz Plot



$$Q_2^2 = (p_2 + k)^2$$

$$Q_1^2 = (p_1 + k)^2$$

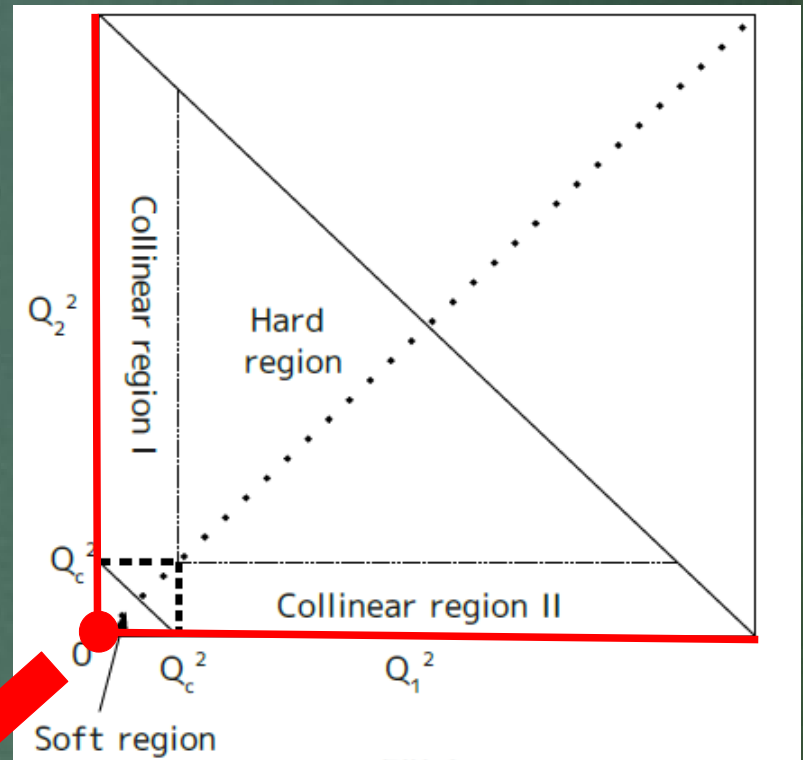
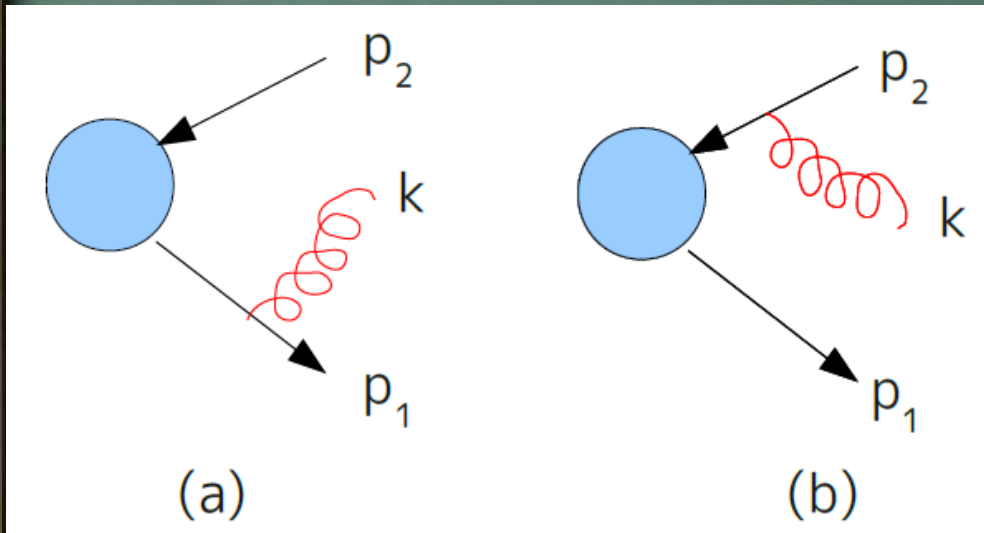
Soft divergence:  $1/\epsilon_{IR}$

