



# Neutrino Physics

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VSOP

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Part 1

NASA Hubble Photo



# What Are Neutrinos Good For?

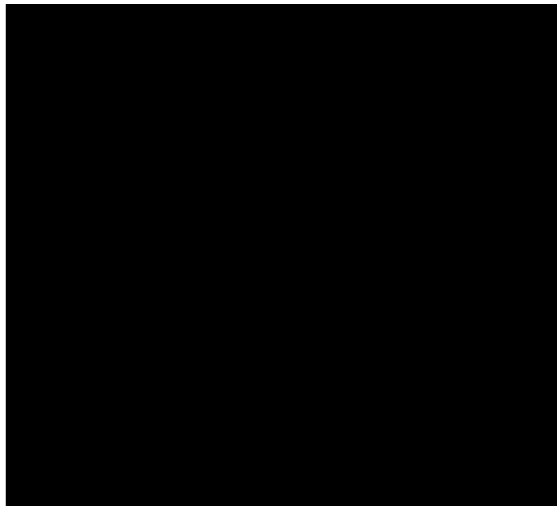
Energy generation in the sun starts with the reaction —

$$p + p \rightarrow d + e^+ + \nu$$

Spin:  $\frac{1}{2} \quad \frac{1}{2} \quad 1 \quad \frac{1}{2} \quad \frac{1}{2}$

Without the neutrino, angular momentum  
would not be conserved.

Uh, oh .....



We would not  
be here!



# The Neutrinos



**Neutrinos and photons are by far the most abundant elementary particles in the universe.  
There are 340 neutrinos/cc.**

The neutrinos are spin –  $1/2$ , electrically neutral, leptons.

The only known forces they experience are  
the weak force and gravity.

This means that their interactions with other matter  
have very low strength.

Thus, neutrinos are difficult to detect and study.

Their weak interactions are successfully described  
by the Standard Model.



# The Neutrino Revolution

(1998 – ...)

Neutrinos have nonzero masses!

Leptons mix!



# The Origin of Neutrino Mass

The fundamental constituents of matter are the *quarks*, the *charged leptons*, and the *neutrinos*.

*Most theorists strongly suspect that the origin of the *neutrino* masses is different from the origin of the *quark* and *charged lepton* masses.*

The Standard-Model *Higgs* may still be involved, but not in the same way as for the quarks and charged leptons.

*More later .....*



The discovery of neutrino mass  
and leptonic mixing  
comes from the observation of  
*neutrino flavor change*  
*(neutrino oscillation)*.



# Introduction to Neutrino Oscillation

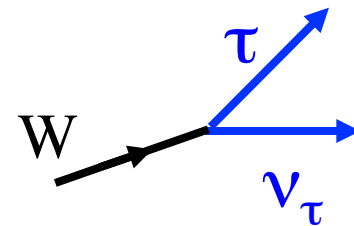
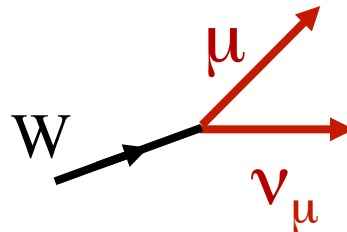
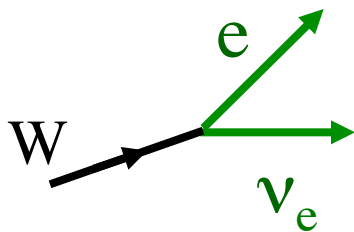


# The Neutrino Flavors

There are three flavors of charged leptons:  $e$ ,  $\mu$ ,  $\tau$

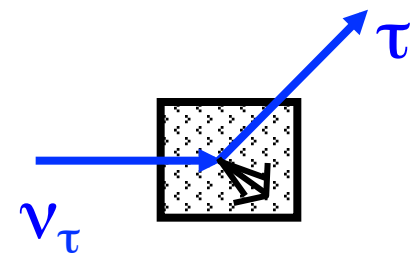
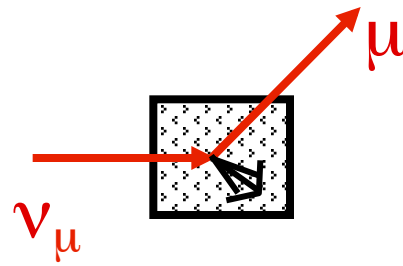
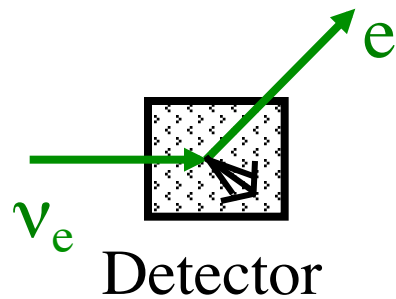
There are three known flavors of neutrinos:  $\nu_e$ ,  $\nu_\mu$ ,  $\nu_\tau$

We *define* the neutrinos of specific flavor,  $\nu_e$ ,  $\nu_\mu$ ,  $\nu_\tau$ , by W boson decays:

