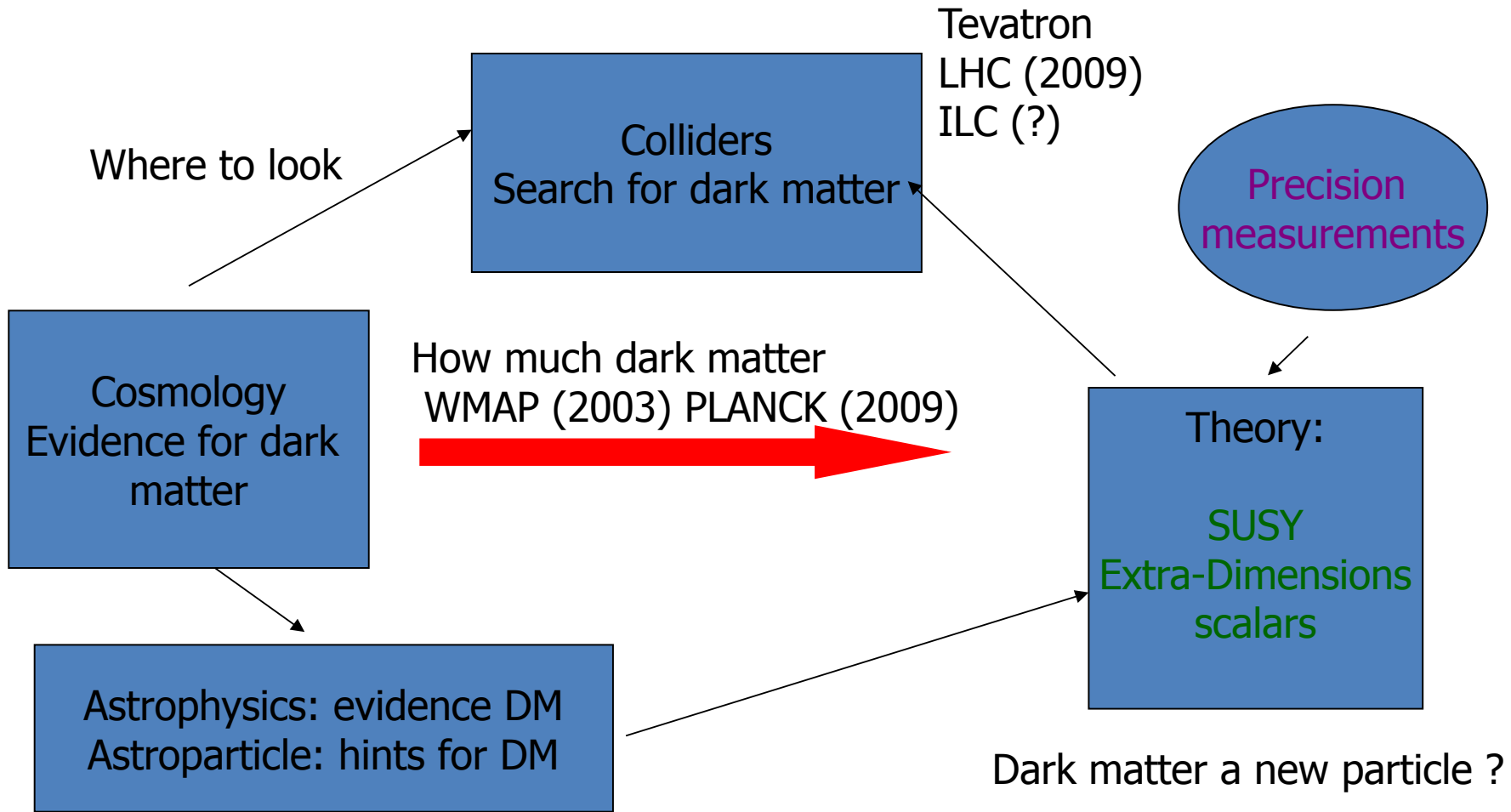


# Searching for Dark matter

- **Direct detection**
  - Elastic scattering of WIMPs off nucleons in a large detector
  - Xenon, CDMS, Dama/Libra
- **Indirect detection**
  - WIMPs annihilation in galaxy, observe decay products
  - e, p,  $\gamma$ : Pamela, Fermi, Hess
  - Neutrinos: IceCube, Km3Net
- **Collider searches**
  - Indirect + Direct searches : Tevatron, LHC, ILC

# Cosmology-(astro)particle-colliders

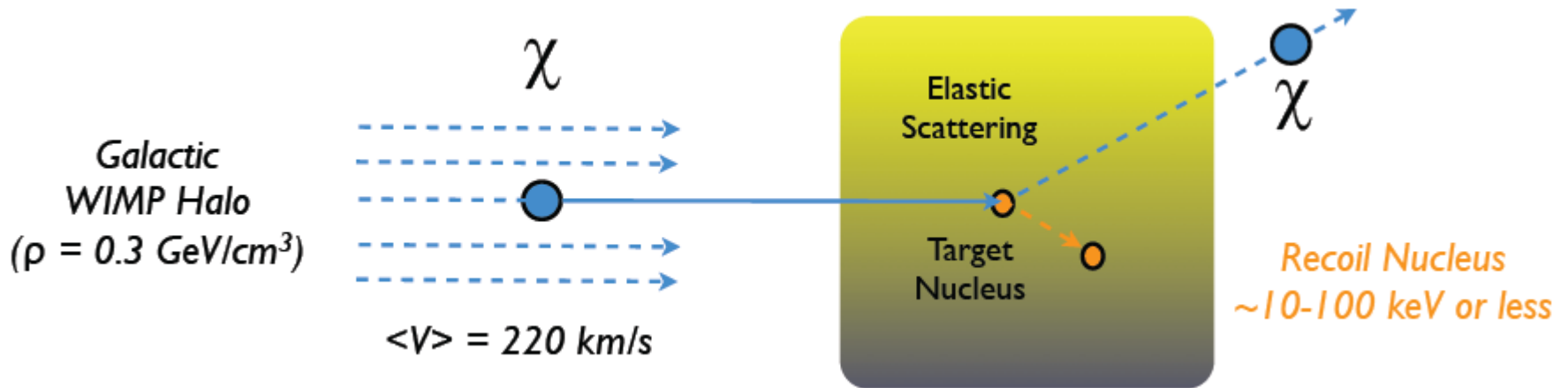


AMS, Egret, PAMELA, Fermi, Hess...  
CDMS, Xenon, DAMA, Kims...

# Direct detection

- Elastic scattering of WIMPs off nuclei in a large detector
- Measure nuclear recoil energy,  $E_R$
- Best way to prove that WIMPs form DM
- Small transfer momentum – typically 100MeV
  - $E_R = q^2/2m_N$   $q$ : transfer momentum
  - $E_R = \mu^2 v^2(1-\cos\theta)/m_N$
  - $\mu = m_\chi m_N / (m_\chi + m_N)$  : reduced mass
  - 100GeV WIMP,  $v=220\text{km/s} \rightarrow E_R < 27\text{keV}$

# Direct detection



- Two types of scattering
  - Coherent scattering on  $A$  nucleons in nucleus, for spin independent interactions
    - Dominant for heavy nuclei
  - Spin dependent int – only one unpaired nucleon
    - Dominant for light nuclei

# Steps to compute nucleus recoil energy

- Wimp-quark/gluon scattering: depend on particle physics model, compute from Feynman diagrams
- Relate WIMP-quark to WIMP-nucleon – quark coefficients in nucleons – determined from first principle + experiments
- WIMP-nucleon  $\rightarrow$  WIMP nucleus : form factor
- Take into account velocity distribution of WIMP
- Recoil energy for WIMP scattering on nucleus
  
- Experimental results are presented in sigma WIMP-proton vs DM mass : easy comparison between exp.