# IR divergence/Structure

- For QCD case: massless quarks

Final state radiation

Daritz Plot



$$Q_2^2 = (p_2 + k)^2$$

$$Q_1^2 = (p_1 + k)^2$$

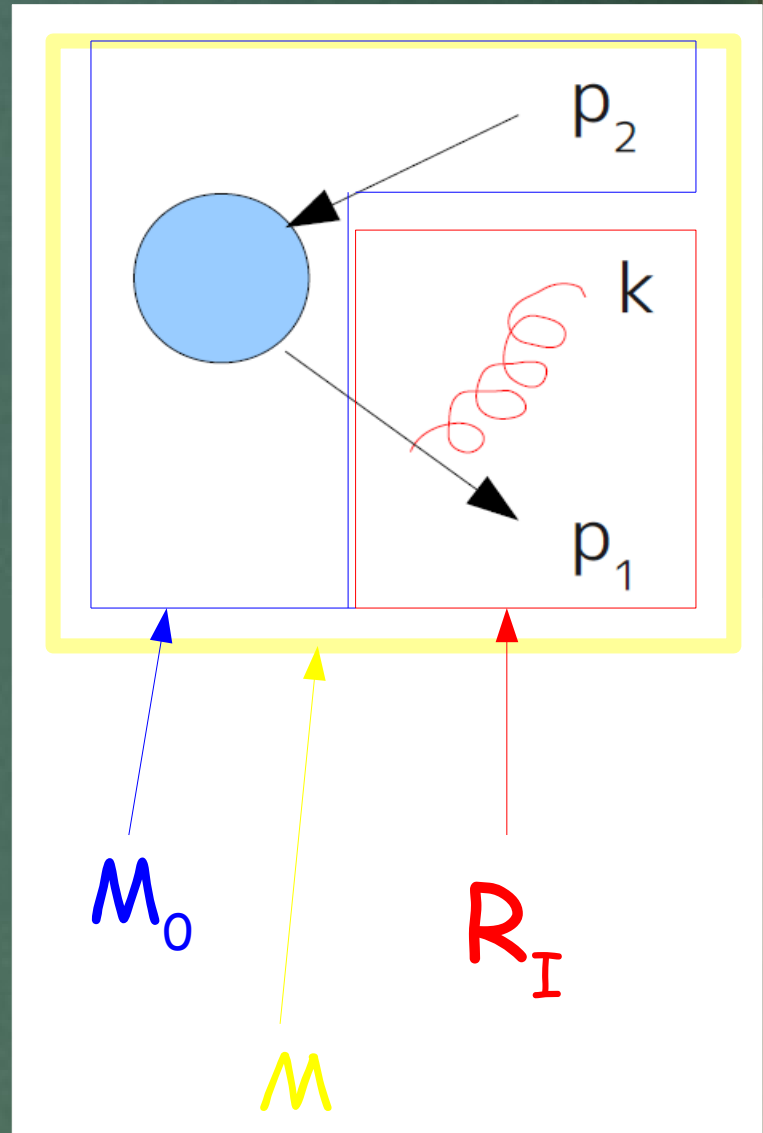
Soft/Collinear divergence:  $1/\epsilon_{IR}^2$

# IR divergence/Example

$$\begin{aligned}
 \sigma_{n+1} &= \frac{1}{(\text{flux})} \int d\Phi_{n+1} |\mathcal{M}|^2 \\
 &\approx \frac{1}{(\text{flux})} \int d\Phi_{n+1} |\mathcal{M}_0|^2 \frac{R_I}{(q_1^2)^2} \\
 &= \frac{1}{(\text{flux})} \int d\Phi_n |\mathcal{M}_0|^2 \int \frac{dq_1^2}{2\pi (q_1^2)^2} \int d\Phi_2 R_I \\
 &= \frac{1}{(\text{flux})} \int d\Phi_n |\mathcal{M}_0|^2 \otimes \delta_{col}, \\
 \delta_{col} &\equiv \int \frac{dq_1^2}{2\pi (q_1^2)^2} \int d\Phi_2 R_I
 \end{aligned}$$