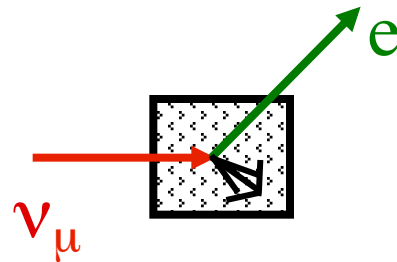




As far as we know, when a neutrino of given flavor interacts and turns into a charged lepton, that charged lepton will always be of the same flavor as the neutrino.



but not

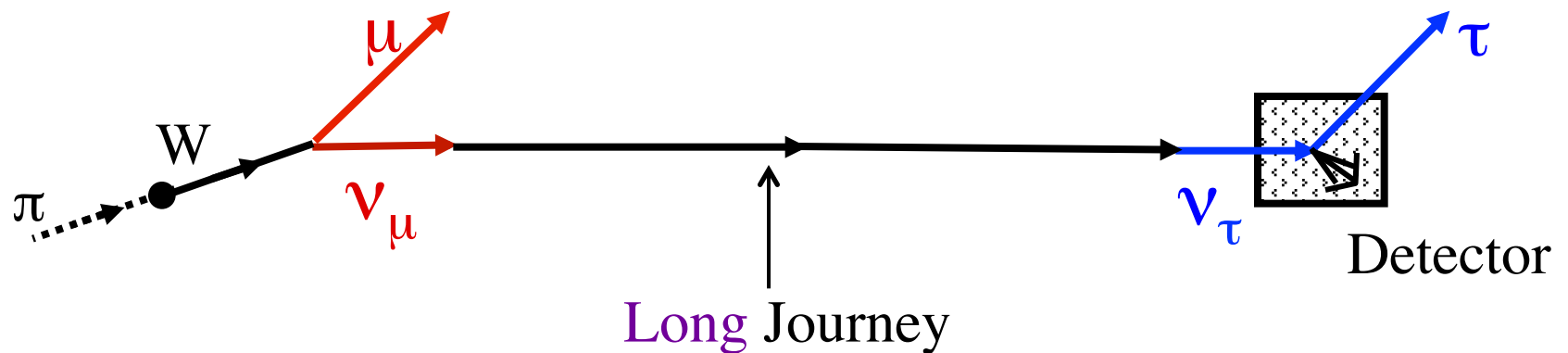


Lederman  
Schartz  
Steinberger



# Neutrino Flavor Change (“Oscillation”)

*If neutrinos have masses, and leptons mix, we can have —*



Give a  $\nu$  time to change character, and you can have

for example:  $\nu_\mu \longrightarrow \nu_\tau$

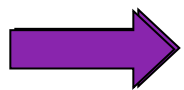
The last 16 years have brought us compelling evidence that such flavor changes actually occur.



Flavor *change* does not add neutrinos to a beam — it just changes the flavor of those already present.

If some of the neutrinos in a beam born as  $\nu_\mu$  are turning into  $\nu_\tau$ , there must be fewer  $\nu_\mu$  left in the beam.

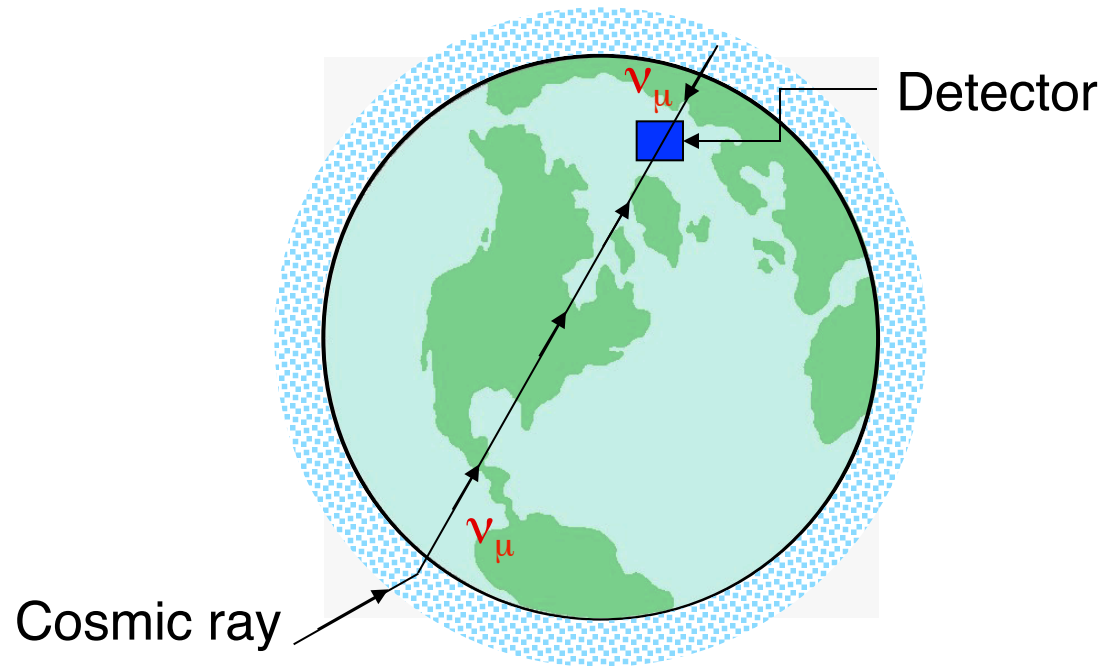
*Disappearance* of some of the old flavors