# WIMP- Nucleon amplitude

- For any WIMP, need effective Lagrangian for WIMP-nucleon amplitude *at small momentum*,
- Generic form for a Majorana fermion

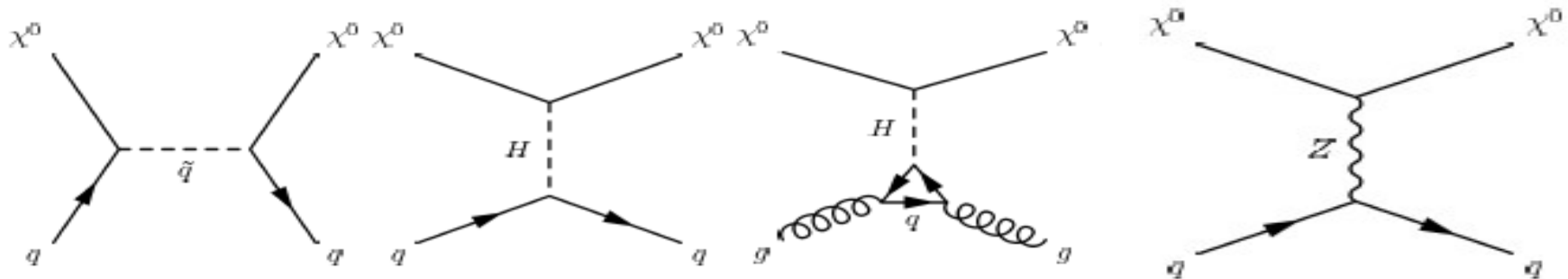
$$\mathcal{L}_F = \lambda_N \bar{\psi}_\chi \psi_\chi \bar{\psi}_N \psi_N + i\kappa_1 \bar{\psi}_\chi \psi_\chi \bar{\psi}_N \gamma_5 \psi_N + i\kappa_2 \bar{\psi}_\chi \gamma_5 \psi_\chi \bar{\psi}_N \psi_N + \kappa_3 \bar{\psi}_\chi \gamma_5 \psi_\chi \bar{\psi}_N \gamma_5 \psi_N$$

$$+ \kappa_4 \bar{\psi}_\chi \gamma_\mu \gamma_5 \psi_\chi \bar{\psi}_N \gamma^\mu \psi_N + \xi_N \bar{\psi}_\chi \gamma_\mu \gamma_5 \psi_\chi \bar{\psi}_N \gamma^\mu \gamma_5 \psi_N$$

- For Majorana fermion only 2 operators survive at small  $q^2$
- First need to compute the WIMP quark amplitudes
  - Computed from Feynman diagrams+ Fierz
  - depends on details of particle physics model
- Effective Lagrangian for WIMP-quark scattering has same generic form as WIMP nucleon

# Direct detection

- Typical diagrams
- Higgs exchange often dominates



For Dirac fermions Z exchange contributes to SI and SD

# Spin independent interactions

- The case of Majorana fermion

$$\mathcal{L}^{SI} = \lambda_N \bar{\psi}_\chi \psi_\chi \bar{\psi}_N \psi_N$$

- Matrix element squared

$$|A_N^{SI}|^2 = 64 (\lambda_N M_\chi M_N)^2$$

- Summing over photons and neutrons

$$|A_A^{SI}|^2 = 64 M_\chi^2 M_A^2 (\lambda_p Z + \lambda_n (A - Z))^2$$

- Cross section for scattering on point like nucleons

$$\sigma_0^{SI} = \frac{4\mu_\chi^2}{\pi} (\lambda_p Z + \lambda_n (A - Z))^2 \quad \mu_\chi = m_{\tilde{\chi}} M_A / (m_{\tilde{\chi}} + M_A).$$



# WIMP-quark to WIMP-nucleon

- Coefficients relate WIMP-quark operators to WIMP nucleon operators
  - Scalar, vector, pseudovector, tensor
  - Extracted from experiments
  - Source of theoretical uncertainties
- Example , scalar coefficients, contribution of q to nucleon mass (heavy quark contribution expressed in terms of gluonic content)

$$\langle N | m_q \bar{\psi}_q \psi_q | N \rangle = f_q^N M_N$$

$$\lambda_{N,p} = \sum_{q=1,6} f_q^N \lambda_{q,p}$$

$$f_Q^N = \frac{2}{27} \left( 1 - \sum_{q \leq 3} f_q^N \right)$$

- Scalar coefficients extracted from ratios of light quark masses, pion-nucleon sigma term and  $\sigma_0$  (size of SU(3) breaking effect)

$$\sigma_{\pi N} = m_l \langle p | \bar{u}u + \bar{d}d | p \rangle$$

$$\sigma_0 = m_l \langle p | \bar{u}u + \bar{d}d - 2\bar{s}s | p \rangle$$

- Large uncertainty in s-quark contribution

$$\sigma_{\pi N} = 55 - 73 \text{ MeV} \quad \text{and} \quad \sigma_0 = 35 \pm 5 \text{ MeV}$$

| Nucleon | $f_{Tu}$ | $f_{Td}$ | $f_{Ts}$ [24] | $f_{Ts}$ [25] | $f_{Ts}$ [20, 26] |
|---------|----------|----------|---------------|---------------|-------------------|
| n       | 0.023    | 0.034    | 0.08          | 0.14          | 0.46              |
| p       | 0.019    | 0.041    | 0.08          | 0.14          | 0.46              |

- 2011: Lattice calculations give new estimates of those coefficients – get s-quark content lower than previously thought ( $\sim 0.02$ )
- The value of the quark coefficient a large impact on the scattering rate - varying coefficients within the range above can in the MSSM lead to almost order of magnitude change in cross section
  - Bottino et al hep-ph/0010203, Ellis et al hep-ph/0502001

# WIMP-nucleon to WIMP-nucleus

- To get rate as a function of the recoil energy must take into account nuclear form factor + velocity distribution
- Ignoring form factor effect expect isotropic scattering in CMS frame - - in lab frame for velocity  $v$  get constant distribution over recoil energy in interval  $0 < E < E_{max}$

$$E_{max}(v) = 2 \left( \frac{v^2 \mu_\chi^2}{M_A} \right)$$

- For fixed  $v$ , recoil energy distribution

$$\frac{d\sigma_A^{SI}}{dE} = \sigma_0^{SI} \frac{\Theta(E_{max}(v) - E) F_A^2(q)}{E_{max}(v)} \quad q = \sqrt{2EM_A}$$

- $F_A(q)$  : form factor (Woods-Saxon form factor)

$$F_A(q) = \int e^{-iqx} \rho_A(x) d^3x \quad \rho_A(r) = \frac{c_{norm}}{1 + \exp((r - R_A)/a)}$$

$$R_A = 1.23A^{\frac{1}{3}} - 0.6 \text{fm}$$

$a=0.52 \text{fm}$  ( extracted from muon scattering data)

- DM have velocity distribution  $f(v)$
- Integrating over incoming velocities  $\rightarrow$  distribution of number of events over the recoil energy

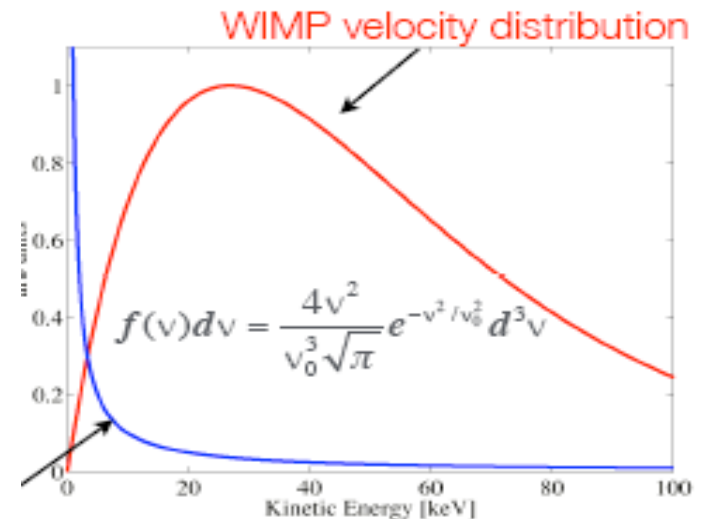
$$\frac{dN^{SI}}{dE} = \frac{2M_{det}t}{\pi} \frac{\rho_0}{M_\chi} F_A^2(q) (\lambda_p Z + \lambda_n (A - Z))^2 I(E)$$

$\rho_0$ : DM density near Earth

- $M_{det}$ : detector mass
- $T$ : exposure time

$$I(E) = \int_{v_{min}(E)}^{\infty} \frac{f(v)}{v} dv$$

$$v_{min}(E) = \left( \frac{EM_A}{2\mu_\chi^2} \right)^{1/2}$$



# WIMP-nucleon to WIMP-nucleus

- Rates (SI and SD) depends on nuclear form factors and velocity distribution of WIMPs + local density

$$\frac{dN^{SI}}{dE} = \frac{2M_{det}t}{\pi} \frac{\rho_0}{M_\chi} F_A^2(q) (\lambda_p Z + \lambda_n (A - Z))^2 I(E)$$

Nuclear form factors

Particle physics  
+ quark content in nucleon

DM velocity  
distribution

$$I(E) = \int_{v_{min}(E)}^{\infty} \frac{f(v)}{v} dv$$
$$v_{min}(E) = \left( \frac{EM_A}{2\mu_\chi^2} \right)^{1/2}$$