$$R_I = 2g_s^2 \frac{k_T}{x(1-x)} (P(x) + (1-x)\epsilon_{IR})$$

$$P(x) = \frac{1+x^2}{1-x}$$

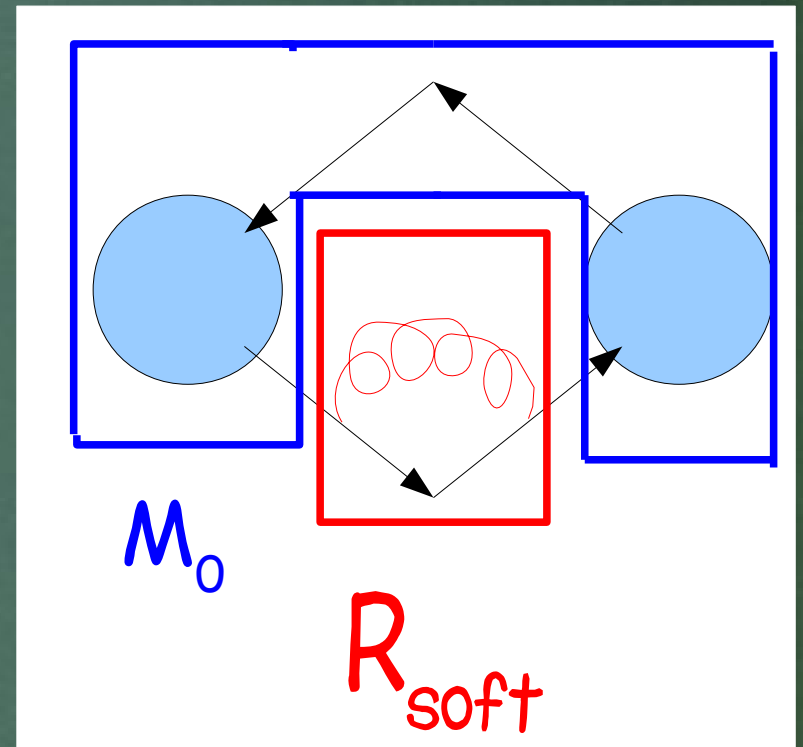


# IR divergence/Example

## Soft Part

$$\sigma_{n+1} = \frac{1}{(\text{flux})} \int d\Phi_{n+1} |\mathcal{M}|^2$$
$$\approx \frac{1}{(\text{flux})} \int d\Phi_n |\mathcal{M}_0|^2 \otimes \delta_{\text{soft}},$$

$$\delta_{\text{soft}} \equiv \int \frac{dk_\theta}{\pi} \int d\Phi_2 \frac{s R_{\text{soft}}}{q_1^2 q_2^2}$$



$$R_{\text{soft}} = -4 p_1 \cdot p_2 (1 + \epsilon_{IR})$$

# IR divergence/Example

Collinear Part:

$$\delta_{col} = \frac{\alpha_s}{\pi} \frac{1}{\epsilon_{IR}^2} + \frac{\alpha_s}{4\pi} \left( -3 + 4 \log(Q_c^2/\mu_F^2) \right) \frac{1}{\epsilon_{IR}}$$

$$- \frac{\alpha_s}{4\pi} \left( -7 + \pi^2 + 3 \log(Q_c^2/\mu_F^2) - \log^2(Q_c^2/\mu_F^2) \right)$$

Soft Part:

$$\delta_{soft} = - \frac{\alpha_s}{\pi} \frac{1}{\epsilon_{IR}^2} - \frac{\alpha_s}{\pi} \left( \log(Q_c^2/s) + \log(Q_c^2/\mu_F^2) \right) \frac{1}{\epsilon_{IR}}$$

$$+ \frac{\alpha_s}{12\pi} \left( \pi^2 - 6 \left( \log(Q_c^2/s) + \log(Q_c^2/\mu^2) \right)^2 \right)$$

