



The Abdus Salam International Centre  
for Theoretical Physics



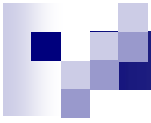
M.N.Kiselev

# Electron transport through nanostructures

## Lecture 1

**Quantum Dots: Coulomb blockade, tunneling, cotunneling**

Regional School on Physics at the Nanoscale, Hanoi, 14-25 December 2009



## Outline of the course:

- Quantum Dots

- Kondo effect in nano-devices
- From Fermi liquid to Luttinger liquid

### For reading:

Transport through QDs: W.G. van der Wiel et al, RMP 75 (2003)

SET and Coulomb blockade: M.A.Kastner, RMP 64 (1992)

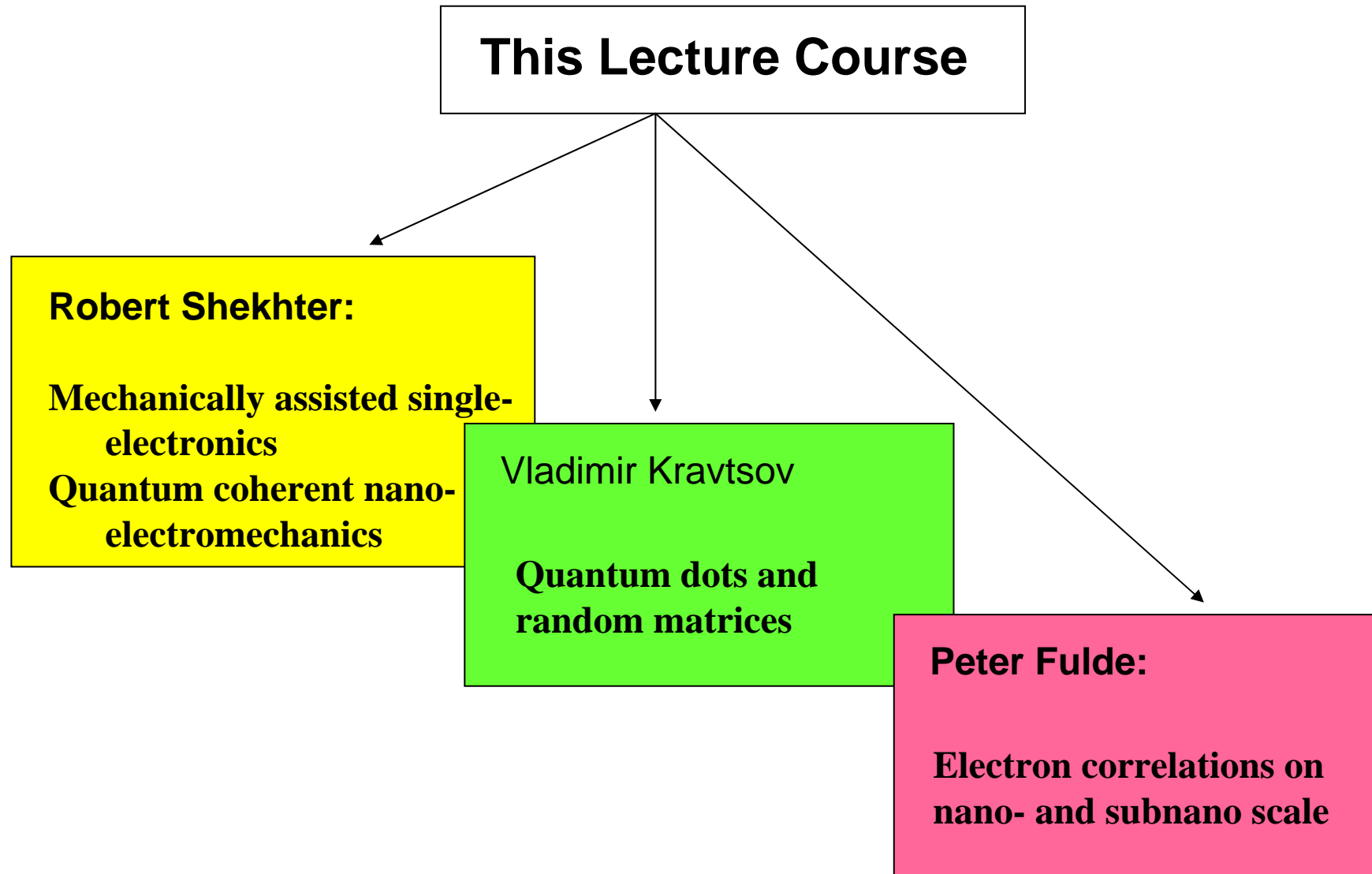
Popular reading: Leo Kouwenhoven and Charles Marcus, Physics World 1998

See also in the web Lecture courses of Ya. Blanter, Y. Gefen, Yu. Galperin

\* Some transparencies are courtesy of Yuval Gefen, Yaroslav Blanter and Yuri Galperin



## Connection with other lectures:



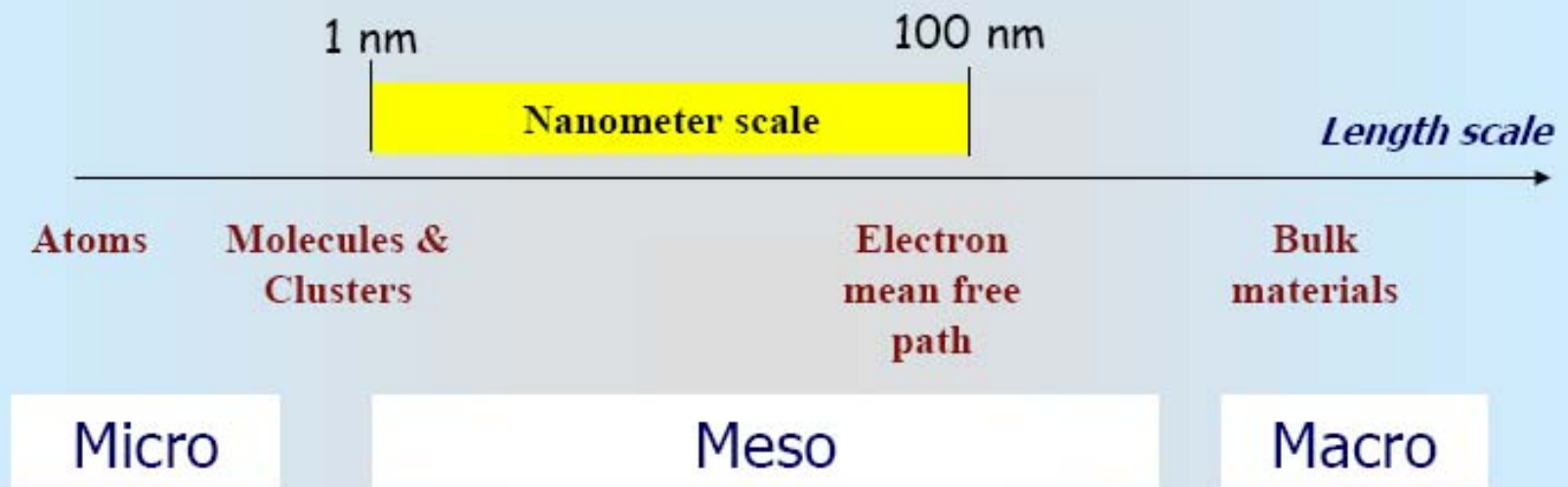
# Outline of this lecture



- What is Nano?
- Examples of QD:
- Vertical vs lateral QDs
- Diffusive vs ballistic QDs
- Metallic vs semiconductor QDs
- Open vs close QDs
- Coulomb blockade
- Sequential tunneling
- Elastic vs inelastic cotunneling
- "Universal" Hamiltonian

## Characteristic scales in nanoscience

$$1 \text{ nm} = 10^{-9} \text{ m}$$



Modern electronic devices belong to mesoscopic scale

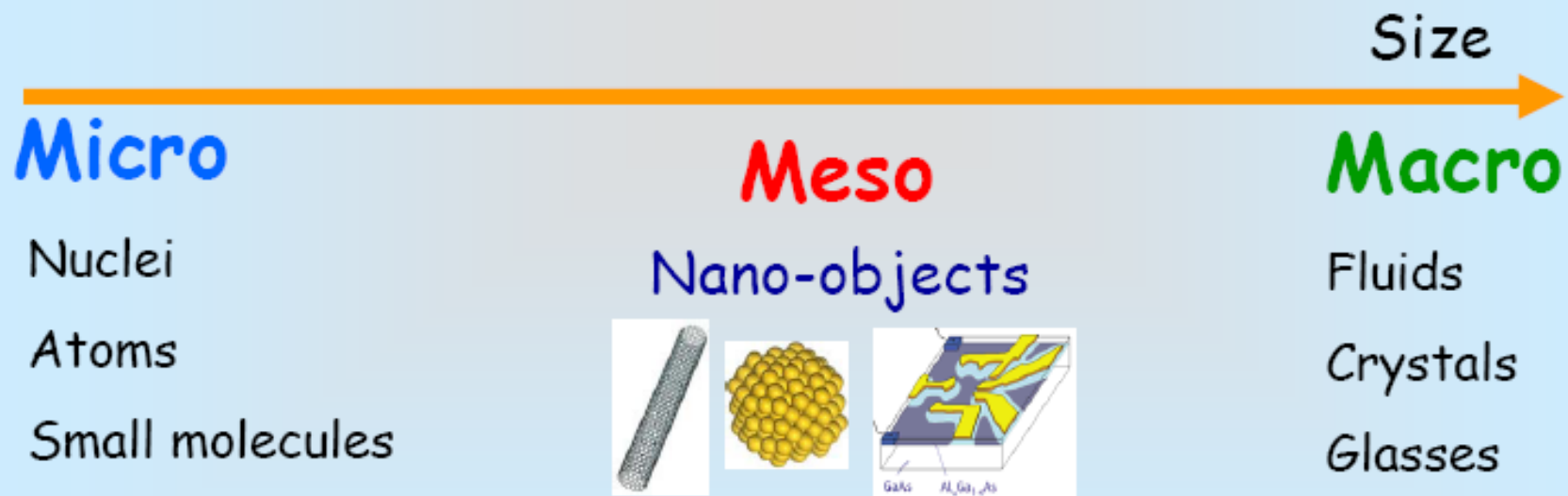
## Nano means Big !?

Nanoscale objects do not fully belong to the microcosm

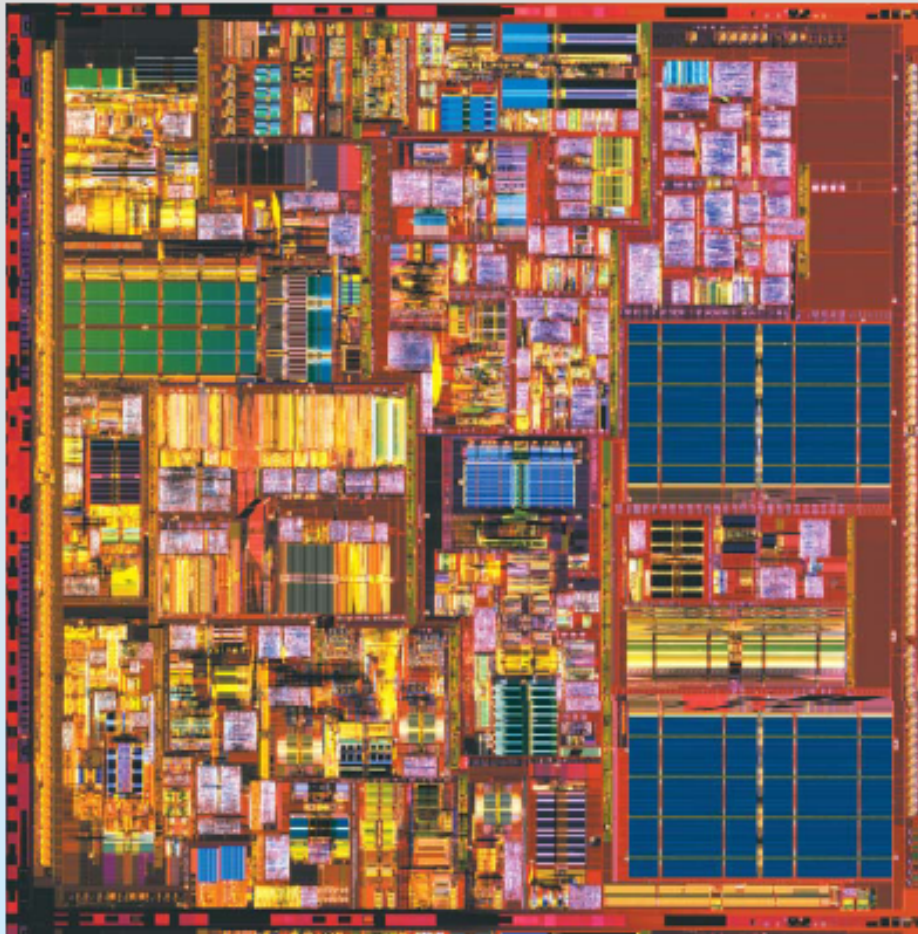
Many atoms, electrons, etc., are involved