# Spin dependent

- Effective Lagrangian for Majorana fermion

$$\mathcal{L}^{SD} = \xi_N \bar{\psi}_\chi \gamma_5 \gamma_\mu \psi_\chi \bar{\psi}_N \gamma_5 \gamma^\mu \psi_N$$

$$|A_N^{SD}|^2 = 192(\xi_N S_N M_\chi M_N)^2$$

- Sum spin currents produced by p and n separately
- $\psi_0$  component vanish  $\rightarrow$  3dim vector current proportional to angular momentum

- $S_p = S_n = 1/2$   $\vec{J}_N^A = S_N^A \vec{J}_A / |J_A|$

- Non trivial summation over spins

$$\begin{aligned} & \sum_{s_\chi, s'_\chi} \sum_{s_A, s'_A} \sum_{1 \leq k, l \leq 3} \langle s_\chi | J_\chi^k | s'_\chi \rangle \langle s'_\chi | J_\chi^l | s_\chi \rangle \langle s_A | J_A^k | s'_A \rangle \langle s'_A | J_A^l | s_A \rangle \\ &= \sum_{1 \leq k, l \leq 3} \text{tr}(J_\chi^k J_\chi^l) \text{tr}(J_A^k J_A^l) = (2J_\chi + 1)J_\chi(J_\chi + 1) \cdot (2J_A + 1)J_A(J_A + 1)/3 \end{aligned}$$

- After average over initial polar,  $(2J_\chi + 1)(2J_A + 1)$  cancels out

- WIMP-nucleus amplitude squared

$$|A^{SD}|^2 = 256 \frac{J_A + 1}{J_A} (\xi_p S_p^A + \xi_n S_n^A)^2 M_\chi^2 M_A^2$$

- Cross section at rest for point-like nucleus

$$\sigma_0^{SD} = \frac{\mu_\chi^2}{\pi} \frac{J_A + 1}{J_A} (\xi_p S_p^A + \xi_n S_n^A)^2$$

- $S_N^A$  are obtained from nuclear calculations or from simple nuclear model  $\sim 0.5$  for nuclei with odd nb of p or n  $\sim 0$  for even nuclei

| Nucleus           | $\langle S_p \rangle_{\text{OGM}}$ | $\langle S_n \rangle_{\text{OGM}}$ | $\langle S_p \rangle$ | $\langle S_n \rangle$ |
|-------------------|------------------------------------|------------------------------------|-----------------------|-----------------------|
| $^{19}\text{F}$   | 0.46                               | 0.0                                | 0.415                 | -0.047                |
|                   |                                    |                                    | 0.368                 | -0.001                |
| $^{27}\text{Al}$  | 0.25                               | 0.0                                | -0.343                | 0.030                 |
| $^{29}\text{Si}$  | 0.0                                | 0.15                               | -0.002                | 0.13                  |
| $^{73}\text{Ge}$  | 0.0                                | 0.23                               | 0.011                 | 0.491                 |
|                   |                                    |                                    | 0.030                 | 0.378                 |
| $^{93}\text{Nb}$  | 0.36                               | 0.0                                | 0.46                  | 0.08                  |
| $^{131}\text{Xe}$ | 0.0                                | -0.166                             | -0.041                | -0.236                |

# Axial vector quark coefficients

- Axial vector current counts the total spin of quarks and antiquarks in nucleon
- Operators for A-V interactions in nucleon related to those in quarks

$$\xi_{N,s} = \sum_{q=u,d,s} \Delta q^N \xi_{q,s} \quad 2s_\mu \Delta q^N = \langle N | \bar{\psi}_q \gamma_\mu \gamma_5 \psi_q | N \rangle$$

- $\Delta q^N$  extracted from lepton-proton scattering , in particular strange contribution to spin of nucleon (measured first by EMC) much larger than expected in naïve quark model

$$\Delta_u^p = 0.842 \pm 0.012, \quad \Delta_d^p = -0.427 \pm 0.013, \quad \Delta_s^p = -0.085 \pm 0.018$$

$$\Delta_u^n = \Delta_d^p, \quad \Delta_d^n = \Delta_u^p, \quad \Delta_s^n = \Delta_s^p$$



# Dirac fermion

- Fermions

$$\begin{aligned} \mathcal{L}_F &= \lambda_{N,e} \bar{\psi}_\chi \psi_\chi \bar{\psi}_N \psi_N + \lambda_{N,o} \bar{\psi}_\chi \gamma_\mu \psi_\chi \bar{\psi}_N \gamma^\mu \psi_N \\ &+ \xi_{N,e} \bar{\psi}_\chi \gamma_5 \gamma_\mu \psi_\chi \bar{\psi}_N \gamma_5 \gamma^\mu \psi_N - \frac{1}{2} \xi_{N,o} \bar{\psi}_\chi \sigma_{\mu\nu} \psi_\chi \bar{\psi}_N \sigma^{\mu\nu} \psi_N \end{aligned}$$

$$\lambda_N = \frac{\lambda_{N,e} \pm \lambda_{N,o}}{2} \quad \text{and} \quad \xi_N = \frac{\xi_{N,e} \pm \xi_{N,o}}{2}$$

- Vector current  $\bar{\psi}_q \gamma_\mu \psi_q$  is responsible for the difference between  $\chi N$  and  $\bar{\chi} N$  interactions. It counts the number of quarks minus antiquarks in the nucleon (valence quarks)

– no uncertainties.

$$\lambda_{N,p} = \sum_{q=u,d} f_{V_q}^N \lambda_{q,p} \quad f_{V_u}^p = 2, f_{V_d}^p = 1, f_{V_u}^n = 1, f_{V_d}^n = 2$$

# Scalar and vector DM

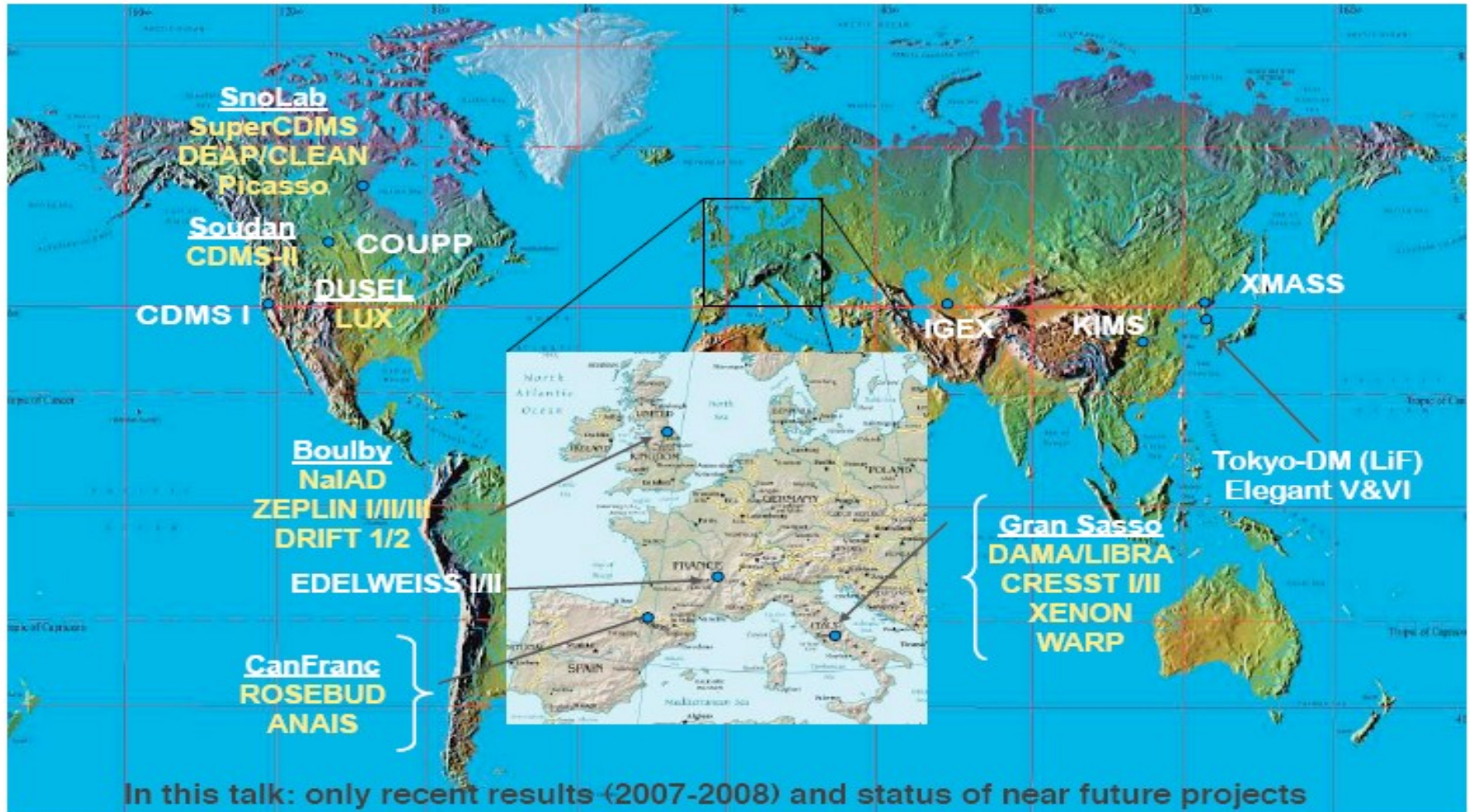
- Complex scalar
  - Only spin independent interactions

$$\mathcal{L}_S = 2\lambda_{N,e}M_\chi\phi_\chi\phi_\chi^*\bar{\psi}_N\psi_N + i\lambda_{N,o}(\partial_\mu\phi_\chi\phi_\chi^* - \phi_\chi\partial_\mu\phi_\chi^*)\bar{\psi}_N\gamma_\mu\psi_N$$