# IR divergence/Example

Examples

$$Z \rightarrow d \bar{d}$$

$$\Gamma_{Z \rightarrow d \bar{d}} = \Gamma_0 (1 + \delta_{\alpha_s}),$$

$$\delta_{\alpha_s} = \delta_v + 2\delta_{col} + \delta_{soft}$$

Virtual  
correction:

$$\delta_v = -\frac{\alpha_s}{\pi} \frac{1}{\epsilon_{IR}^2} + \frac{\alpha_s}{2\pi} (3 - 2 \log(s/\mu_F^2)) \frac{1}{\epsilon_{IR}} \\ + \frac{\alpha_s}{2\pi} \left( -8 + \frac{7}{6} \pi^2 + 3 \log(s/\mu_F^2) - \log^2(s/\mu_F^2) \right)$$

Result:

$$\delta_{\alpha_s} = \frac{\alpha_s}{2\pi} \left( -1 + \pi^2/3 - 3 \log(Q_c^2/s) - 2 \log^2(Q_c^2/s) \right)$$

# IR divergence/Example

Examples

$$Z \rightarrow d \bar{d}$$

$$\Gamma_{Z \rightarrow d \bar{d}} = \Gamma_0 (1 + \delta_{\alpha_s}),$$

$$\delta_{\alpha_s} = \delta_v + 2\delta_{col} + \delta_{soft}$$

Virtual  
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$$\delta_v = -\frac{\alpha_s}{\pi} \frac{1}{\epsilon_{IR}^2} + \frac{\alpha_s}{2\pi} (3 - 2 \log(s/\mu_F^2)) \frac{1}{\epsilon_{IR}} + \frac{\alpha_s}{2\pi} \left( -8 + \frac{7}{6} \pi^2 + 3 \log(s/\mu_F^2) - \log^2(s/\mu_F^2) \right)$$

Result:

$$\delta_{\alpha_s} = \frac{\alpha_s}{2\pi} \left( -1 + \pi^2/3 - 3 \log(Q_c^2/s) - 2 \log^2(Q_c^2/s) \right)$$

+ Hard Corr.  $\rightarrow$

$$\Gamma_{\alpha_s} = \Gamma_0 \frac{\alpha_s}{\pi} = 0.01445 \text{ GeV}$$



# IR divergence/KLN Theorem

- For QCD case: massless quarks

Kinoshita-Lee-Nauenberg (KLN) Theorem:

After summing up contributions from all possible degenerate states, the s-matrix has no IR-divergence.

If one looks at some specific initial state, some collinear divergence is still remaining.

Parton Distribution Function



# IR divergence/Initial Radiation

Factorization: Coll. Approx.

Matrix Element

$$\left| \mathcal{M}_{N+1}^{(d)} \right|^2 = \left| \mathcal{M}_N^{(4)} \left( q \rightarrow \sum_{i=1}^N q_i \right) \right|^2 \frac{16\pi}{s\mu^{2\varepsilon_{IR}}} f_c \frac{\alpha_s}{2\pi} P(x) \frac{1}{k_T^2} \left( \frac{1-x}{x} \right)$$
$$q^\mu = p_1^\mu + p_2^\mu - k^\mu,$$

Phase Space

$$d\Phi_{N+1}^{(d)} = d\Phi_N^{(4)} \left( q \rightarrow \sum_{i=1}^N q_i \right)$$
$$\times \frac{1}{16\pi^2 \Gamma(1 + \varepsilon_{IR})} \left( \frac{k_T^2}{4\pi x^2} \right)^{\varepsilon_{IR}} \frac{1}{1-x} dx dk_T^2$$

# IR divergence/Initial Radiation

Factorization: Coll. Approx.