Number of degrees of freedom is large



## CMOS TECHNOLOGY



Intel's Norwood (Pentium 4 - 130 nm) processor

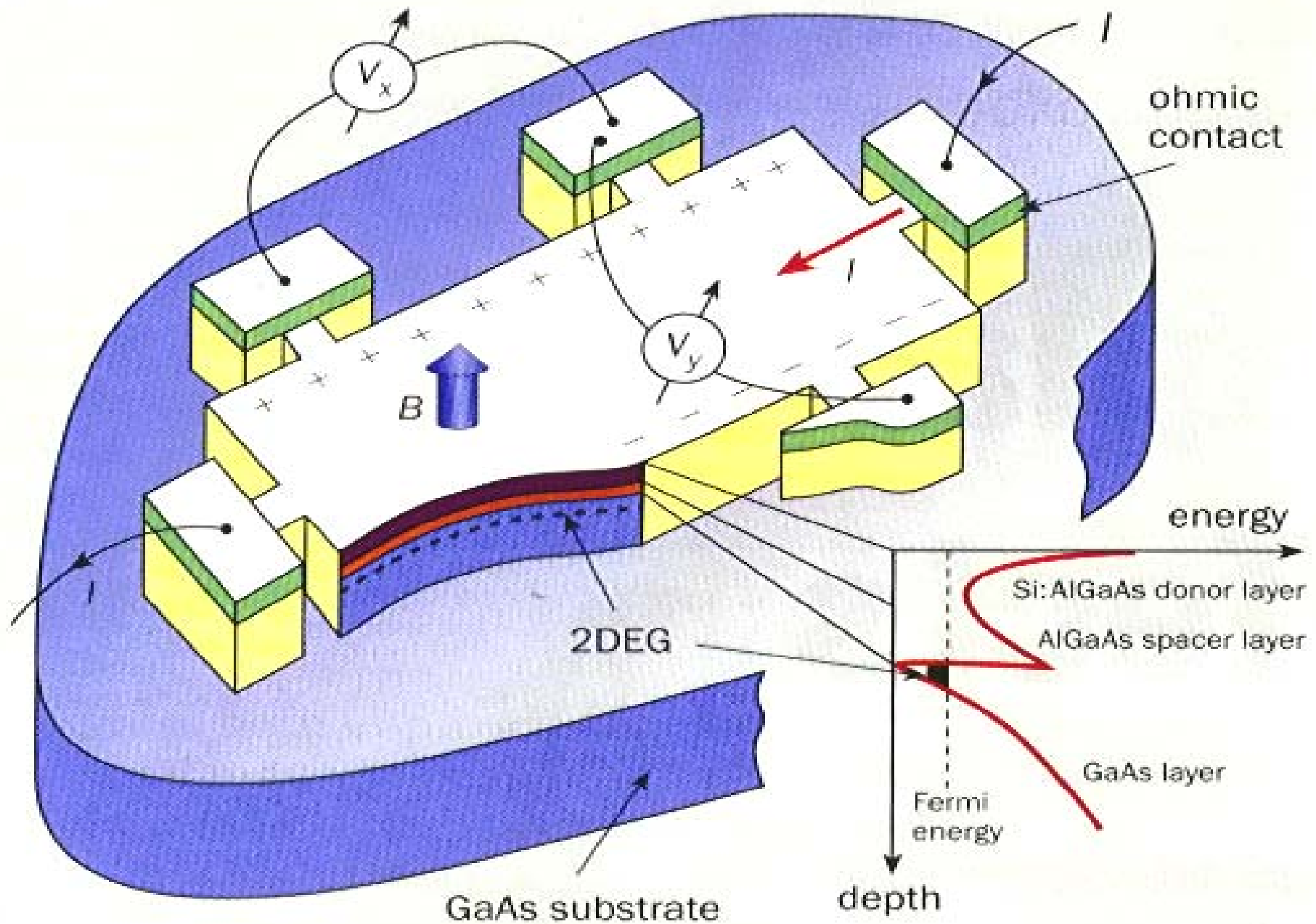
Intel's Prescott processor  
(released March 2004):

- 150 million transistors
- 90 nm design rules
- 3.4 GHz clock frequency

DRAM chips:

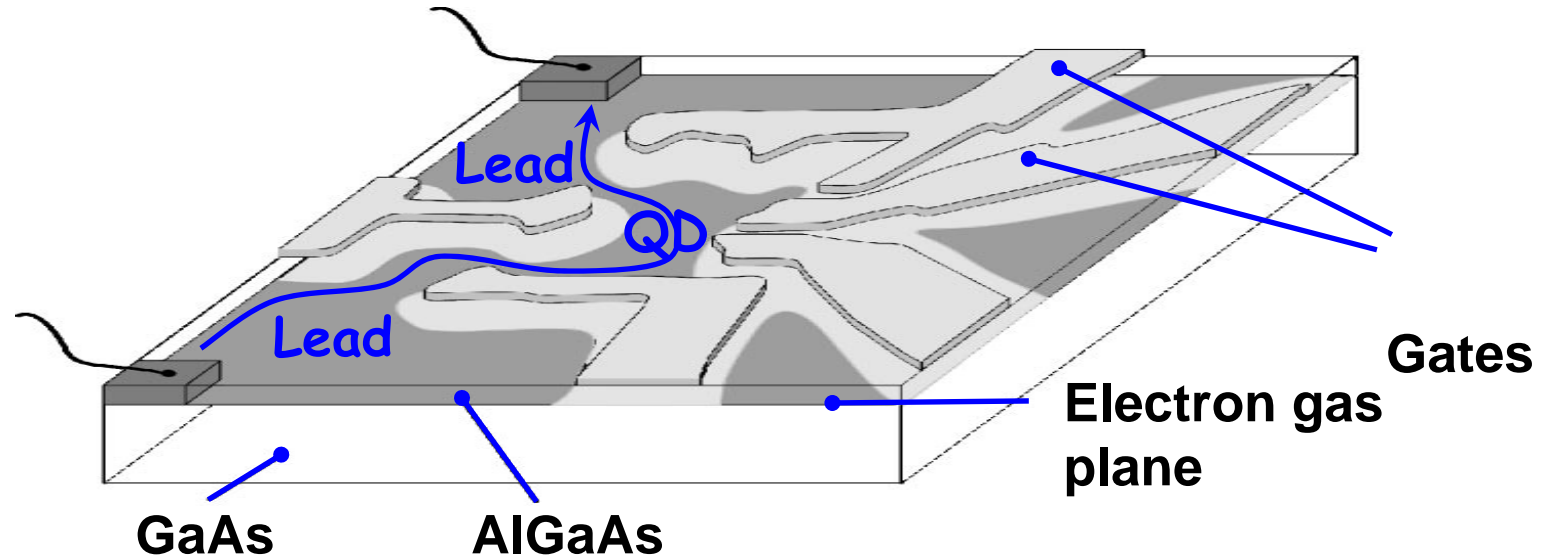
4 Gb chips demonstrated  
( $\sim 10^9$  transistors/cm<sup>2</sup>)

# Two Dimensional Electron Gas (2DEG)



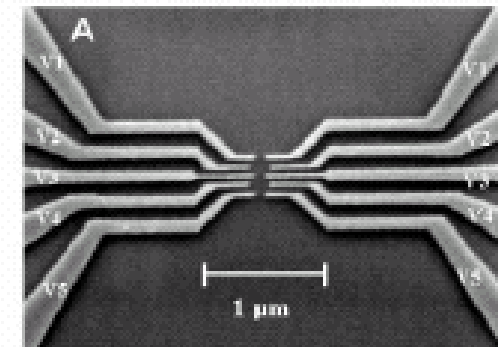
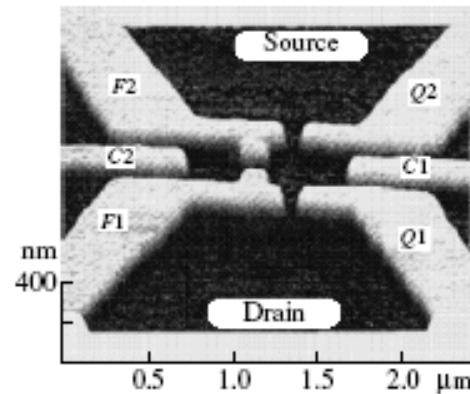
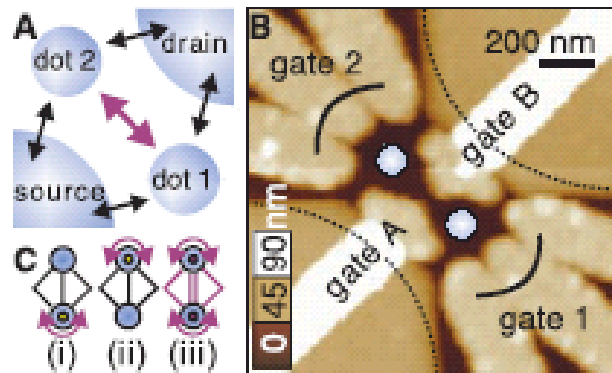
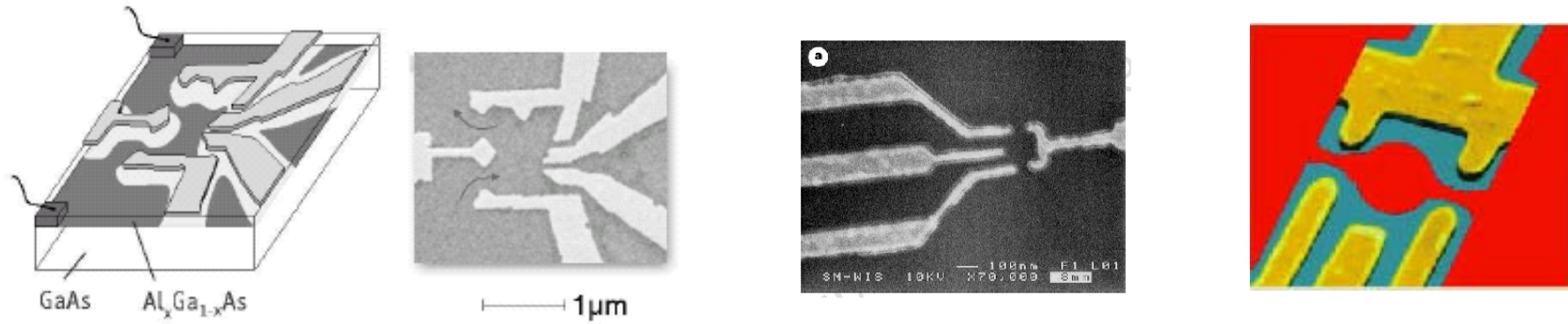


# Quantum dots



- Tune: gate potentials, temperature, field...
- Measure: I-V curves, conductance  $G$ ...
- Aharonov-Bohm interferometry, dephasing, coherent state manipulation...

# Quantum dots: from simple to complex



D.Goldhaber-Gordon et al (1998)

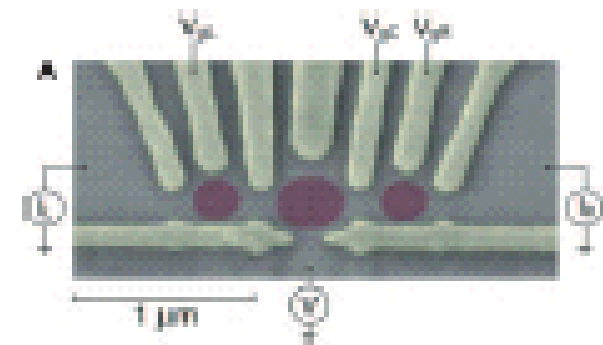
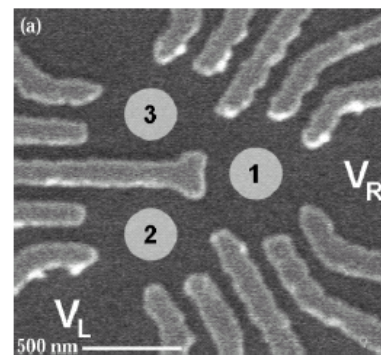
J.P.Kotthaus (1995)

A.Holleitner et al (2002)

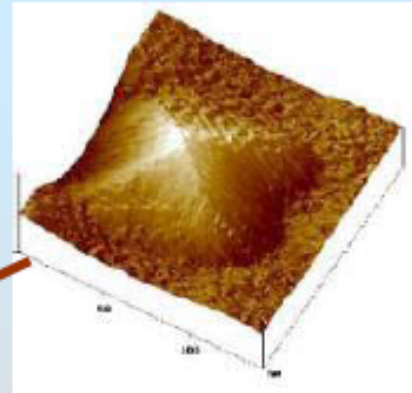
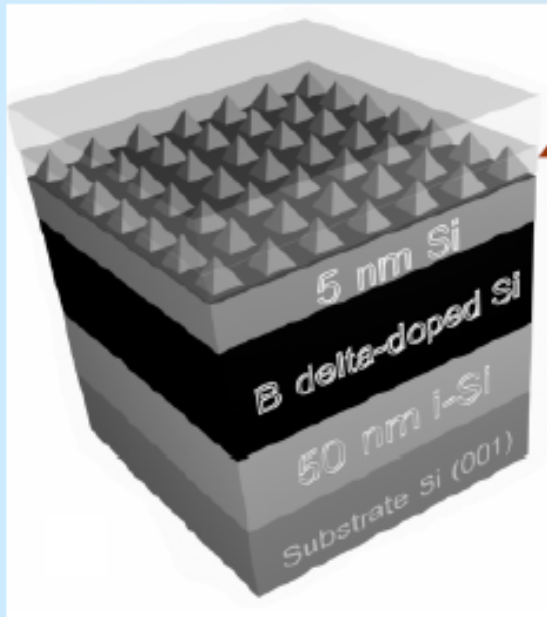
L.W.Molenkamp et al (1995)