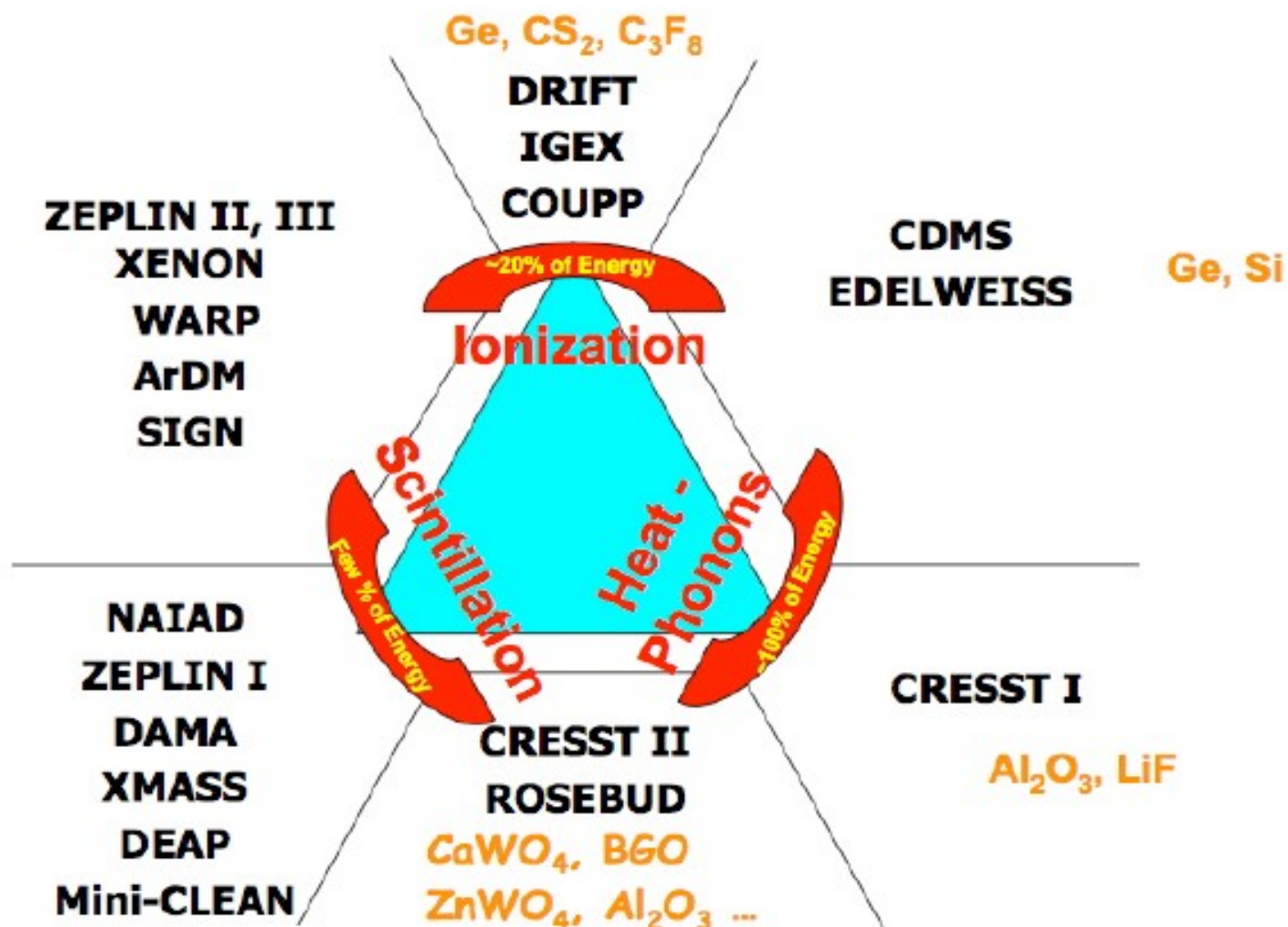
- Complex vector (SI and SD)

$$\begin{aligned}\mathcal{L}_V &= 2\lambda_{N,e}M_\chi A_{\chi,\mu}A_\chi^\mu\bar{\psi}_N\psi_N + \lambda_{N,o}i(A_\chi^{*\alpha}\partial_\mu A_{\chi,\alpha} - A_\chi^\alpha\partial_\mu A_{\chi\alpha}^*)\bar{\psi}_N\gamma_\mu\psi_N \\ &+ \sqrt{6}\xi_{N,e}(\partial_\alpha A_{\chi\beta}^*A_{\chi\gamma} - A_{\chi\beta}^*\partial_\alpha A_{\chi\gamma})\epsilon^{\alpha\beta\gamma\mu}\bar{\psi}_N\gamma_5\gamma_\mu\psi_N \\ &+ i\frac{\sqrt{3}}{g}\xi_{N,o}(A_{\chi\mu}A_{\chi\nu}^* - A_{\chi\mu}^*A_{\chi\nu})\cdot\bar{\psi}_N\sigma_{\mu\nu}\psi_N\end{aligned}$$

# World Wide Wimp searches



# Direct Detection Techniques



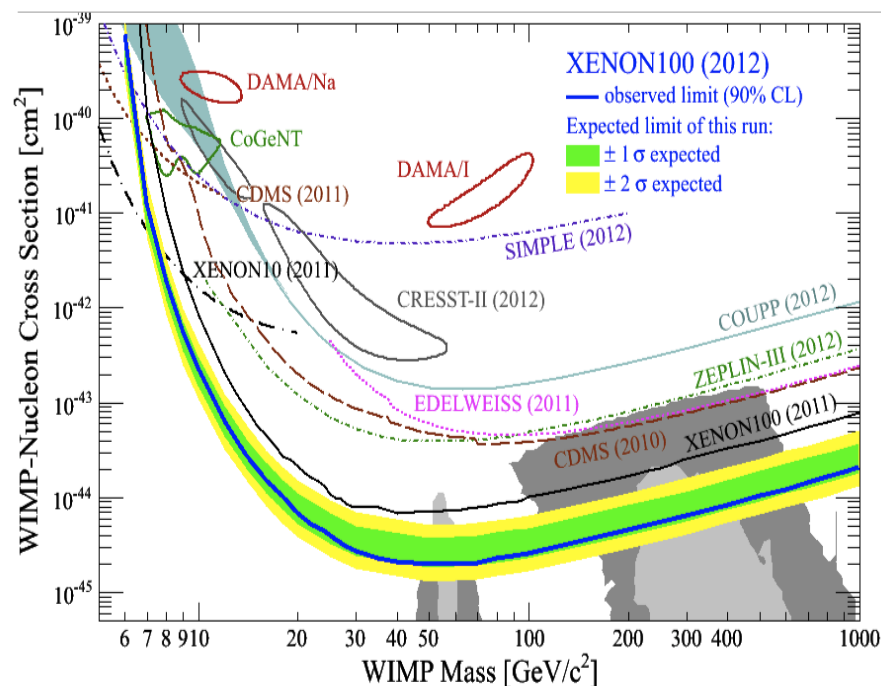
# Direct detection -results

SI

- For easy comparison between experiments – extract  $\sigma_{\chi p}$

$$\sigma_p^{\text{SI}} = \lim_{m_\chi \rightarrow \infty} \sigma \{ m_N = m_p, m_\chi \}$$