Splitting functions:

$$P_{qq}^{(0)}(z) = C_F \left[ \frac{1+z^2}{(1-z)_+} + \frac{3}{2} \delta(1-z) \right],$$

$$P_{qg}^{(0)}(z) = T_F [z^2 + (1-z)^2],$$

$$P_{gq}^{(0)}(z) = C_F \frac{1+(1-z)^2}{z},$$

$$P_{qq}^{(0)}(z) = 2C_A \left[ \frac{z}{(1-z)_+} - \frac{1-z}{z} + z(1-z) \right] + 2\pi\beta_0 \delta(1-z),$$

# IR divergence/Initial Radiation

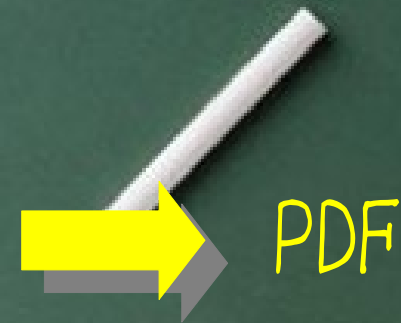
Factorization: Coll. Approx.

Cross Section

$$\begin{aligned}\sigma &= \frac{1}{(2p_1^0)(2p_2^0)v_{rel}} \int_{\Omega_{full}} d\Phi_{N+1}^{(d)} \left| \mathcal{M}_{N+1}^{(d)} \right|^2, \\ &= \left( \frac{s}{4\pi\mu^2} \right)^{\epsilon_{IR}} \frac{B(\epsilon_{IR}, \epsilon_{IR})}{2\Gamma(1 + \epsilon_{IR})} f_c \frac{\alpha_s}{2\pi} \int_0^1 dx \sigma_0(xs) \mathbf{P}(x) \left( \frac{1-x}{x} \right)^{2\epsilon_{IR}}\end{aligned}$$

$1/\epsilon_{IR}$

$$\frac{1}{\epsilon_{IR}} f_c \frac{\alpha_s}{2\pi} \mathbf{P}(x)$$



# IR divergence/Initial Radiation

Factorization: Coll. Approx.

Cross Section

No IR-divergence

$$\sigma_{\text{NLO}} = \left[ \sigma_{\text{tree}} \left( 1 + \overset{\cdot}{\underbrace{\quad}_{\text{v}}} + \overset{\cdot}{\underbrace{\quad}_{\text{s/c}}} \right) + \sigma_{\text{vis}} \right] \otimes \text{PDF/PS}$$

$1/\epsilon_{\text{IR}}^2, 1/\epsilon_{\text{IR}}$   
cancellation

$$\frac{1}{\epsilon_{\text{IR}}} f_c \frac{\alpha_s}{2\pi} P(x)$$

$P(x)$ : Splitting function      Space/time dimension:  $d = 4 + 2\epsilon_{\text{IR}}$

# IR divergence/DGLAP Equation

DGLAP Equation

Splitting function ← pQCD

$$\frac{dD(x, Q^2)}{d \ln Q^2} = \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} P_+(x/y) D(y, Q^2)$$

$$D(x, Q^2) = \Pi(Q^2, Q_s^2) D(x, Q_s^2) + \frac{\alpha}{2\pi} \int_{Q_s^2}^{Q^2} \frac{dK^2}{K^2} \Pi(Q^2, K^2) \int_x^{1-\epsilon} \frac{dy}{y} P(y) D(x/y, K^2)$$