H.Jeong et al (2001)

C.Marcus et al (2003)

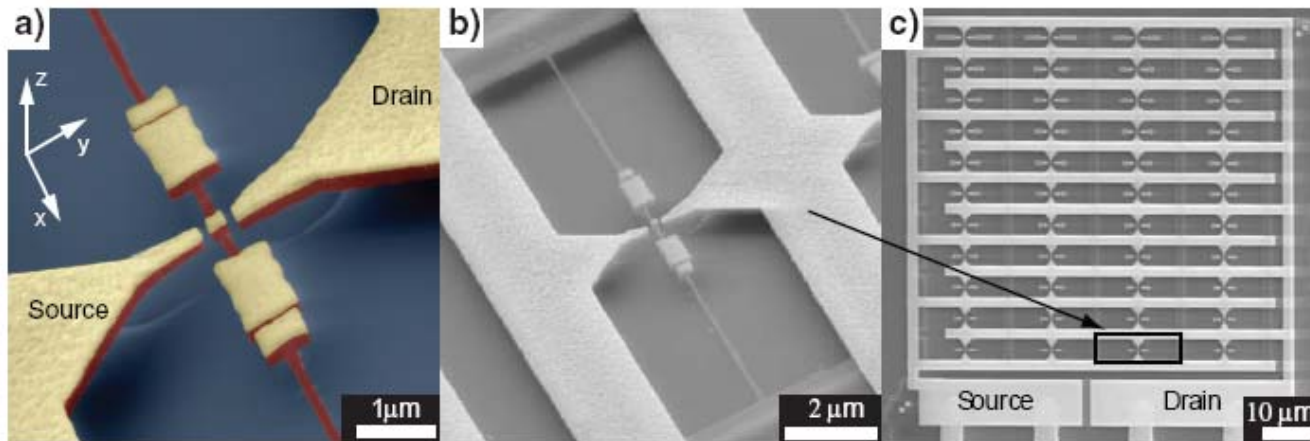
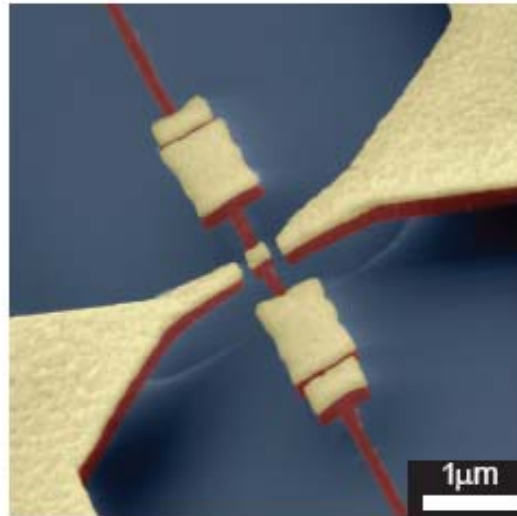


# Self-assembled QD



Self-assembled quantum dots are periodic arrays of "artificial atoms". They are considered to be promising systems for heterostructure lasers.

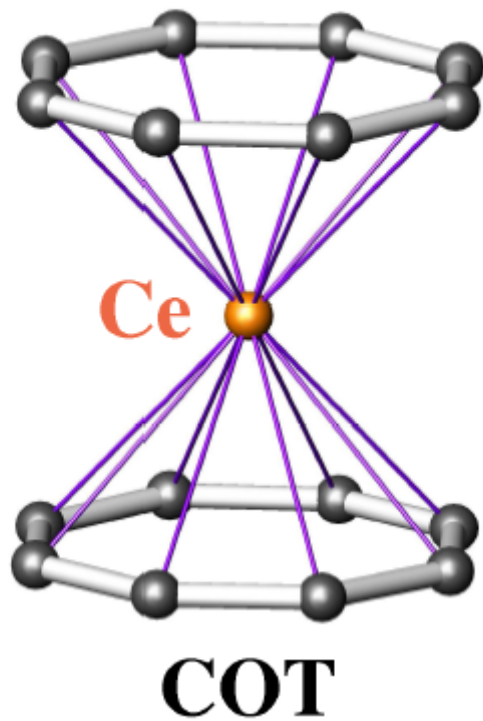
## Nanoelectromechanical shuttling: QD devices



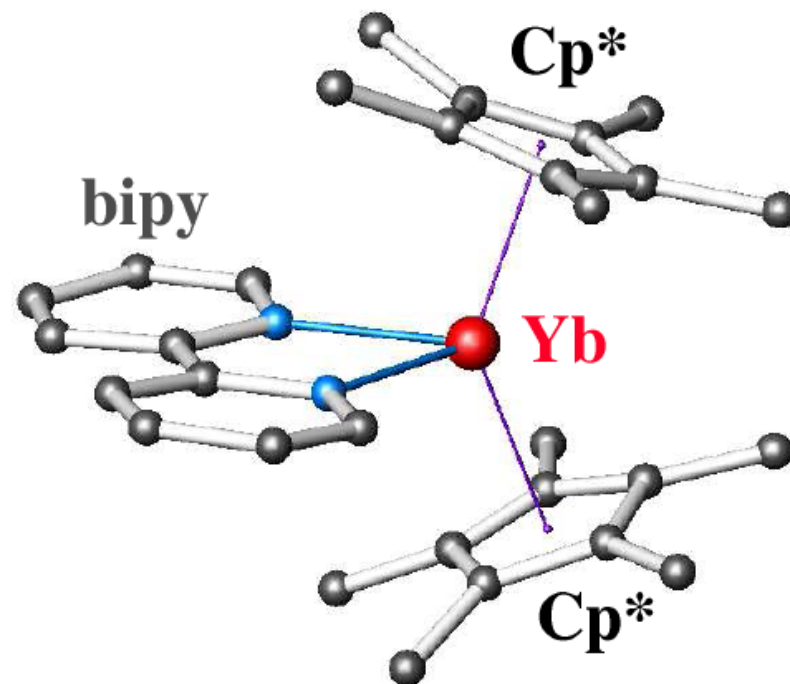
Cerocene



Ytterbocene



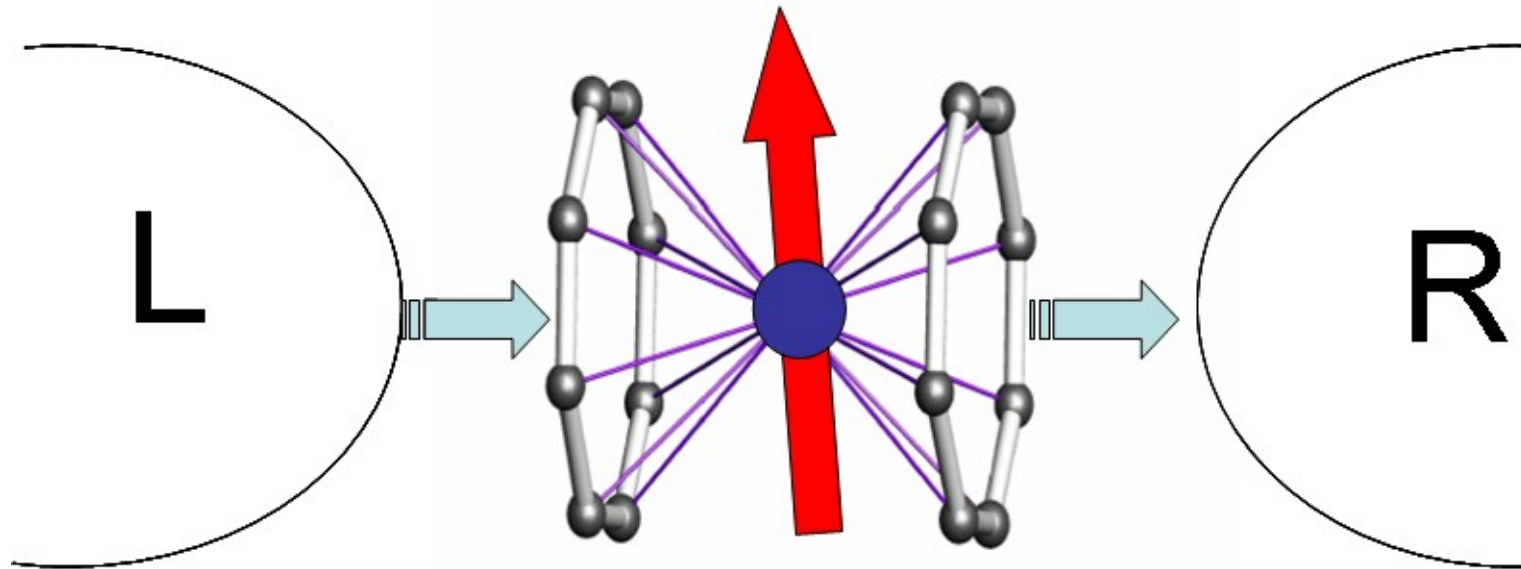
COT = C<sub>8</sub>H<sub>8</sub>



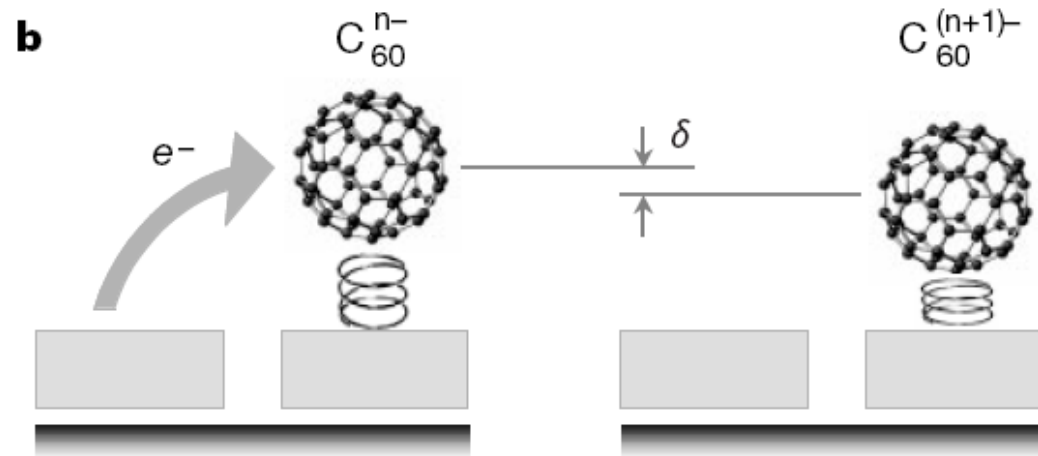
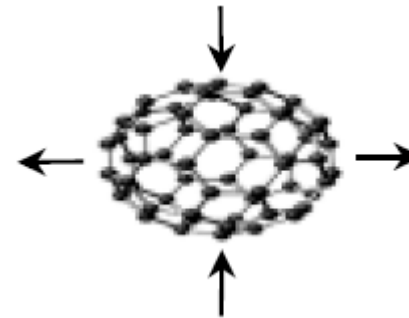
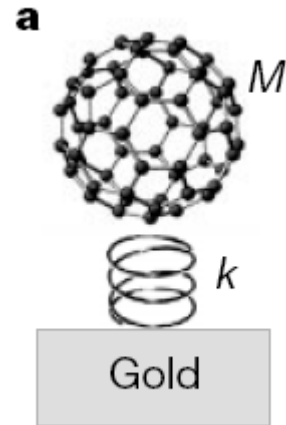
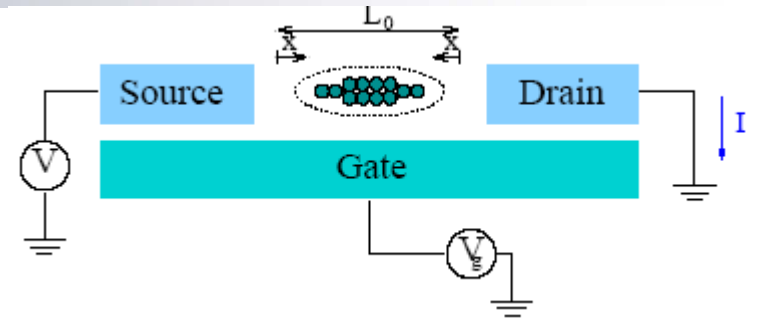
Cp\* = C<sub>5</sub>Me<sub>5</sub>,