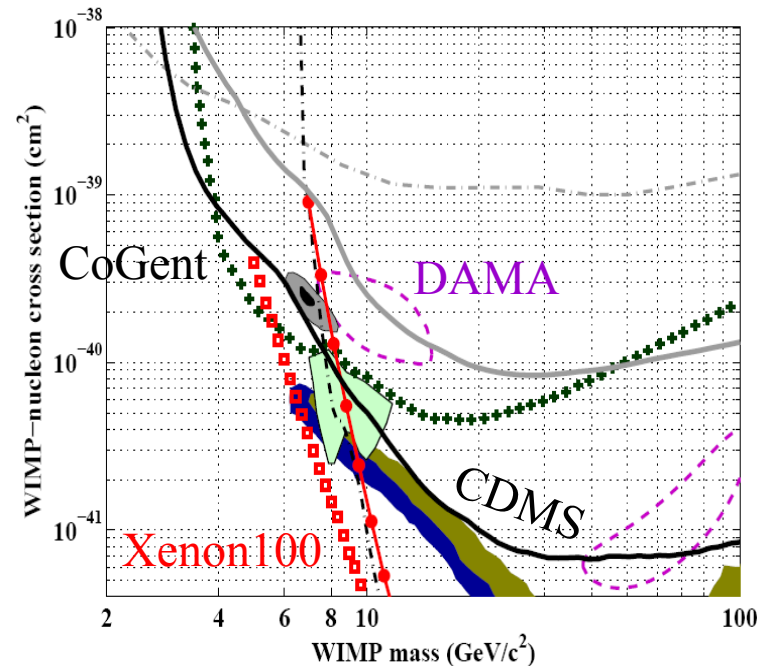
- Assume velocity distribution
- Limits are improving every year
  - Best limits Xenon (2012)
  - DAMA confirm their annual modulation signal



# Direct Detection

- DAMA : signal in annual modulation compatible with light DM ( $8.9\sigma$  )
- Recently CoGent, CDMS , CRESST also reported some signals compatible with 'light' DM
- Some of the favoured regions are excluded by Xenon10, Xenon100, CDMS
  - theoretical uncertainties

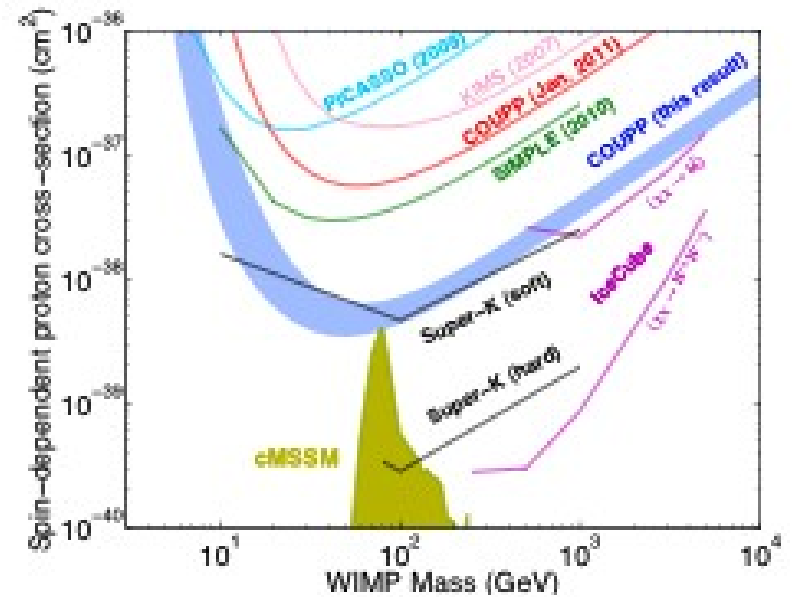
Akerib et al, CDMS, 1010.4290



DM proton scattering  
cross section :  
experimental results

# Spin dependent

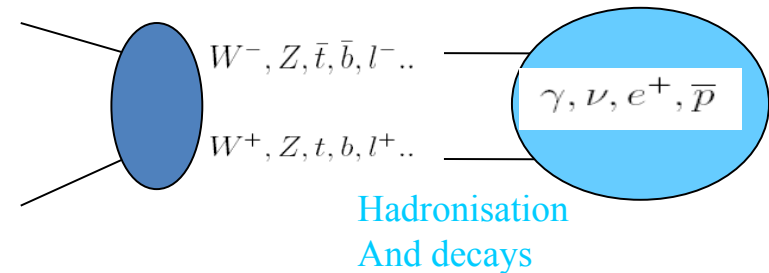
- Also KIMS 1204.2646



Coupp : 1204.3094

# Indirect detection

- Annihilation of pairs of DM particles into SM : decay products observed
- Searches for DM in 4 channels
  - Antiprotons and
  - Positrons from galactic halo/center
  - Photons from galactic halo/center
  - Neutrinos from Sun
- Rate for production of  $e^+, p, \gamma$ 
  - Dependence on the DM distribution ( $\rho$ ) – not well known in center of galaxy
- Typical annihilation cross section at freeze-out



$$Q(x, \mathbf{E}) = \frac{\langle \sigma v \rangle}{2} \left( \frac{\rho(x)}{m_\chi} \right)^2 \frac{dN}{dE}$$

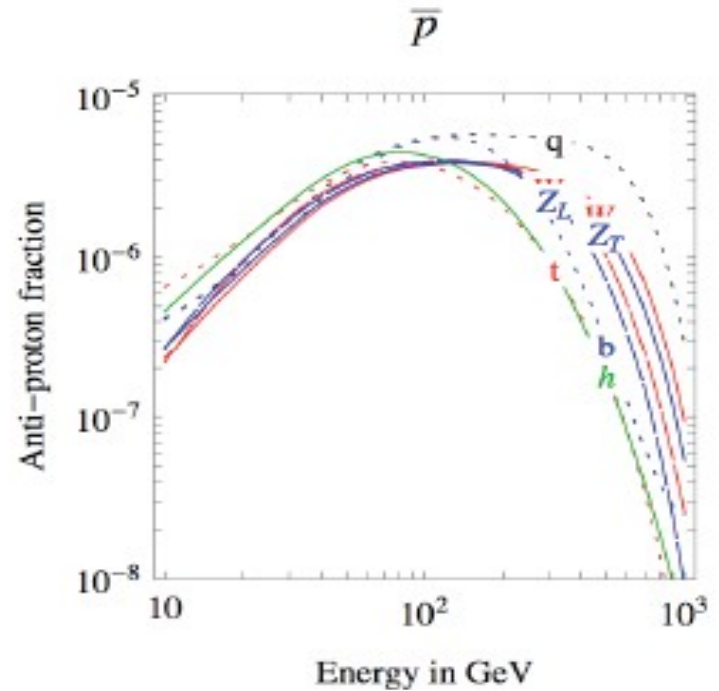
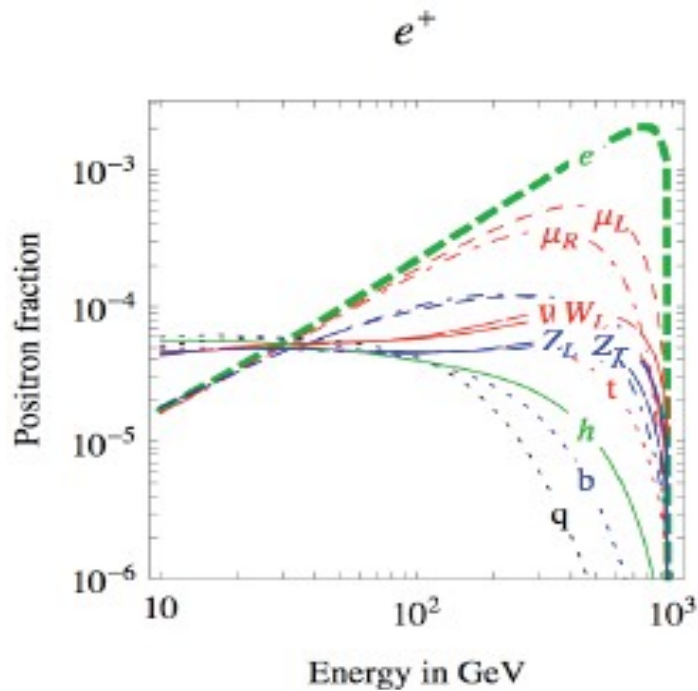
$v=0.001c$

