

## Quantum Hall effect in Bernal-stacked Bilayer graphene

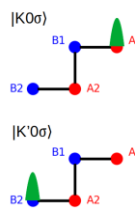
- Landau level energy:

$$E_N = \text{sign}(N)\hbar\omega_c\sqrt{|N|(|N|-1)}$$

- Landau Levels  $N = 0$  and  $N = 1$  are nearly degenerate ( $E_0 = E_1 = 0$ )
- 8-fold nearly degenerate Zero-energy Landau level
- 3 degrees of freedom: spin  $\sigma = \uparrow/\downarrow$ , valley  $\xi = K/K'$ , orbital (Landau Level)  $N=0, 1$
- By neglecting all mixing with higher Landau levels, the situation is reminiscent of extreme Landau level mixing

→ Favor the formation of Wigner crystal [1]

### Single-particle wavefunction

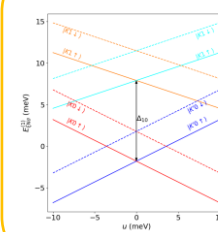


In the basis (A1, B2, A2, B1):

$$\begin{aligned}\Psi_0^K &= (\phi_0, 0, 0, 0) \\ \Psi_1^K &= (\sqrt{1-\gamma}\phi_1, 0, \sqrt{\gamma}\phi_0, 0) \\ \Psi_0^{K'} &= (0, \phi_0, 0, 0) \\ \Psi_1^{K'} &= (0, \sqrt{1-\gamma}\phi_1, 0, \sqrt{\gamma}\phi_0)\end{aligned}$$

### Parameter $\gamma$

### Single-particle energy levels



- Single electron energy:

$$E_{\xi N \sigma}^{(1)} = -\xi \alpha_N \frac{u}{2} + N \Delta_{10}^{(0)} - E_{Z \sigma}$$

- $u$  = energy bias between the two layers (interlayer bias)
- The levels  $|K0 \uparrow\rangle$  and  $|K1 \uparrow\rangle$  cross at large value of  $u$
- In the vicinity of the crossing, we define the orbital splitting energy:
 
$$\Delta_{10} = E_{K0 \uparrow}^{(1)} - E_{K1 \uparrow}^{(1)}$$

$$= -\Delta_{10}^{(0)} + (\alpha_1 - \alpha_0) \frac{u}{2}$$

### Parameter $\Delta_{10}$

### Hamiltonian

We consider the regime in the vicinity of the crossing between the levels  $|K0 \uparrow\rangle$  and  $|K1 \uparrow\rangle$ :

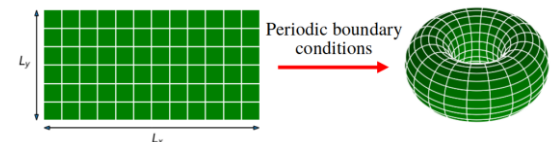
$$\mathcal{H} = \frac{1}{2} \sum_{N_1, N_2, N_3, N_4=0}^1 \sum_{j_1, j_2, j_3, j_4=1}^{N_\phi} \mathcal{A}_{j_1 j_2 j_3 j_4}^{N_1 N_2 N_3 N_4} c_{N_1 j_1}^\dagger c_{N_2 j_2}^\dagger c_{N_3 j_3} c_{N_4 j_4} + \Delta_{10} \tilde{N}_0$$

$$\mathcal{A}_{j_1 j_2 j_3 j_4}^{N_1 N_2 N_3 N_4} = \frac{1}{L_x L_y} \sum_{(s, \ell) \neq (0,0)} \delta_{q_x, s \frac{2\pi}{L_x}} \delta_{q_y, \ell \frac{2\pi}{L_y}} \frac{2\pi e^2}{\epsilon q} e^{-\frac{q^2 l^2}{2}} e^{i2\pi(j_1 - j_3) \frac{s}{N_\phi}} F_{N_1 N_4}(\mathbf{q}) F_{N_2 N_3}(-\mathbf{q}) \delta'_{\ell, j_1 - j_4} \delta'_{j_1 + j_2, j_3 + j_4}$$

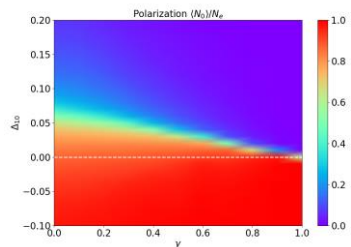
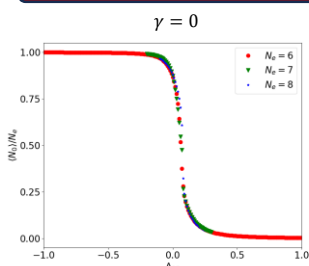
### Form factors

$$\begin{aligned}F_{00} &= 1 \\ F_{11} &= \gamma + (1-\gamma) \left(1 - \frac{q^2 l^2}{2}\right) \\ F_{01} &= \sqrt{1-\gamma} \frac{(iq_x + q_y)l}{\sqrt{2}} \\ F_{10} &= \sqrt{1-\gamma} \frac{(iq_x - q_y)l}{\sqrt{2}}\end{aligned}$$