bipy = (NC<sub>5</sub>H<sub>4</sub>)<sub>2</sub>

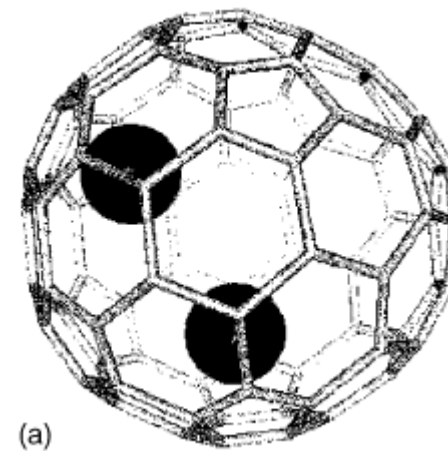
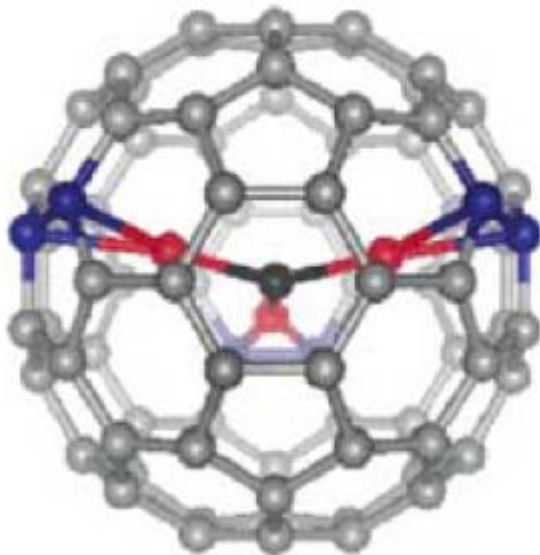
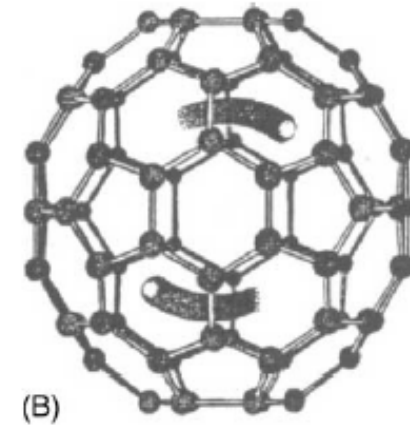
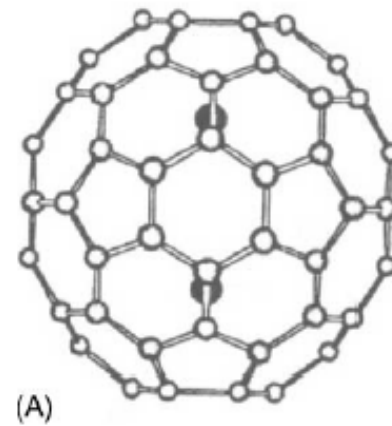
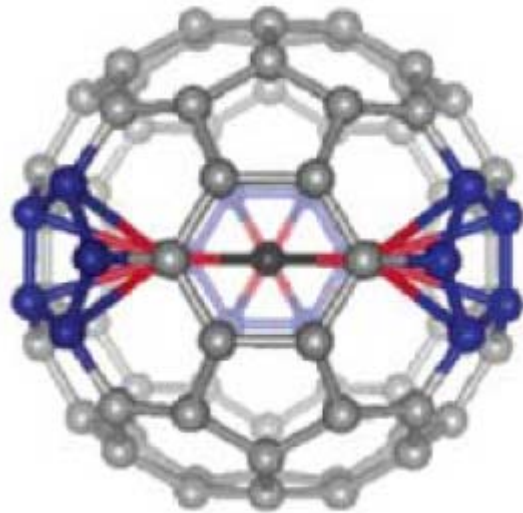
# Molecular Transistor



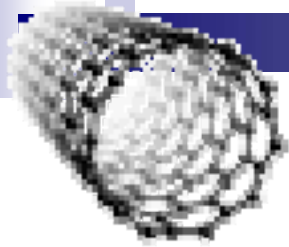
# Fullerenes



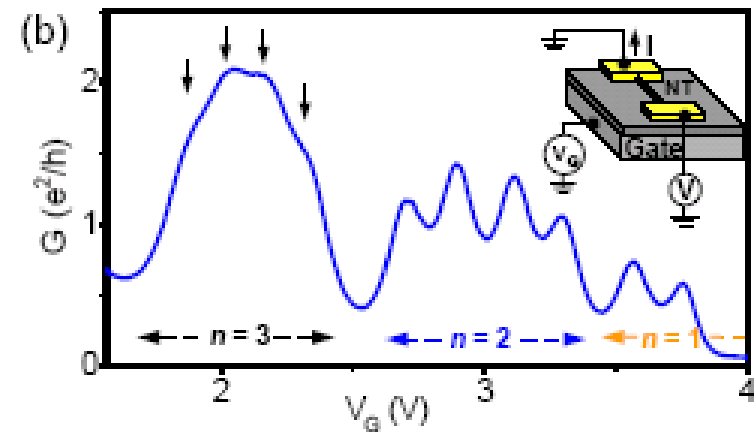
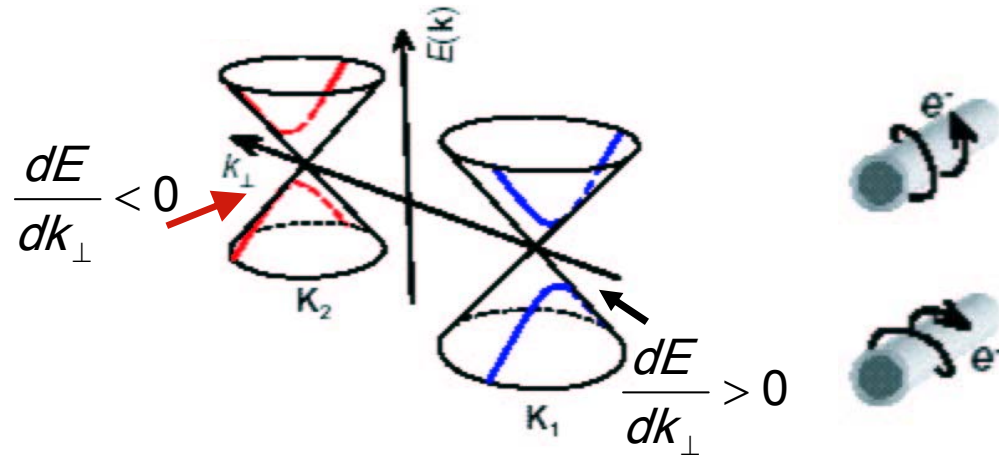
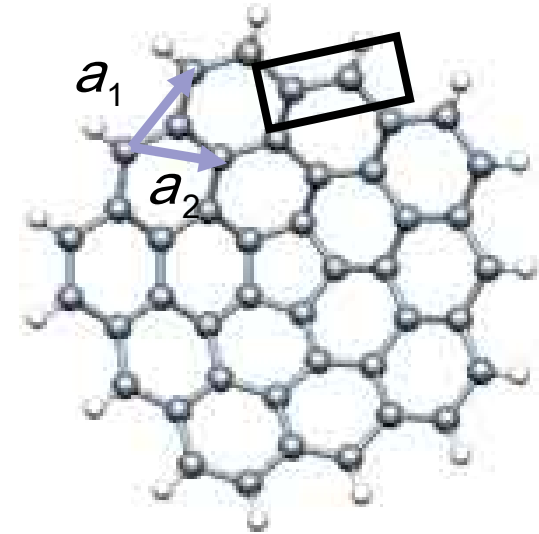
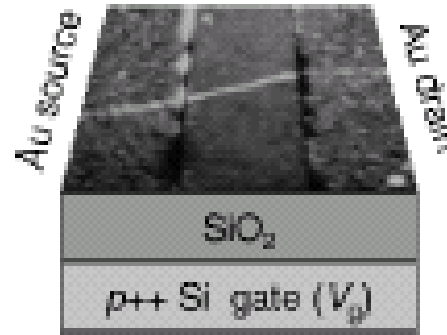
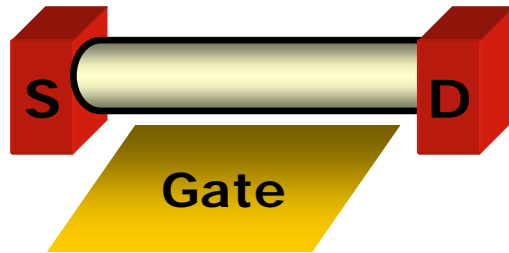
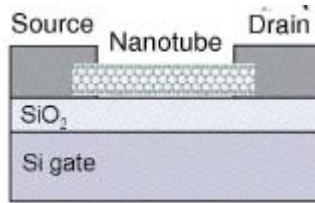
# Transition metals inside fullerenes



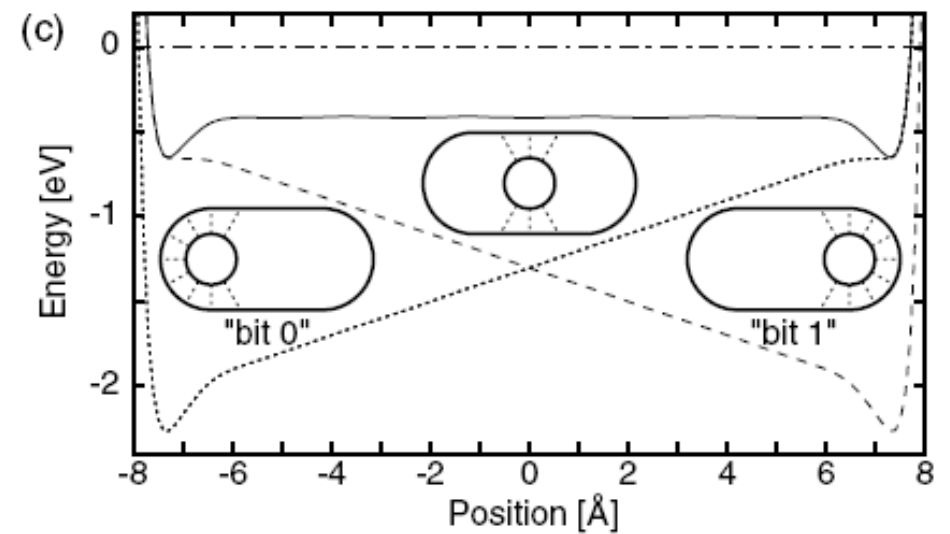
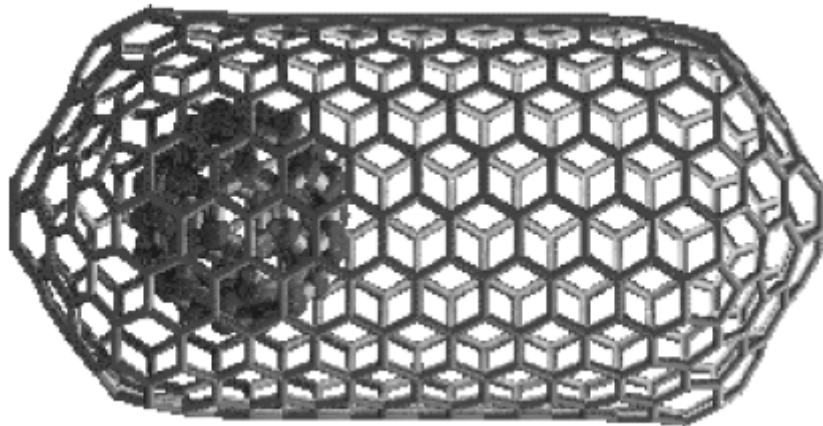
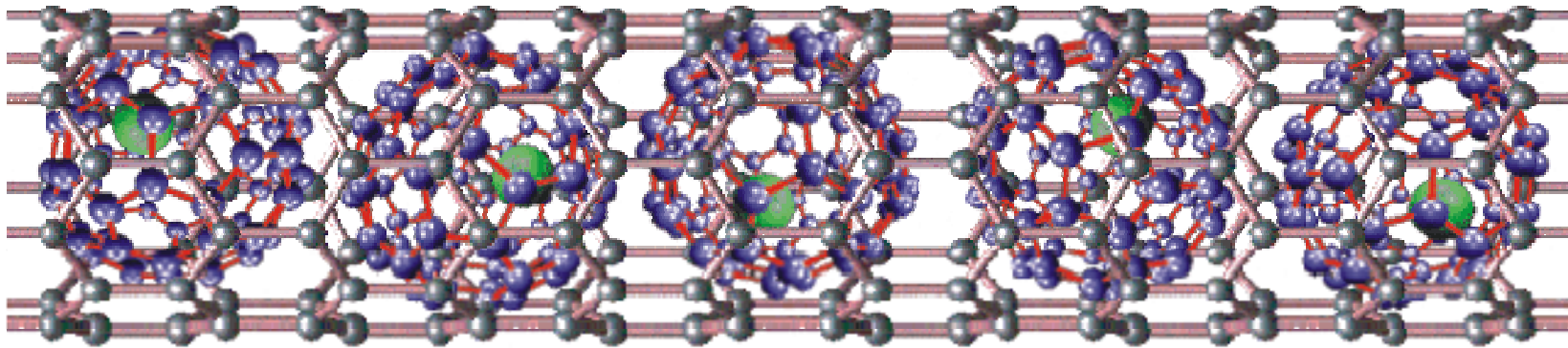