$$\langle \sigma v \rangle = 3 \times 10^{-26} \text{ cm}^3 / \text{sec}$$



# dN/DE

- Spectrum depends
  - mass of DM
  - primary annihilation channels

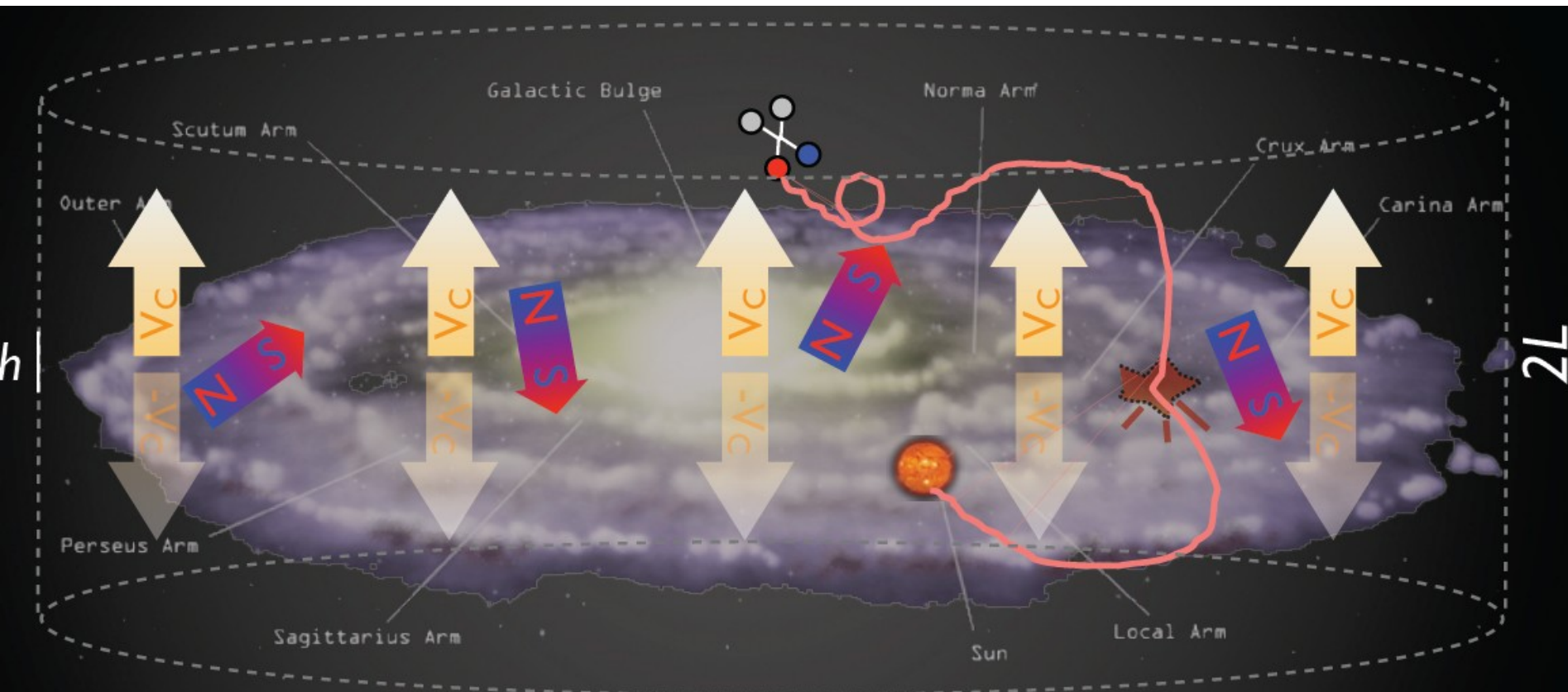


# Propagation of cosmic rays

- *For Charged particle spectrum detected different than spectrum at the source*
- Charged cosmic rays are deflected by irregularities in the galactic magnetic field
  - For strong magnetic turbulence, MC simulations show that effect similar to space diffusion
- Energy losses due to interactions with interstellar medium
- Convection driven by galactic wind
- Reacceleration due to interstellar shock wave

# Antiprotons and positrons from DM annihilation in halo

M. Cirelli, Pascos2009



$$\frac{\partial N}{\partial t} - \nabla \cdot [K(\mathbf{x}, E) \nabla N] - \frac{\partial}{\partial E} [b(E) N] = q(\mathbf{x}, E)$$

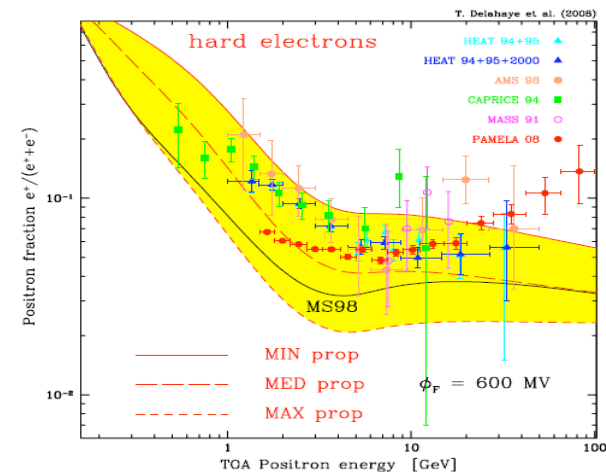
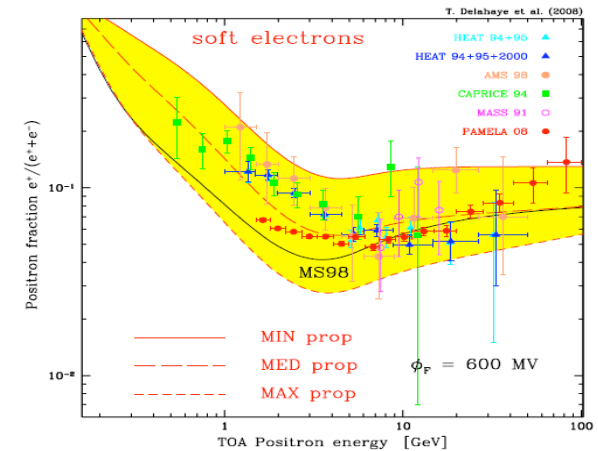
diffusion
Energy losses
Source

# Indirect detection

- For charged particles : solve propagation equation

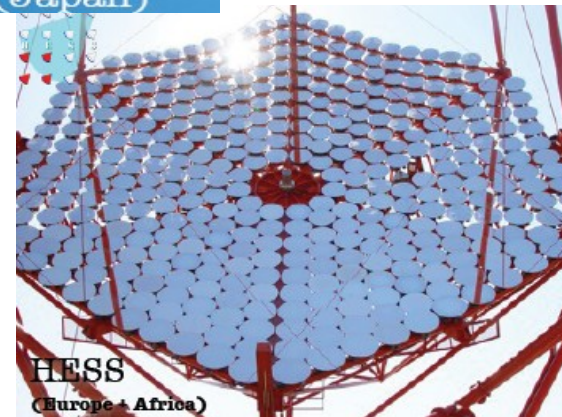
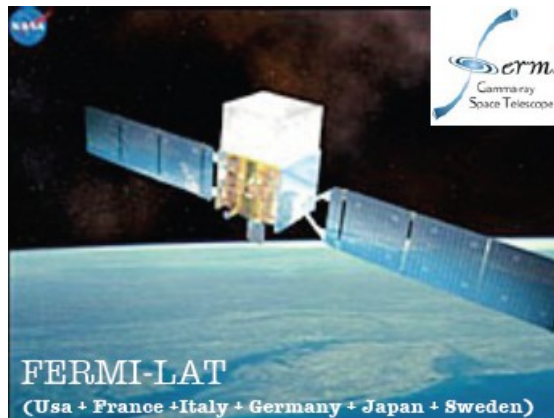
$$\frac{\partial N}{\partial t} - \nabla \cdot [K(\mathbf{x}, E) \nabla N] - \frac{\partial}{\partial E} [b(E) N] = q(\mathbf{x}, E)$$

- Theoretical computation of spectrum of secondary charged particle and from DM annihilation
  - GALPROP – Strong and Moskalenko
  - T. Delahaye, P. Salati et al
- Background spectrum
  - Astro sources: supernova explosions, interaction between cosmic ray nuclei in interstellar medium