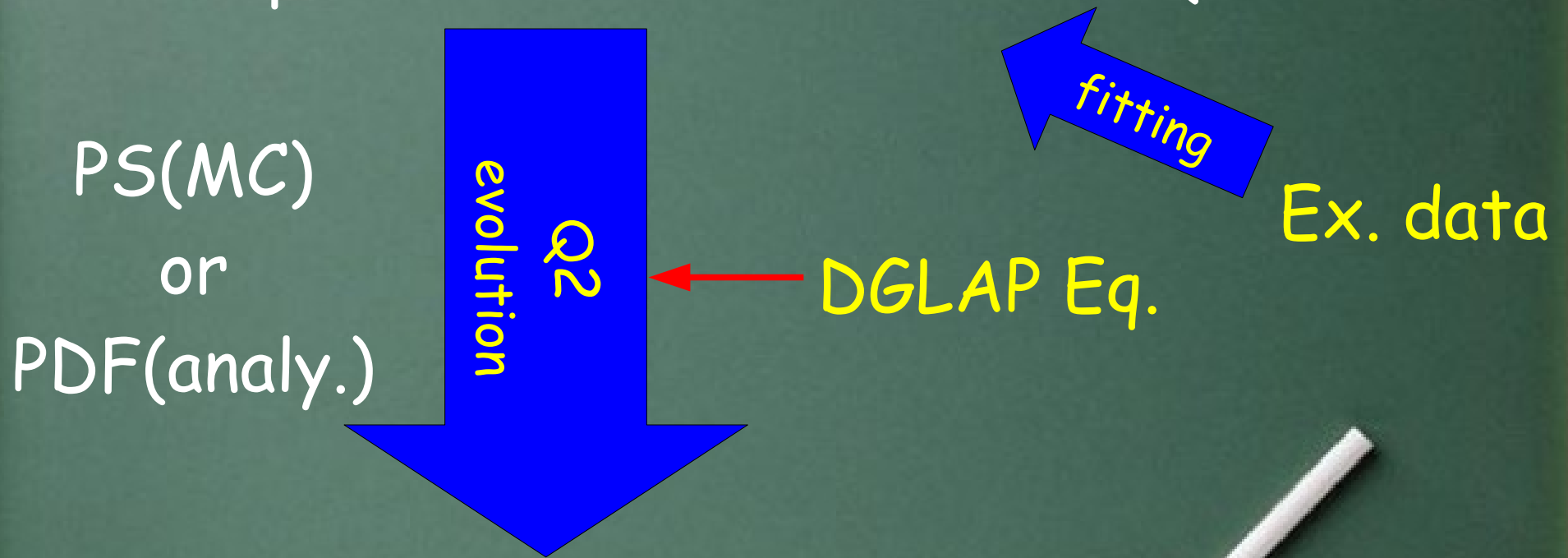
$$\Pi(Q^2, Q_s^2) = \exp\left(-\frac{\alpha}{2\pi} \int_{Q_s^2}^{Q^2} \frac{dK^2}{K^2} \int_0^{1-\epsilon} dx P(x)\right)$$

Sudakov Factor

non-branch provability

# IR divergence/DGLAP Equation

Initial parton distribution @ Low  $Q^2$



Parton distribution @ High  $Q^2$

# IR divergence/DGLAP Equation

## • PDF:

- Initial distribution: data fitting
- $Q^2$  evolution: Analytic solution of DGLAP Eq.
- No kinematical information

## • PS:

- Initial distribution: from PDF
- $Q^2$  evolution: MC method to solve DGLAP Eq.
- generate  $P_t$  distribution

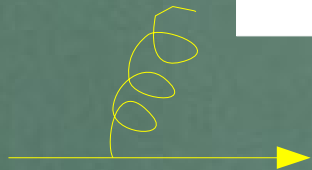


# IR divergence/DGLAP Equation

DGLAP Eq. (Integral-differential equation)

$$\frac{d D(x, Q^2)}{d \log(Q^2)} = \frac{\alpha}{2\pi} \int_0^x \frac{dy}{y} P(x/y) D(y, Q^2)$$

$q \rightarrow q+g :$   $P(x) = \left( \frac{1+x^2}{1-x} \right) \rightarrow \theta(1-\epsilon-x) - \delta(1-x) \int_0^{1-\epsilon} dy P(y)$



$$\frac{d D}{d K^2} K^2 + \frac{\alpha}{2\pi} c(\epsilon) D = \frac{\alpha}{2\pi} P^{\sqrt{5}} D$$

$$c(\epsilon) = \int_0^{1-\epsilon} dz p(z)$$



# IR divergence/DGLAP Equation


$$\frac{d \Pi}{d K^2} K^2 + \frac{\alpha}{2\pi} c \Pi = 0$$


Sudakov form factor  
↓  
non-branching prob.