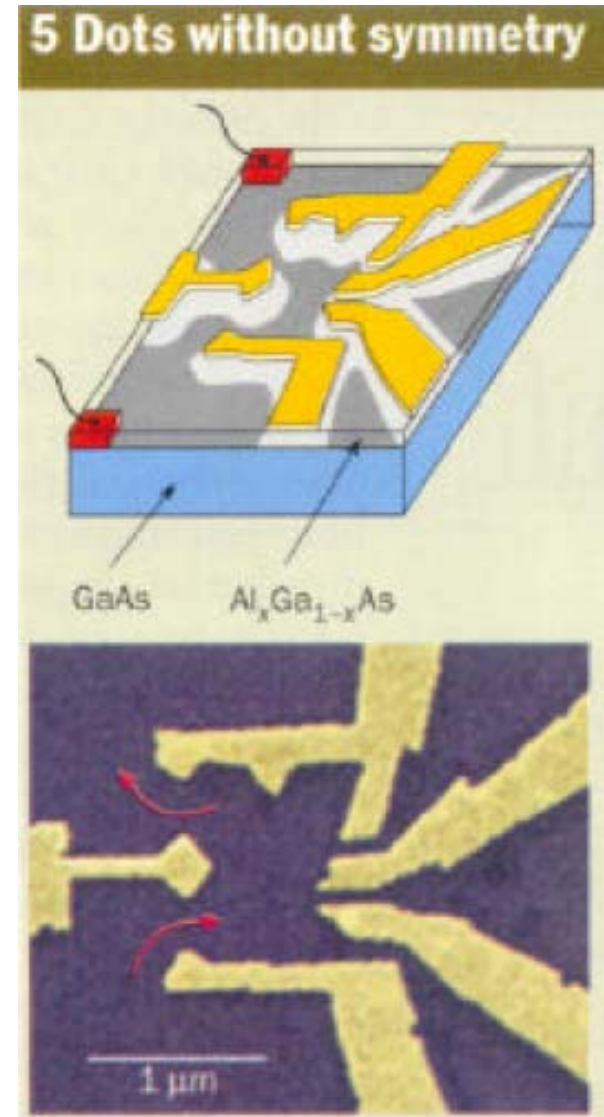
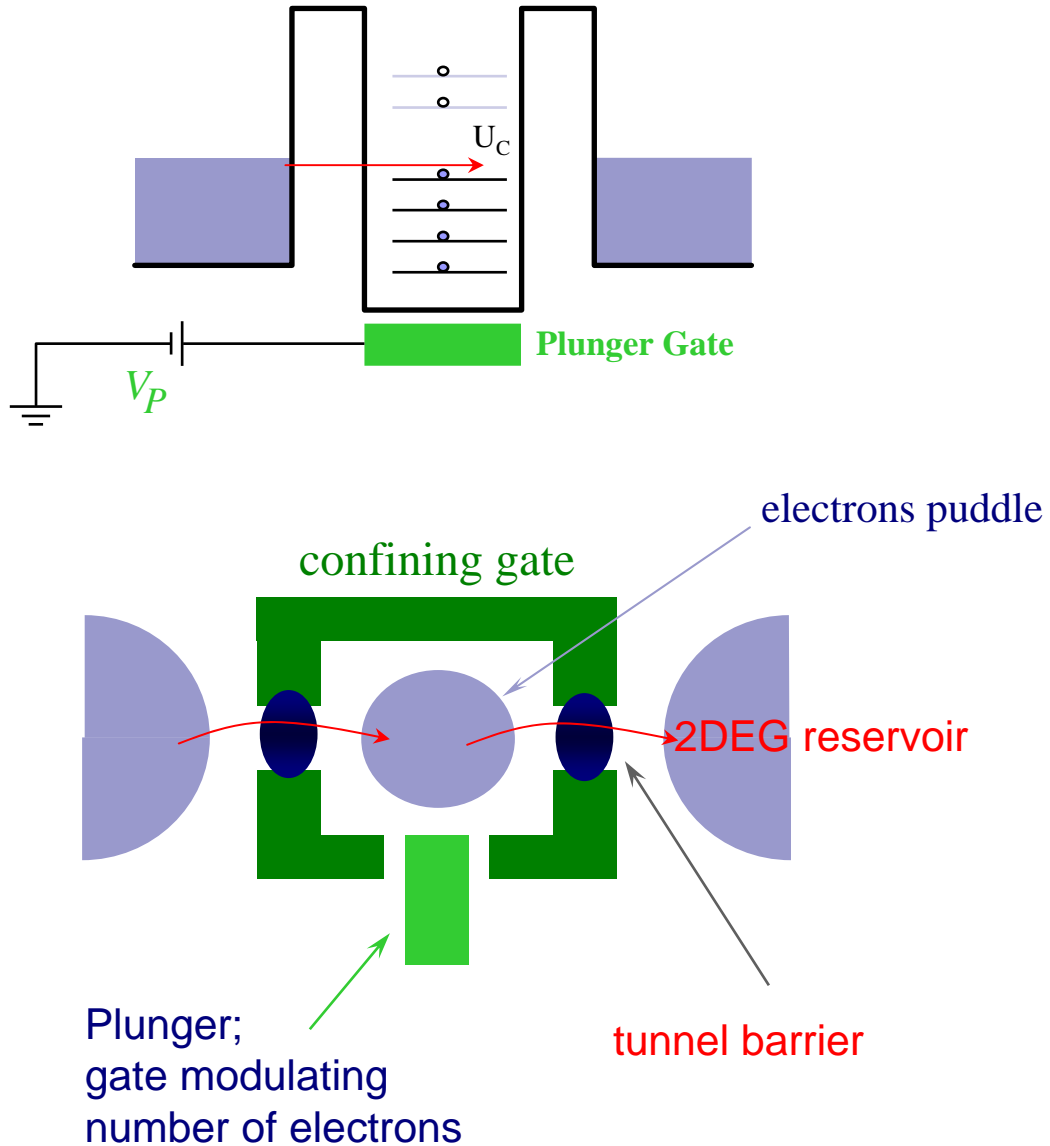
# Carbon Nanotubes



# Nanotube peapods: $C_{60}$ @ CNT



# Realization of lateral QD in 2DEG



## **Characteristic parameters:**

**size:**  $100 \text{ \AA} \rightarrow 2 \mu\text{m}$

**# electrons:**  $0 \rightarrow \text{hundreds}$

### **mobility**

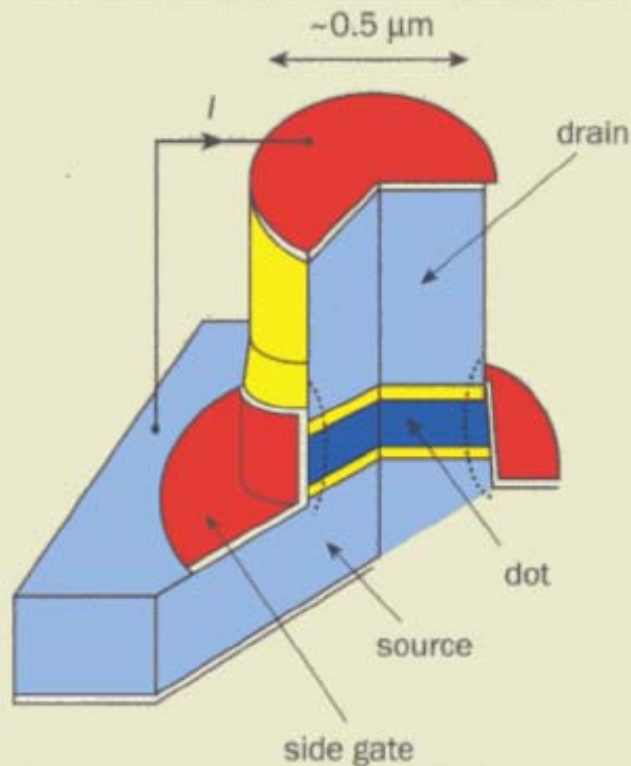
(of 2DEG in strong magnetic fields)  $\sim (30 - 50) \cdot 10^3 \text{ cm}^2 / \text{V} \cdot \text{sec}$   
original Integer Quantum Hall Effect  
current world record (Weizmann)  $\sim 36 \cdot 10^6$

## **CONTROL:**

- ✦ **size of QD**
- ✦ **density of electrons  $\rightarrow$  # electrons**
- ✦ **(mobility; disorder)**
- ✦ **shape**
- ✦ **contact to leads**

# Vertical QDs

## 1 Vertical quantum dot structure

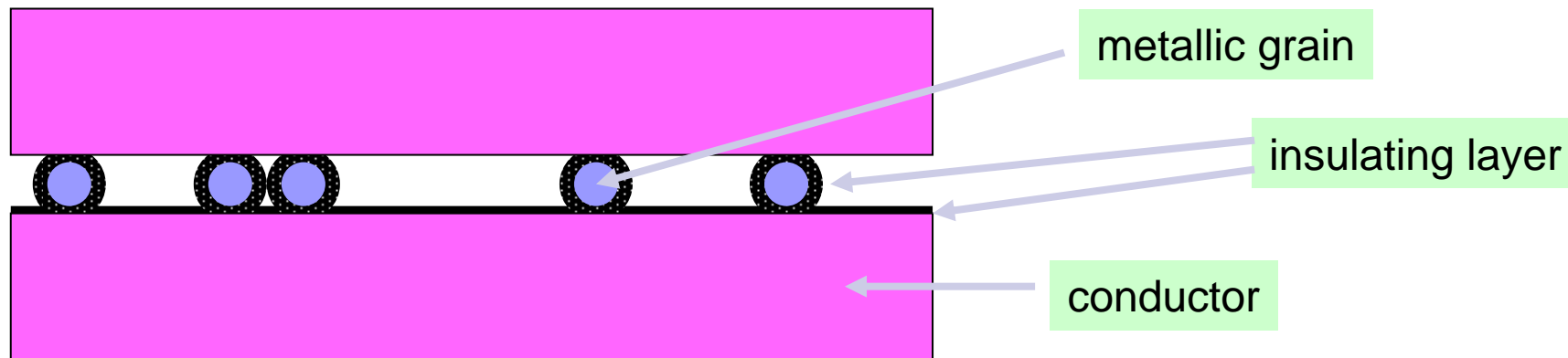
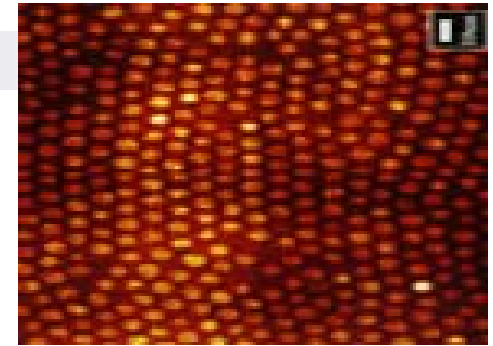


The quantum-dot structure studied at Delft and NTT in Japan is fabricated in the shape of a round pillar. The source and drain are doped semiconductor layers that conduct electricity, and are separated from the quantum dot by tunnel barriers 10 nm thick. When a negative voltage is applied to the metal side gate around the pillar, it reduces the diameter of the dot from about 500 nm to zero, causing electrons to leave the dot one at a time.

**advantages:** easy access to small # electrons, symmetric QDs

**disadvantages:** hard to control shape/size; dot-lead coupling

## Metallic QDs



**size:**  $30\text{\AA}$  and up

$\lambda_F$  : a few  $\text{\AA}$

**# electrons:**  $>$  many hundreds

*originally: statistics of an ensemble*

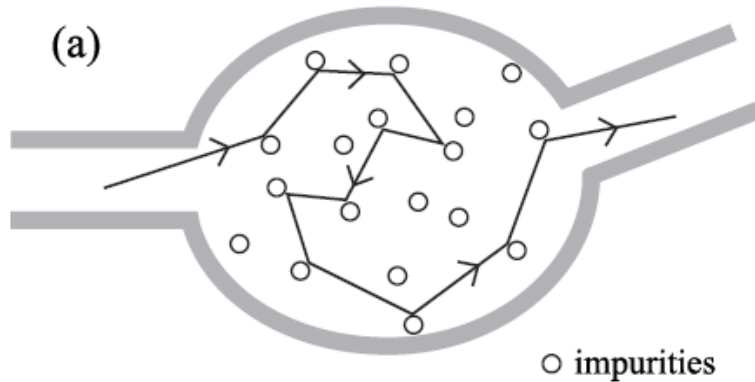
*today: can attach leads to a single QD*