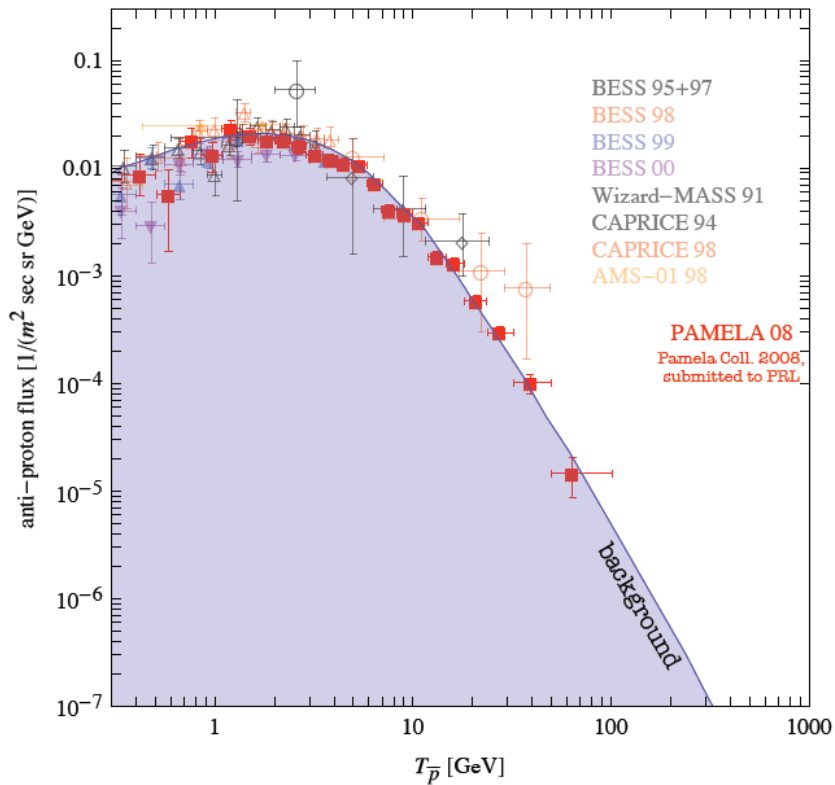
# Indirect DM searches

Payload for  
Anti  
Matter  
Exploration and  
Light nuclei  
Astrophysics



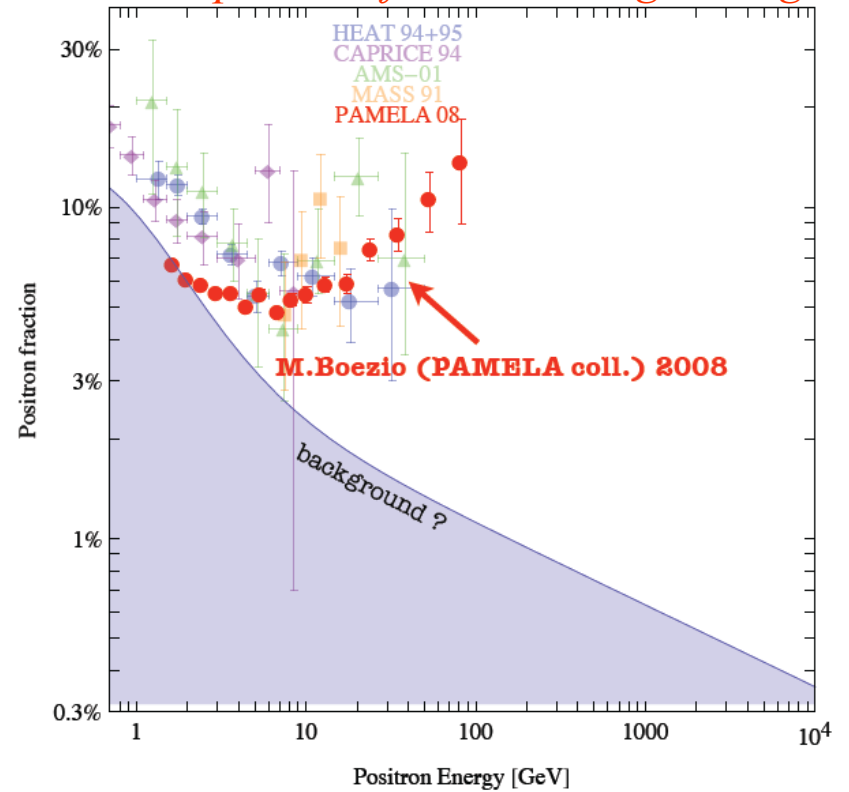
Ground base large array gamma ray telescope

# PAMELA - results



O. Adriani, 0810.4994

*Excess in positron fraction at high energies*

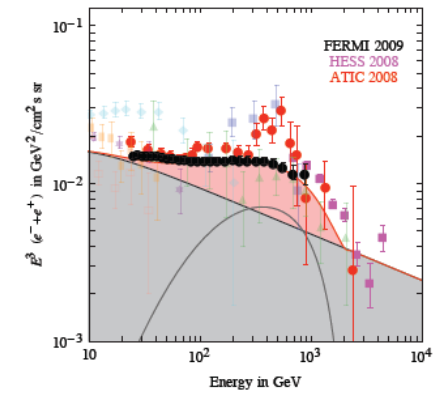
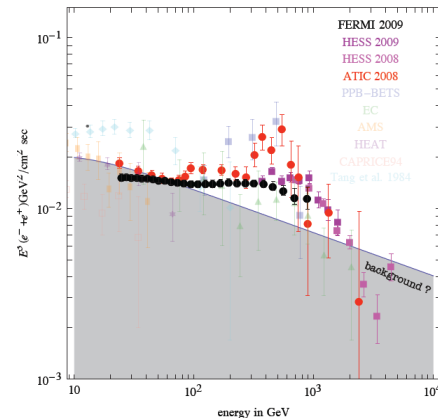


O. Adriani 0810.4995

# DM indirect detection

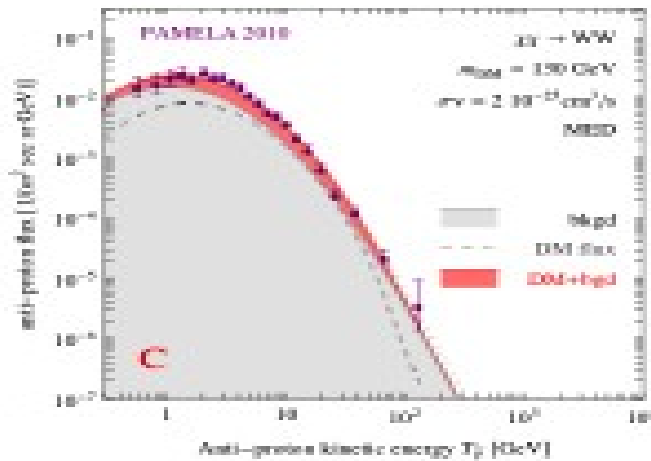
Fermi-LAT 0905.0025

- Results on *total* electron positron spectrum
  - Higher energies than PAMELA
  - Excess over background
- Fit Pamela, Fermi, Hess with e.g. heavy DM (2TeV) annihilating into taus
- Careful investigation of secondary spectrum
- Astro sources (pulsars ) give similar signal

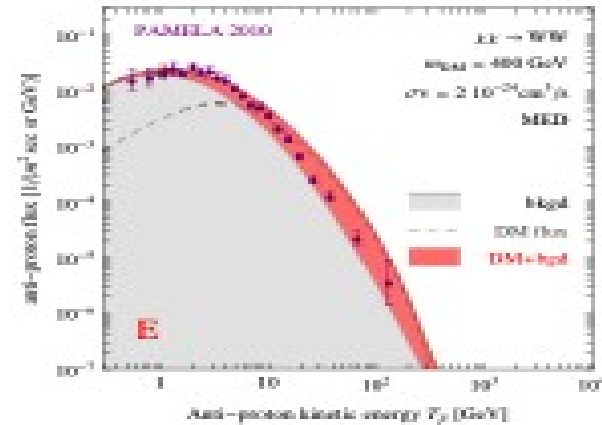


# DM in antiprotons - example

*WW channel*



O. Adriani, 0810.4994



O. Adriani 0810.4995



# Photons

- Flux calculation

$$\Phi_{\gamma,\nu} = \frac{1}{8\pi} \left( \frac{\langle \sigma_{ann} v \rangle}{m_{\chi}^2} \right) \sum_{f.s.} \left( \frac{dN_{\gamma,\nu}}{dE} \right)_{f.s.} \int_{l.o.s.} \rho_s^2$$

- Photon production
  - In decay of SM particles
  - Monochromatic gamma rays ( $\gamma\gamma, \gamma Z$ )
  - Internal bremsstrahlung
- Integral over line of sight depends strongly on the galactic DM distribution