Inhomogeneous Eq.

$$\frac{d D}{d K^2} K^2 + \frac{\alpha}{2\pi} c D = \frac{\alpha}{2\pi} P^{\sqrt{5}} D$$

$$D(x, K^2) = \hat{D}(x, K^2) \Pi(K^2, Q^2)$$


$$\frac{d \hat{D}}{d K^2} K^2 = \frac{\alpha}{2\pi} P^{\sqrt{5}} \hat{D}$$

$$\hat{D} = \frac{\alpha}{2\pi} \int P^{\sqrt{5}} \hat{D} \frac{d K^2}{K^2} + c$$


# IR divergence/DGLAP Equation

Integral equation

$$D(x, Q^2) = \Pi(Q^2, Q_s^2) D(x, Q_s^2) + \frac{\alpha}{2\pi} P^{\sqrt{5}} D(x, K^2) \frac{dK^2}{K^2}$$

$$\Pi(Q_1^2, Q_0^2) = -\exp\left(\int_{Q_0^2}^{Q_1^2} \frac{\alpha}{2\pi} \frac{dK^2}{K^2} \int_0^{1-\epsilon} dx P(x)\right)$$

Method of successive substitution

Equation  $\phi(x) = f(x) + \lambda \int_a^x K(x, y) \phi(y) dy$

0<sup>th</sup> approx.

$$\phi_0(x) = f(x)$$

1<sup>st</sup>

$$\phi_1(x) = f(x) + \lambda \int_a^x K(x, y) \phi_0(y) dy$$

2<sup>nd</sup>

$$\phi_2(x) = f(x) + \lambda \int_a^x K(x, y) \phi_1(y) dy$$

# IR divergence/DGLAP Equation

Method of successive substitution

No emission

1 parton emission

$$\begin{aligned}\phi_2 = & f(x) + \lambda \int_a^x K(x, y) \phi(y) dy \\ & + \lambda^2 \int_{y_1}^x dy \int_a^{y_1} dy_1 K(x, y_1) \phi(y_1) K(y_1, y) \phi(y)\end{aligned}$$

2 parton emission

# IR divergence/DGLAP Equation

Method of successive substitution

$$\phi(x) = f(x) + \lambda \int_a^x K(x, y) \phi(y) dy$$

$$\begin{aligned} \phi_2 = & f(x) + \lambda \int_a^x K(x, y) \phi(y) dy \\ & + \lambda^2 \int_{y_1}^x dy \int_a^{y_1} dy_1 K(x, y_1) \phi(y_1) K(y_1, y) \phi(y) \end{aligned}$$

$$\phi_n(x) = f(x) + \sum_{l=1}^n \lambda^l \int_a^x K_l(x, y) f(y) dy$$

$$K_l(x, y) = \int \cdots \int dy_1 \cdots dy_{l-1} K(x, y_1) \cdots K(y_{l-1}, y)$$

# IR divergence/DGLAP Equation

The diagram illustrates the DGLAP equation as a sum of diagrams representing gluon emissions. The first row shows a sum of diagrams with one, two, and three gluon emissions, followed by an equals sign and a single gluon emission diagram. The second row shows a plus sign followed by a horizontal line with two gluon emissions above it, a plus sign, and a horizontal line with two gluon emissions above it. The third row shows a plus sign followed by a horizontal line with three gluon emissions above it, a plus sign, and three dots. The fourth row shows a plus sign followed by three dots. The final row shows an equals sign followed by the expression  $\exp(\text{diagram})$  in red, where the diagram is a single gluon emission.

$$\begin{aligned} & \text{[Diagram with 1, 2, 3 gluon emissions]} = \text{[Diagram with 1 gluon emission]} \\ & + \frac{\text{[Diagram with 2 gluon emissions]}}{2} \\ & + \frac{\text{[Diagram with 3 gluon emissions]} + \dots}{3!} + \dots \\ & = \exp(\text{[Diagram with 1 gluon emission]}) \end{aligned}$$