*little control: QD-lead coupling; size of QD*

*special appeal: QDs with special properties: SC; magnetic...*

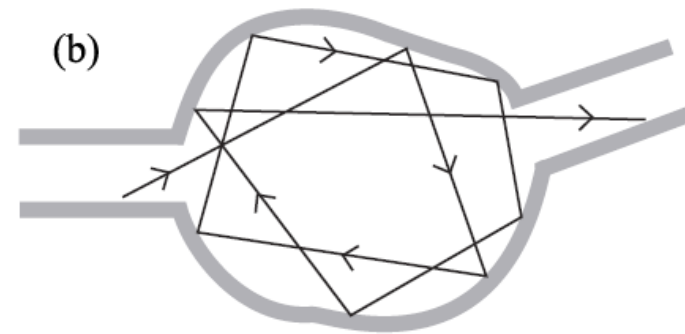
# NON INTERACTING ELECTRONS

## *diffusive vs. ballistic*



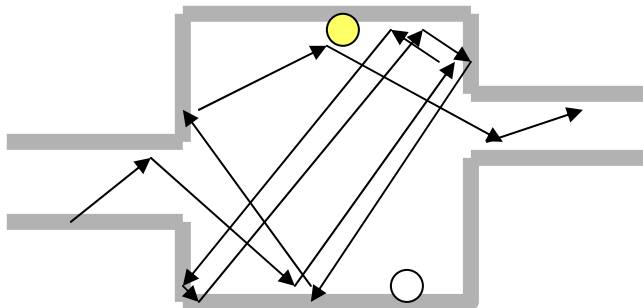
**diffusive**

$$E_{th} = \hbar / (\text{diffusion time}) = \hbar / (L^2 / D)$$



**ballistic**

$$E_{th} = \hbar / (\text{time of flight}) = \hbar / (L / v_F)$$

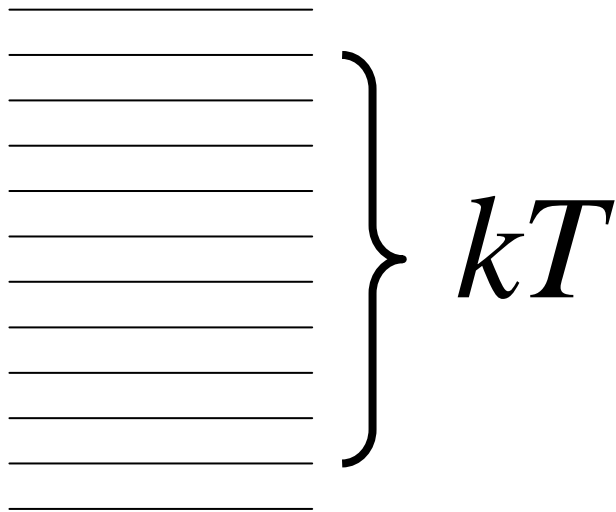


**dirty ballistic**

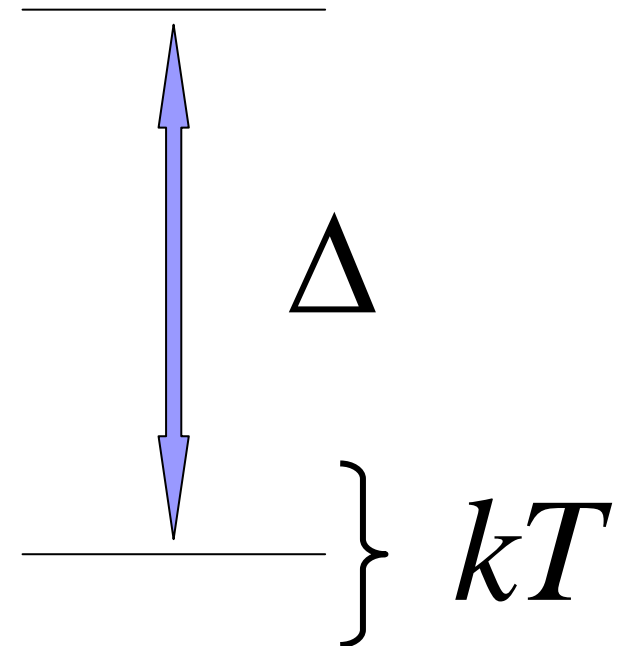
time of flight  $\neq$  Thouless energy  
( Altland , Gefen, Montambaux)

# *“Metallic” vs. discrete spectrum*

$$\Delta \ll kT$$



$$\Delta \gg kT$$







## *Open vs. Closed*

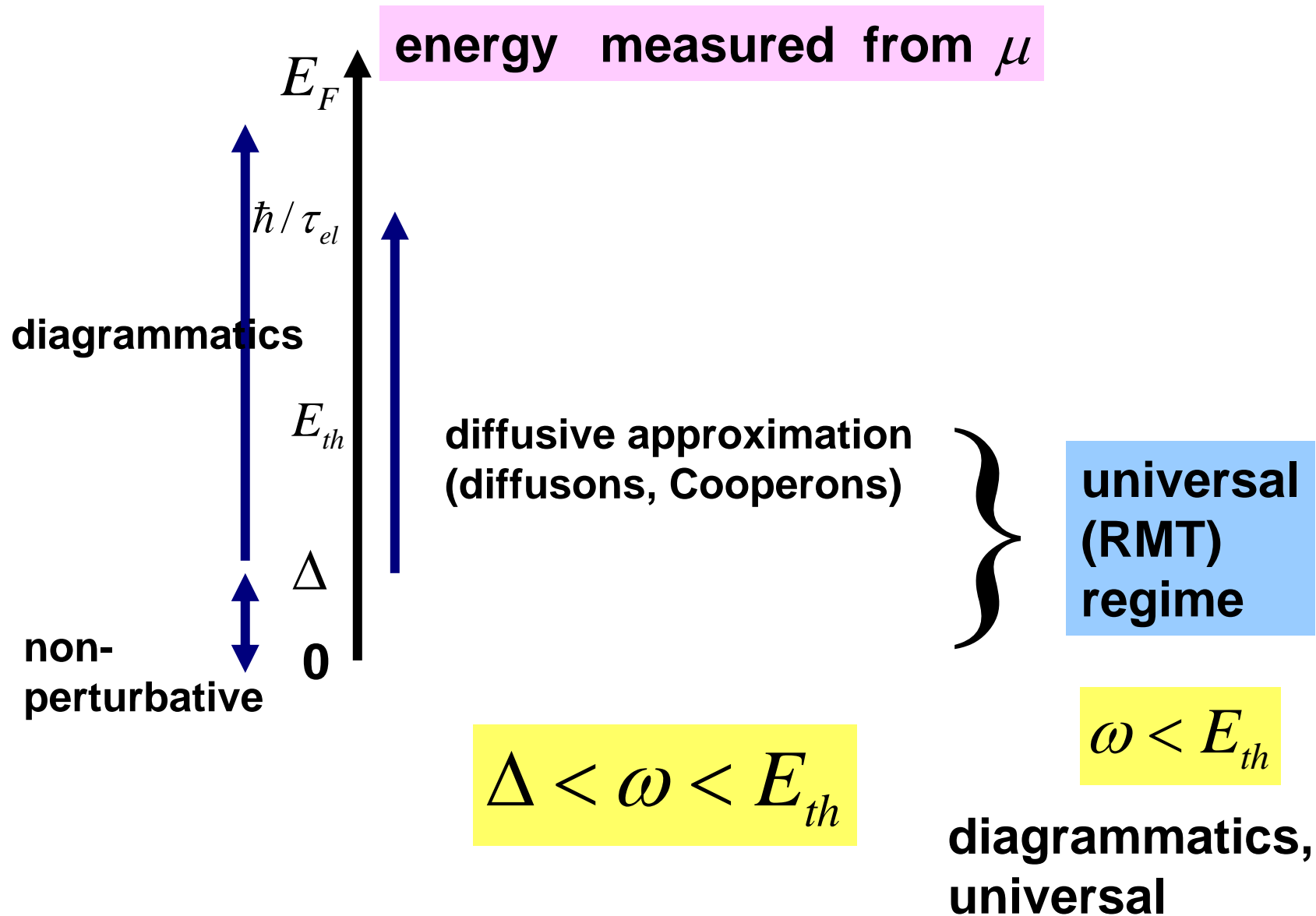
$\frac{\Gamma_c}{\hbar}$  = decay rate of a QD level into channel c

$$\text{total level width} = \Gamma = \sum_c \Gamma_c$$

closed QD (charge on the dot is nearly quantized)

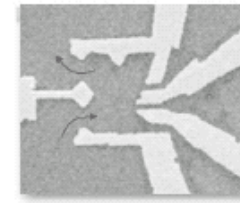
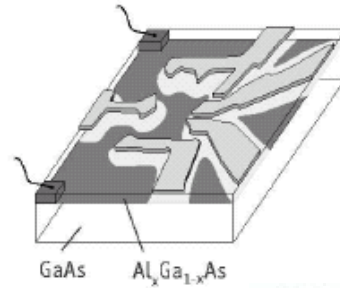
$$\Gamma < E_c$$

# NON INTERACTING ELECTRONS



# Metallic quantum dots: many-electron system

Random Matrix Theory



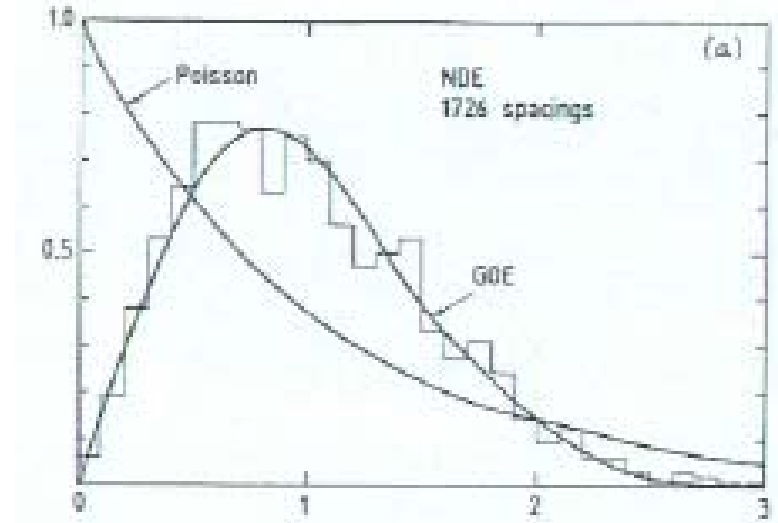
= artificial atom



GOE

GUE

Wigner-Dyson statistics



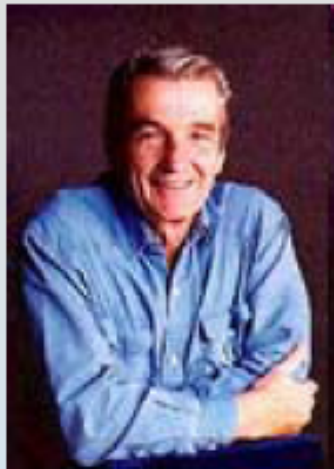
$$P(\{E_\mu\}) \propto \exp\left(\frac{\beta}{2} \sum_{\nu \neq \mu} \ln \left[ \frac{|E_\mu - E_\nu|}{\delta} \right]\right)$$

$\beta = 1$  Orthogonal (GOE)

$\beta = 2$  Unitary (GUE)

$\beta = 4$  Symplectic (GSE)

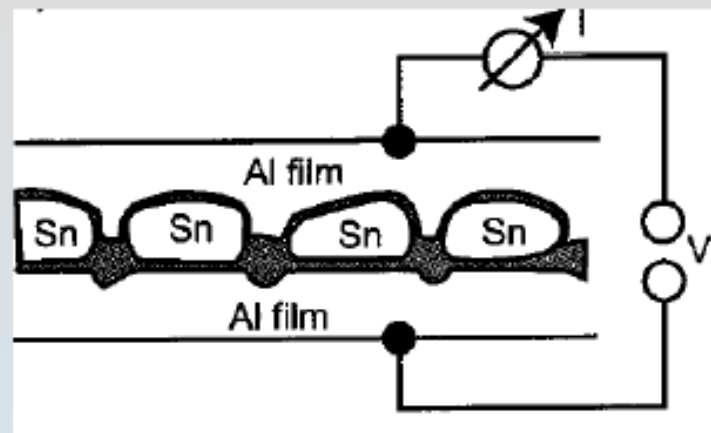
# Coulomb blockage



Ivar Giaever



1973



To move an electron to a confined region one has to pay for its repulsion from existing electrons

# The principle of the Coulomb blockade

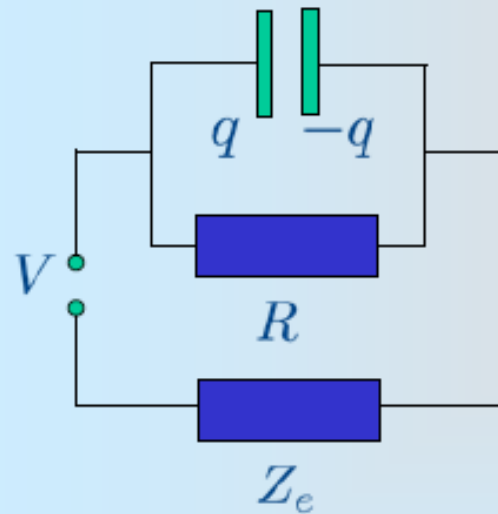


Why R matters?

time delay  $\delta t = eR/V$

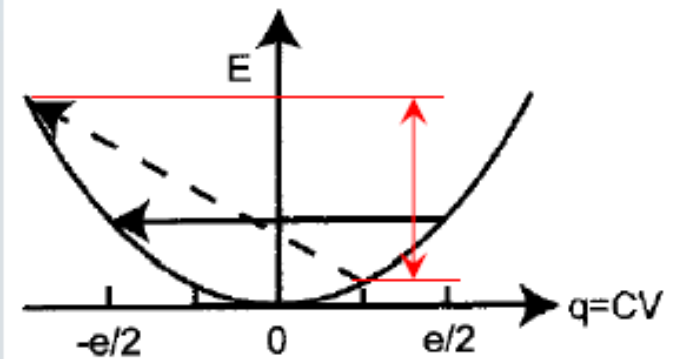
duration  $\tau \sim \hbar/eV$

$\delta t \gg \tau \rightarrow R \gg \hbar/e^2$



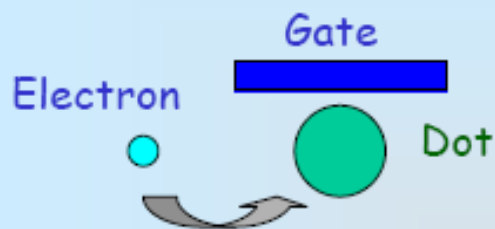
Because of environment capacitances it is difficult to observe CB in single junctions

Energy stored is  $q^2/2C$



At  $|q| < e/2$  the electron tunneling will increase the energy stored in the barrier - one has to pay for the tunneling by the bias voltage