### Method: exact diagonalization on the torus geometry

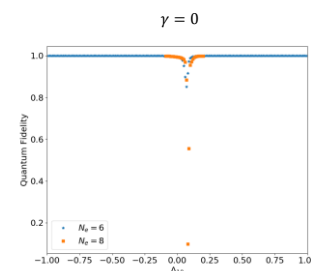


### Orbital transition (filling factor 1/3)



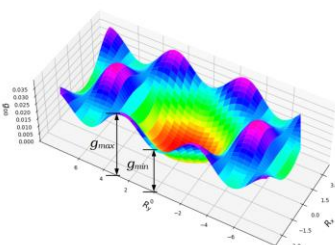
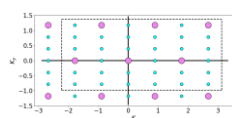
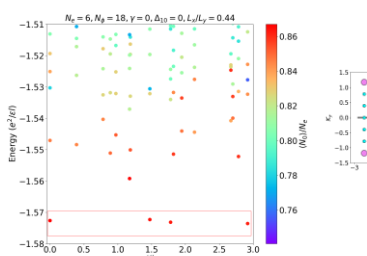
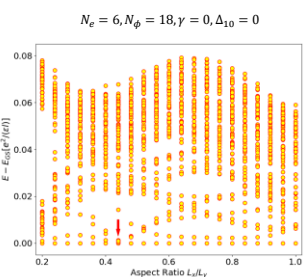
- Polarization in the orbital  $N = 0$  of the ground state:
 
$$\frac{\langle N_0 \rangle}{N_e} = \frac{\langle \Psi_{GS} | \tilde{N}_0 | \Psi_{GS} \rangle}{N_e}$$
- An abrupt change in  $\langle N_0 \rangle / N_e$  when  $\Delta_{10}$  increases while keeping  $\gamma$  unchanged
- Electrons transfer from orbital  $N = 0$  to orbital  $N = 1$

### Quantum fidelity



- Quantum fidelity:
 
$$\mathcal{F} = |\langle \Psi_{GS}(\Delta_{10} + d\Delta_{10}) | \Psi_{GS}(\Delta_{10}) \rangle|^2$$
- A dip in  $\mathcal{F}$  at the same value of  $\Delta_{10}$  of the polarization transition

### The Wigner crystal (filling factor 1/3)



$$N_e = 6, N_\phi = 18, \gamma = 0, \Delta_{10} = 0, \frac{L_x}{L_y} = 0.44$$

- Pair correlation function between electrons in orbitals  $\alpha$  and  $\beta$  separated by a distance  $\mathbf{R}$ :

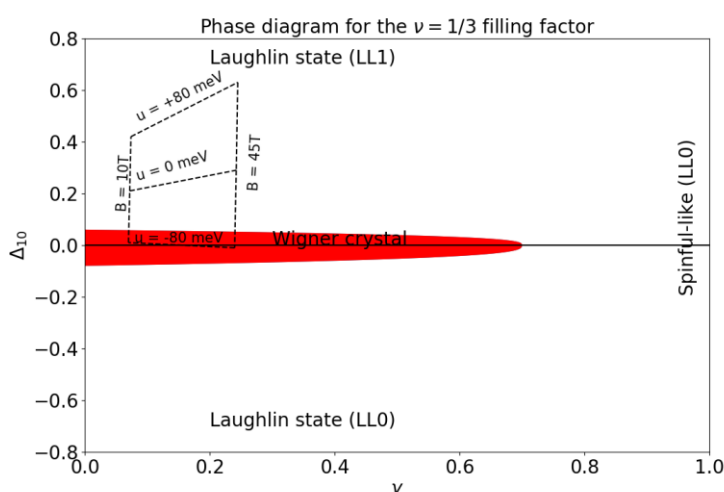
$$g_{\alpha\beta}(\mathbf{R}) = \frac{L_x L_y}{N_e^2} \sum_{i,j} \delta^2(\mathbf{r}_i - \mathbf{r}_j - \mathbf{R}) \langle |\alpha\rangle \langle \alpha|_i \langle \beta| \langle \beta|_j \rangle$$

- The overdensities in  $g_{00}$  show the real-space distribution of Wigner crystal
- Situation reminiscent to extreme Landau level mixing
  - favoring Wigner crystal formation
- In the Wigner crystal phase, electrons occupy the orbital  $N = 0$

Ground state degeneracy → Wigner crystal

Many-body Brillouin zone

### The phase diagram



### Conclusion

- Competition between Wigner crystal and quantum Hall liquid in the vicinity of the crossing between the levels  $|K0 \uparrow\rangle$  and  $|K1 \uparrow\rangle$  of filling factor  $\nu = 1/3$
- The phase transition can be controlled by tuning the electric energy bias  $u$  between the two layers (interlayer bias)
- Electrically-induced Wigner crystal

### References

- Phys. Rev. Lett. **121**, 116802 (2018)
- Nature Communications, 8:948 (2017)
- Phys. Rev. Lett. **124**, 097604 (2020)
- Phys. Rev. B **107**, 125129 (2023)