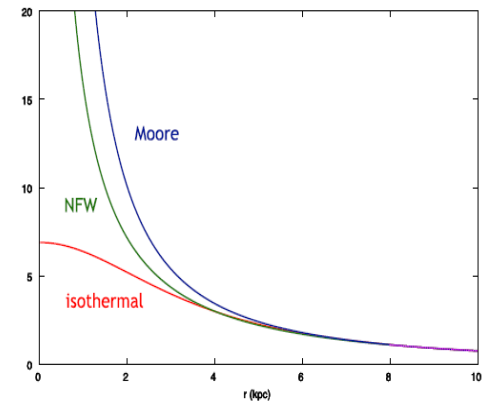
# Dark matter profile

- Dark matter profile

$$\rho_s(r) = \rho_{\odot} \left[ \frac{r_{\odot}}{r} \right]^{\gamma} \left[ \frac{1 + (r_{\odot}/a)^{\alpha}}{1 + (r/a)^{\alpha}} \right]^{\frac{\beta-\gamma}{\alpha}}$$

$$r_{\odot} = 8 \text{ kpc}$$

$$\rho_{\odot} = 0.3 \text{ GeV.cm}^{-3}$$

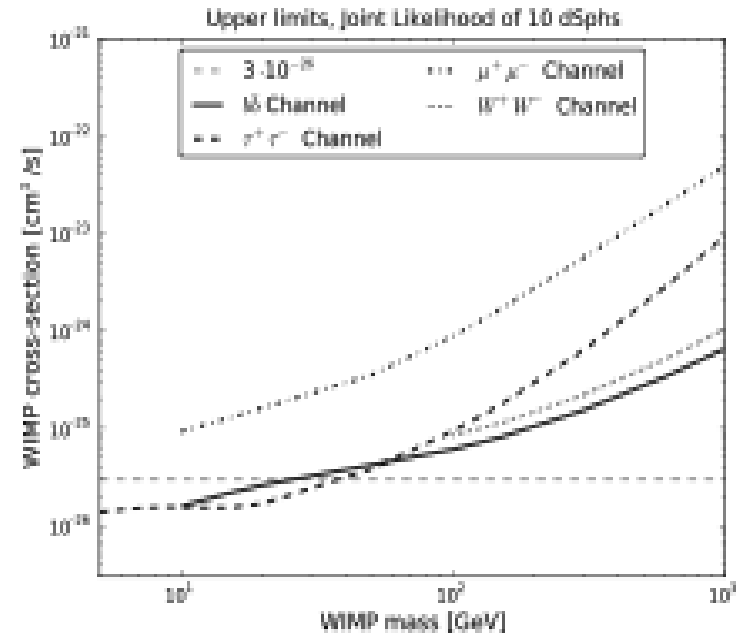


- N-body simulation
- Different halo profile rather similar except in center of galaxy

| Halo model           | $\alpha$ | $\beta$ | $\gamma$ | a (kpc) |
|----------------------|----------|---------|----------|---------|
| Isothermal with core | 2        | 2       | 0        | 4       |
| NFW                  | 1        | 3       | 1        | 20      |
| Moore                | 1.5      | 3       | 1.5      | 28      |

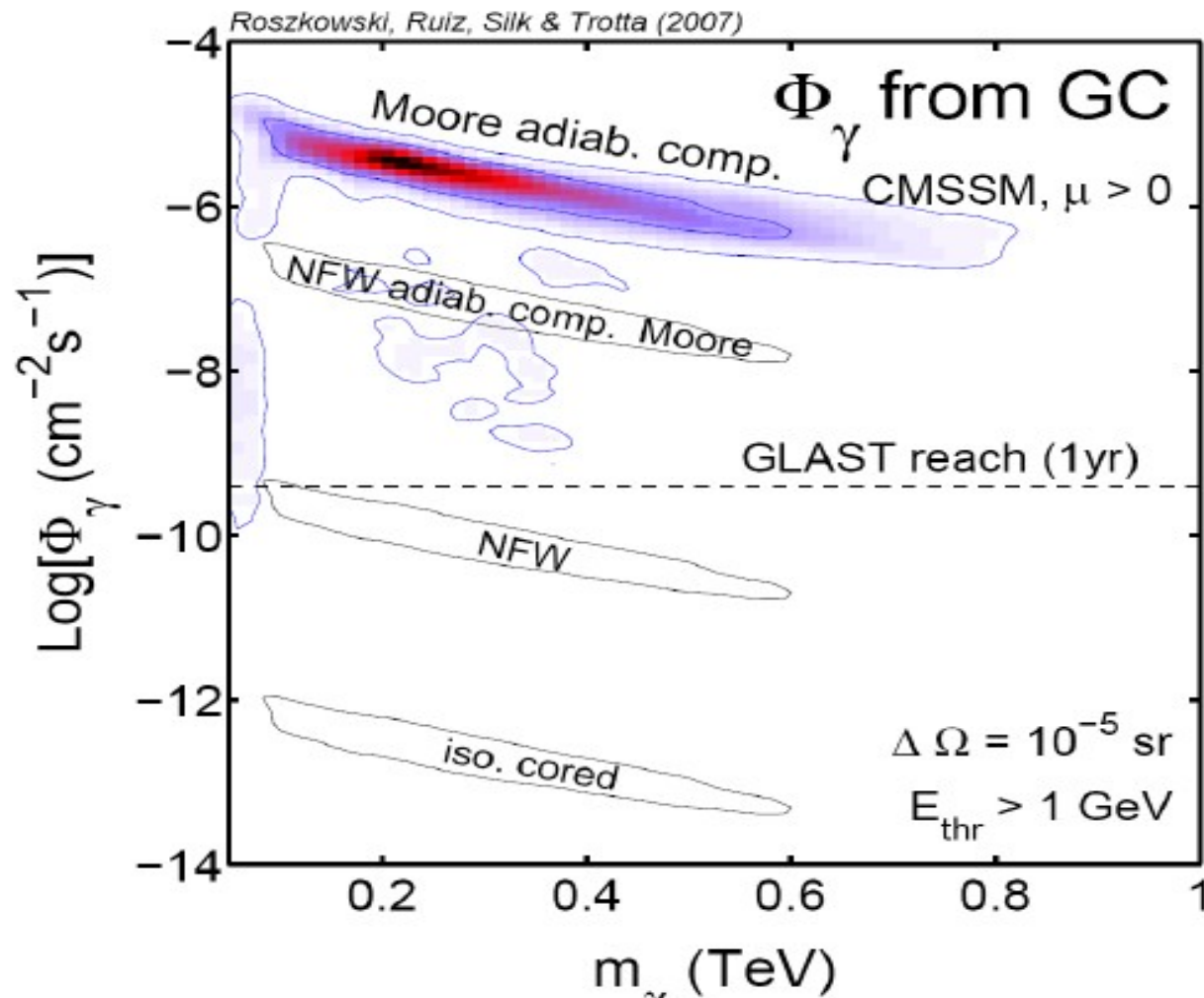
# Photons from dwarf galaxies

- Dwarf galaxies are dominated by DM – good probe
- Not as strong dependence on the density profile (profiles differ strongly only in Galactic center)
- Fermi has derived limits on photon flux and DM cross section for different channels
- Low masses probe the relic density favoured value



# Impact of DM profile on rate

## A SUSY example



# Summary

- A number of direct and indirect detection offer good prospects to probe dark matter (probe  $\sigma v$  and  $M_{\text{DM}}$ )
- Photons and antiprotons sensitive to light DM with expected cross section
- Direct detection can probe both SI and SD interactions in protons and neutrons using different detectors
- Already constrain some favoured models
- Theoretical uncertainties are important
- Hints of signals, not clear it is DM: DAMA, Pamela, Fermi gamma-ray line

- Extra notes
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# Velocity distribution of DM

- Nuclear recoil energy measured depends on WIMP velocity distribution in rest frame of detector → distribution in rest frame of galaxy + Earth velocity in this frame

$$v_0 = 220 \pm 20 \text{ km/s}$$

- Velocity of rotation in LSR
- Peculiar velocity of the Sun
- Earth velocity in Galactic frame:  $\vec{v}_1 = \vec{v}_0 + \vec{v}_{\text{pec}} + \vec{v}_E$  (Earth in solar system)

$$\vec{v}_e = v_e (-\sin(2\pi t), \sin \gamma \cos(2\pi t), \cos \gamma \cos(2\pi t))$$

- Velocity of DM particles on Earth = obtained from velocity of DM particles in Galactic Rest Frame

$$f(v) = \int \delta(v - |\vec{V}|) F_{GRF}(\vec{V} - \vec{v}_0 - \vec{v}_{\text{pec}} - \vec{v}_e) d^3 \vec{V}$$

- Mass Galaxy is in  $498 \text{ km/s} < v_{\text{max}} < 608 \text{ km/s}$  for which

$$F_{GRF} = \rho_{DM} v$$

- Several DM velocity distribution, they are correlated with DM density distribution
- Simplest : isothermal sphere model

$$F_{GRF}(\vec{V}) \sim \exp(-|\vec{V}|^2/\Delta V^2)\Theta(v_{max} - |\vec{V}|)$$

- Lead to

$$f(v) = c_{\text{norm}} \left[ \exp\left(-\frac{(v - v_1)^2}{\Delta V^2}\right) - \exp\left(-\frac{\min(v + v_1, v_{max})^2}{\Delta V^2}\right) \right]$$

$$\Delta V = v_0$$

- Note: Earth motion around Sun leads to 7% modulation effect of  $v_1$  and to modulation of signal in DD experiments
- DM velocity distribution near the sun could be quite different