# Coulomb blockade



$$Q = -Ne$$

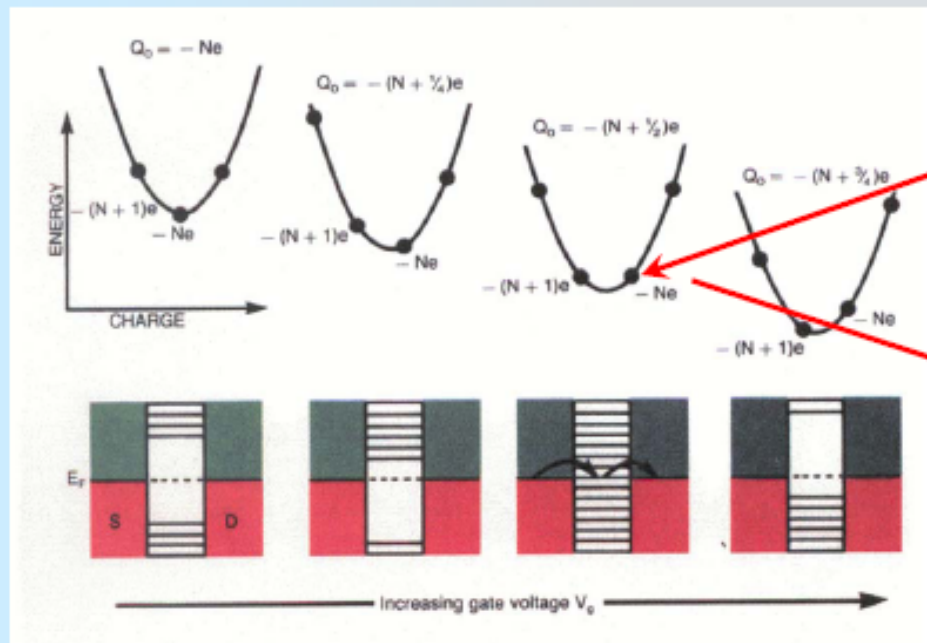
Cost

$$E = QV_g + \frac{Q^2}{2C}$$

Repulsion at the dot



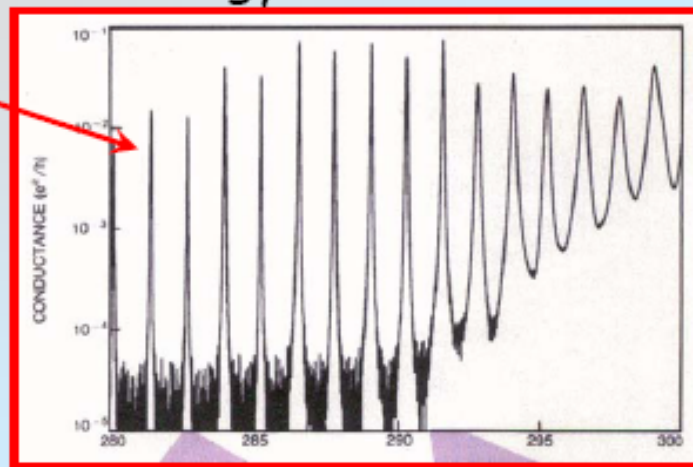
Attraction to the gate



At

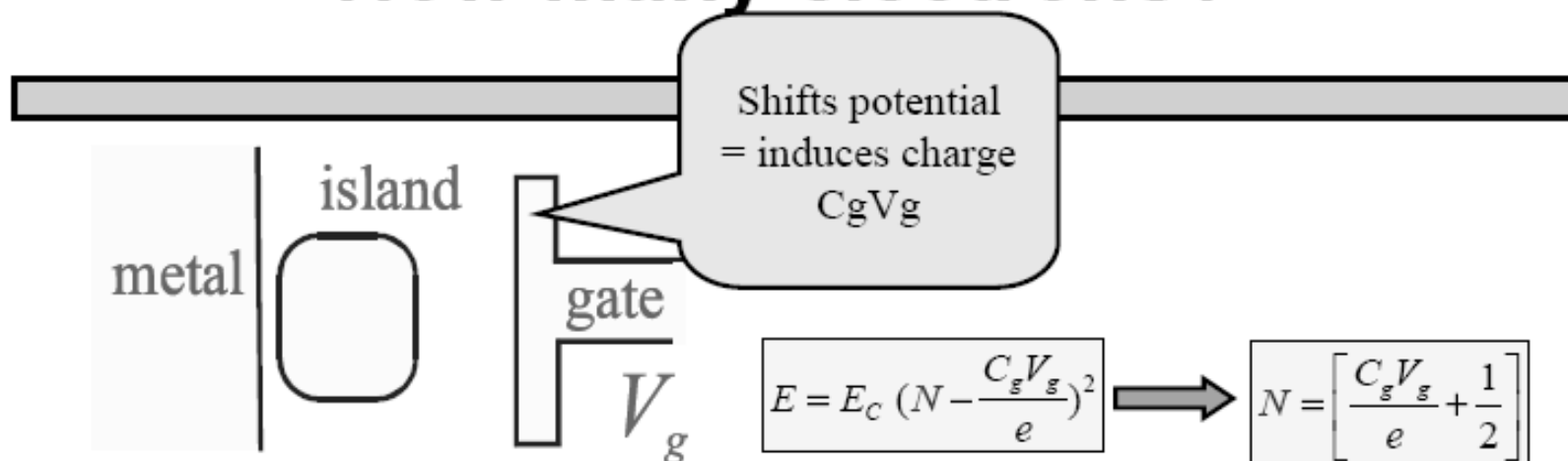
$$V_g = - \left( N + \frac{1}{2} \right) \frac{e}{C}$$

the energy cost vanishes !

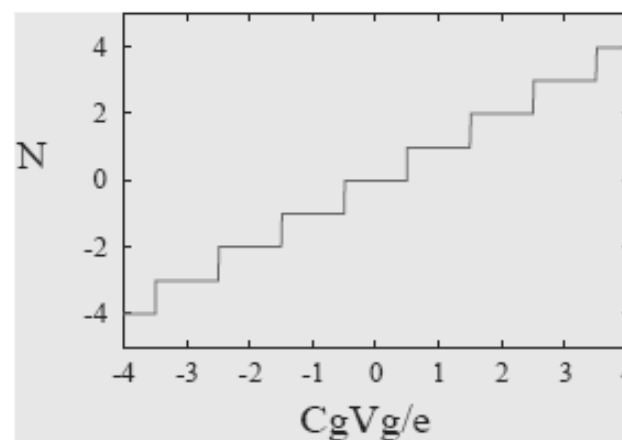
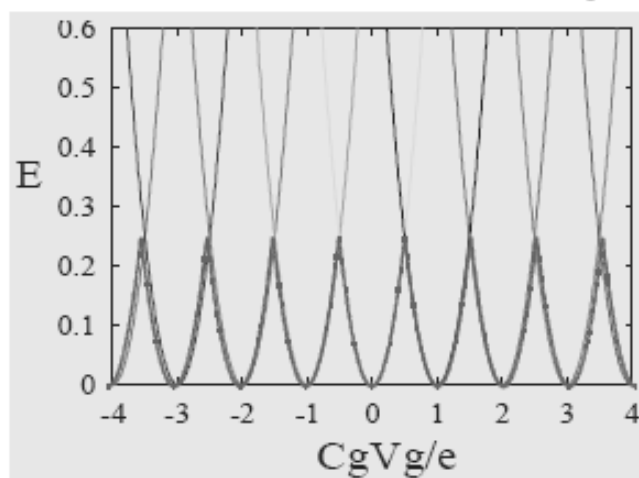


Single-electron transistor (SET)

# How many electrons?

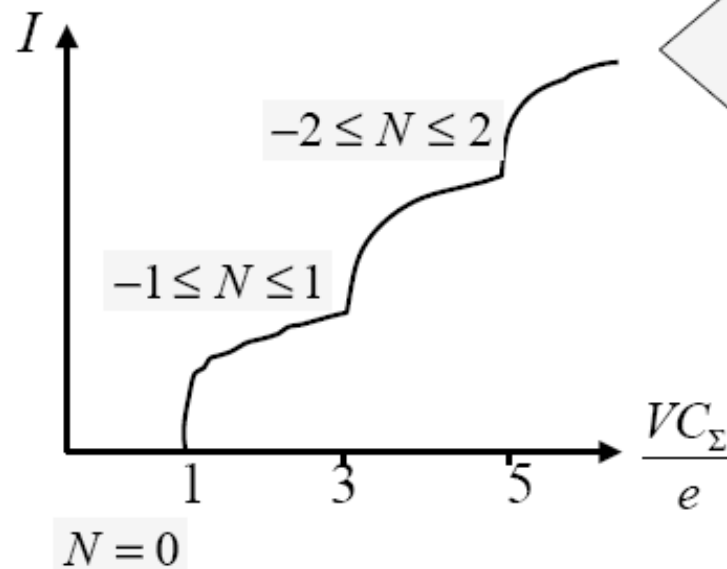


$$E = E_c \left( N - \frac{C_g V_g}{e} \right)^2 \longrightarrow N = \left[ \frac{C_g V_g}{e} + \frac{1}{2} \right]$$



## Coulomb staircase

Nonlinear transport:  
upon increasing  $V$  more charging states become  
available



Ohm's law:  $I = GV$

$$V_{th,n} = (2n + 1)e / C_\Sigma$$

$$Q_g = 0$$





## ***LATERAL QDs : possible parameters***

Temperature  $< 1$  K ( as low as 10-30 mK)

elastic mean free path  $\approx 1 - 150 \mu\text{m}$

$n_s \approx 10^{11} - 10^{12} \text{ cm}^{-2}$

$E_F \approx 10 - 20 \text{ meV}$

$\lambda_F \approx 50 \text{ nm}$

# electrons: 0 - hundreds

single particle level spacing  $= \Delta \sim (0.01 \text{ meV} \sim 0.1 \text{ K})$

Thouless energy  $= E_{\text{th}} \sim (0.3 \text{ meV} \sim 3 \text{ K})$

charging energy  $= E_c \sim (1 \text{ meV} \sim 10 \text{ K})$

# THE “UNIVERSAL” HAMILTONIAN

$$H = H_{sp} + H_{int}$$

$$H_{int} = \frac{1}{2} \sum_{\alpha, \beta, \gamma, \delta} H_{\alpha\beta\gamma\delta} \hat{\Psi}_{\alpha\sigma_1}^\dagger \hat{\Psi}_{\beta\sigma_2}^\dagger \hat{\Psi}_{\gamma\sigma_2} \hat{\Psi}_{\delta\sigma_1}$$

$$H_{\alpha\beta\gamma\delta} = \int d\mathbf{r}_1 d\mathbf{r}_2 V(\mathbf{r}_1 - \mathbf{r}_2) \phi_\alpha(\mathbf{r}_1) \phi_\beta(\mathbf{r}_2) \phi_\gamma^*(\mathbf{r}_2) \phi_\delta^*(\mathbf{r}_1)$$

Note: only orbital indices ( no spin-orbit)

$$H_{int} = H_{int}^{(0)} + H_{int}^{(1/g)}$$

↑  
universal

↑  
non-universal, fluctuating

# CHARGING HAMILTONIAN

$$H = H_{sp} + H_{int}$$

$$H_{int} = H_{int}^{(0)} + H_{int}^{(1/g)}$$

$$H_{int}^{(0)} \Rightarrow E_c (\hat{n} - N_0)^2$$

$$\rightarrow E_c \left( \sum \Psi_{\alpha}^{\dagger} \Psi_{\alpha} \right) \cdot \left( \sum \Psi_{\beta}^{\dagger} \Psi_{\beta} \right) - 2E_c N_0 + E_c N_0^2$$

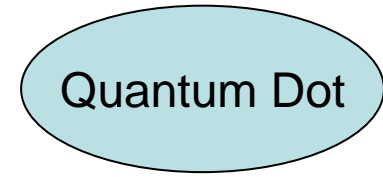
↑  
interaction

↑  
external gate V

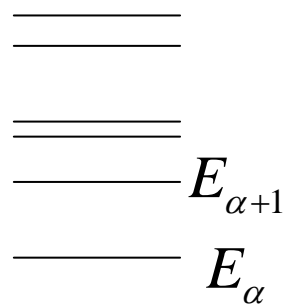
↑  
constant

# Metallic Quantum Dot: Universal Hamiltonian

Metallic grain or small island of electron gas



## Electron-electron interactions in isolated metallic grains



Mean-level spacing

$$\Delta = \langle E_{\alpha+1} - E_{\alpha} \rangle \quad (\text{kinetic energy})$$

Thouless energy

$$E_T \sim D \cdot L^{-2} \quad \text{diffusive regime}$$

$$E_T \sim v_F L^{-1} \quad \text{ballistic regime}$$

$$g = E_T / \Delta \gg 1$$

**metallic grain**

**GUE**

$$H_0 = \sum_{\alpha} E_{\alpha} n_{\alpha}$$

$$E_c = \frac{e^2}{2C}$$

$$H_{\text{int}} = E_c (\hat{n} - N)^2 - J (\vec{S})^2 - \lambda_{\text{BCS}} \hat{T}^{\dagger} \hat{T}$$

$$\hat{n} = \sum_{\alpha, \sigma} d_{\alpha\sigma}^{\dagger} d_{\alpha\sigma} \quad S^{\gamma} = \frac{1}{2} \sum_{\alpha, \sigma, \sigma'} d_{\alpha\sigma}^{\dagger} \sigma_{\sigma\sigma}^{\gamma} d_{\alpha\sigma'} \quad T = \sum_{\alpha} d_{\alpha\uparrow} d_{\alpha\downarrow}$$

charge                      spin                      superconducting

**Coulomb blockade**

Short-range interaction                       $E_c = 4 |J| \sim \Delta$   
 Scaling:

Coulomb interaction                       $E_c = r_s (k_F L)^{d-1} \Delta \gg |J|$

Kurland, Aleiner, Altshuler (2000)  
 Aleiner, Brouwer, Glazman (2002)

# TRANSPORT THROUGH a QD: thermally activated conduction

$$H = H_{sp} + E_c (\hat{n} - N_0)^2 + H_{leads} + H_{tunneling}$$

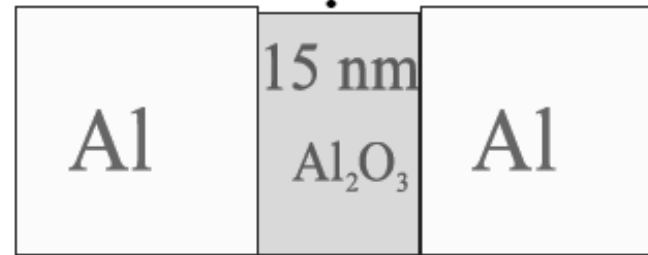
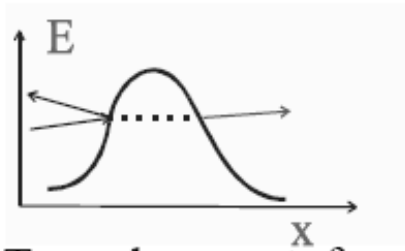
$$H_{leads} = \sum \xi_k c_k^\dagger c_k \quad \text{for each lead}$$

$$H_{tunneling} = \sum_{\alpha, k, n, s} t_{\alpha n} c_{\alpha ks}^\dagger d_{ks} + h.c.$$

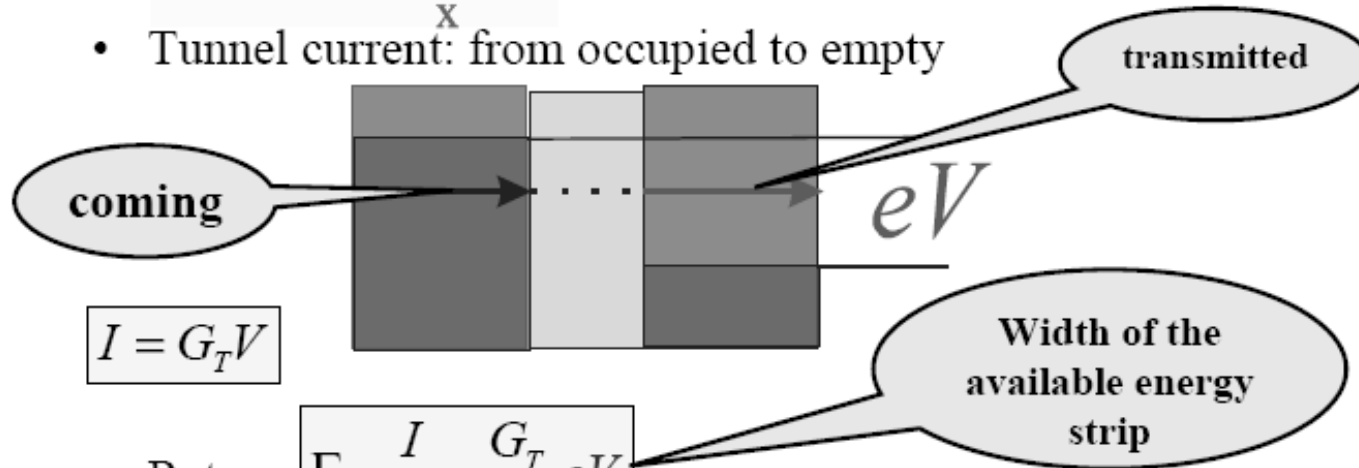
*lead*      *lead energy*      *dot level*      *spin*

# Tunneling in metals (No CB)

- Tunnel barrier



- Tunnel current: from occupied to empty

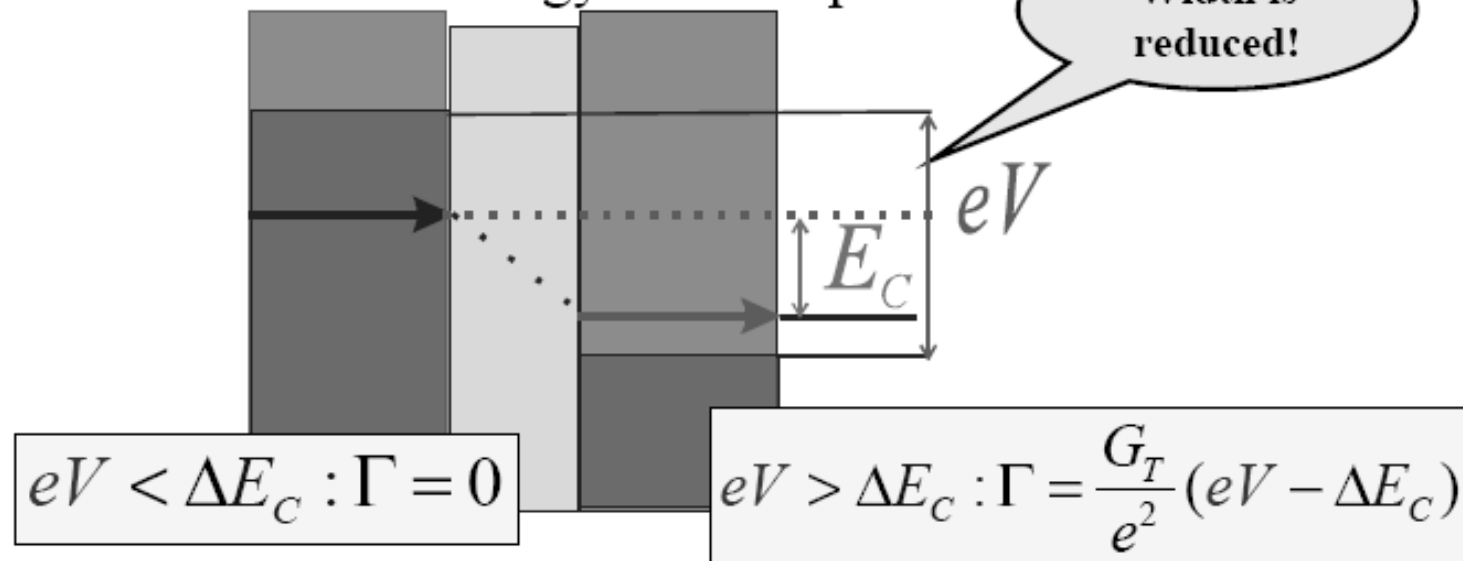


$$I = G_T V$$

- Rate: 
$$\Gamma = \frac{I}{e} = \frac{G_T}{e^2} eV$$

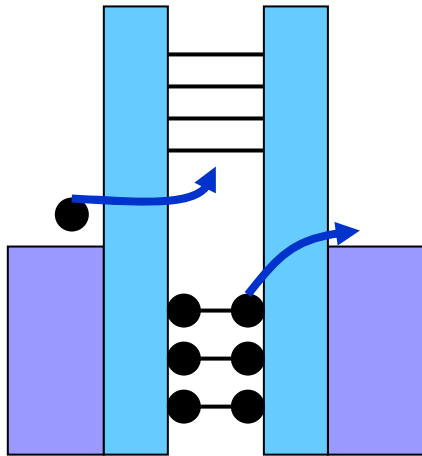
# Tunneling and Coulomb blockade

- Now the same with Coulomb blockade
- Electrostatic energy must be paid



- blockade

# Inelastic cotunneling



*inelastic cotunneling*

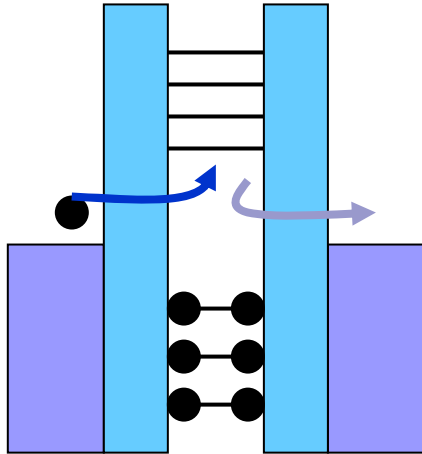
*(excitations left behind)*

state  $k$  on Left  $\rightarrow$

state  $k'$  on Right + Dot ( $n$  filled;  $m$  empty)



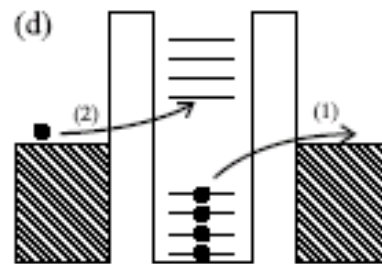
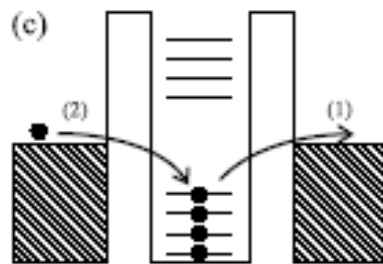
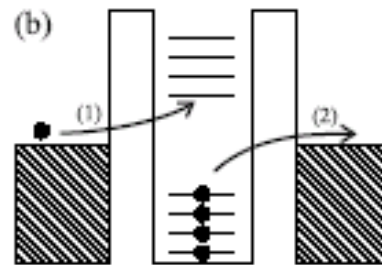
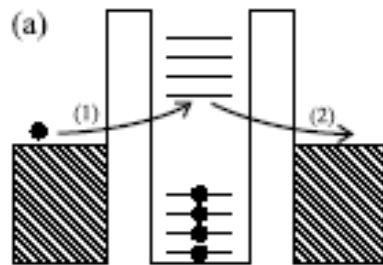
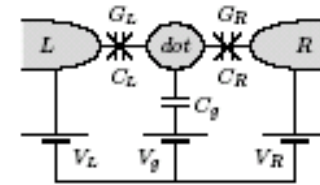
# Elastic cotunneling



*elastic cotunneling*

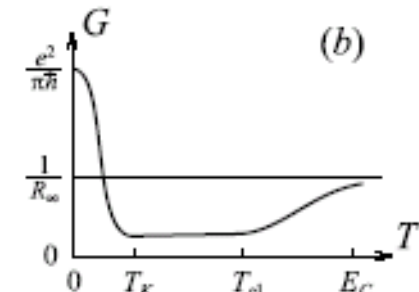
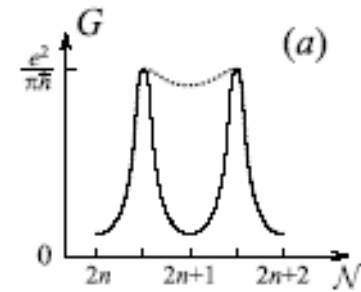
*(no excitations left behind)*

# Tunneling and co-tunneling (summary)



Elastic

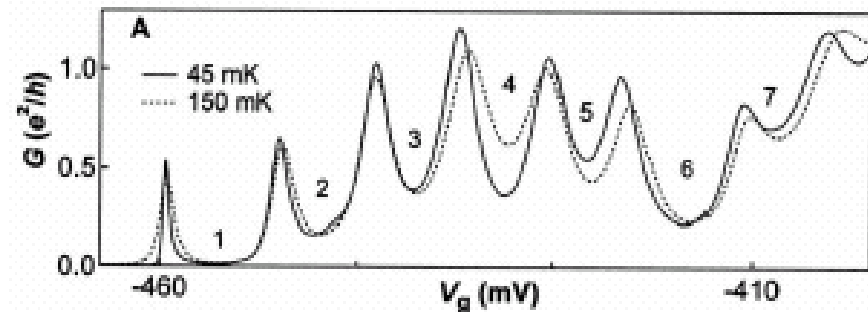
Inelastic



Electron-like

co-tunneling

Hole-like

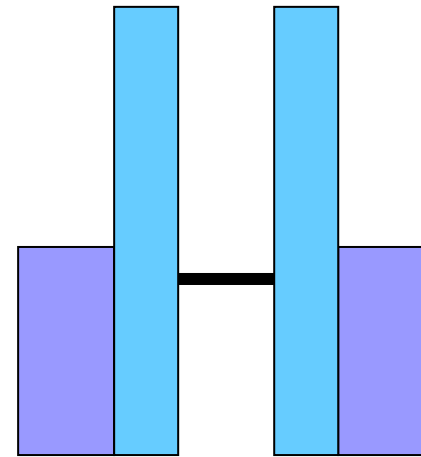


# DISCRETE SPECTRUM QD: coherent vs. Incoherent transport

$$H = H_L + H_R + H_{dot} + H_{tun}$$

$$H_{dot} = \varepsilon \sum_{\sigma} n_{\sigma} + E_c n_{\uparrow} n_{\downarrow}$$

$$H_{tun} = t_{\alpha} c_{\alpha k \sigma}^{\dagger} d_{\sigma} + h.c.$$



a single orbital level  
spin  $\uparrow$  or  $\downarrow$

This is a story for next lecture!