*Appearance* of some of the new flavors



# — Atmospheric Neutrinos — The First Compelling Evidence

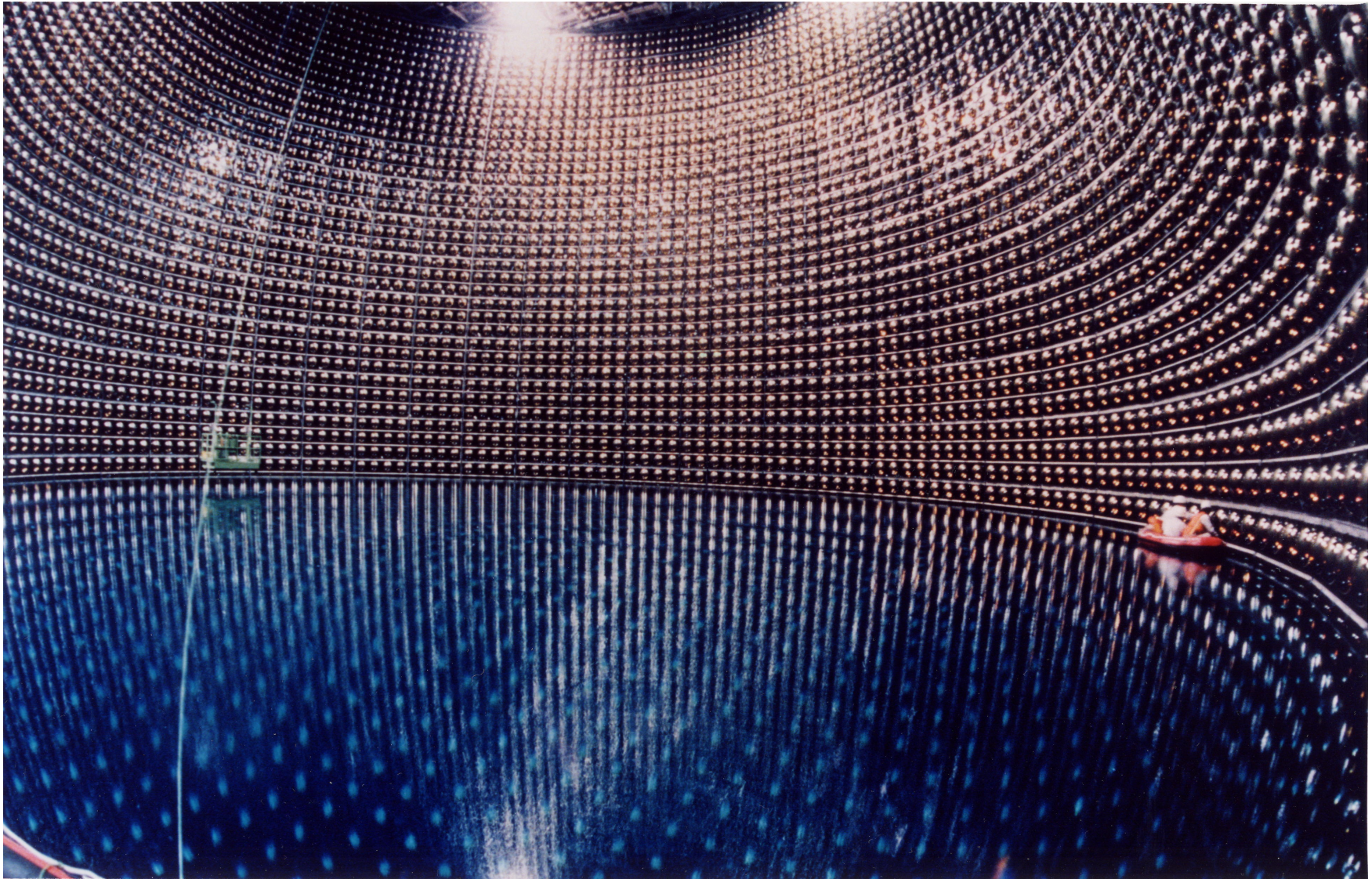


Isotropy of the  $\gtrsim 2$  GeV cosmic rays + Gauss' Law + No  $\nu_\mu$  disappearance

$$\Rightarrow \frac{\phi_{\nu_\mu}(\text{Up})}{\phi_{\nu_\mu}(\text{Down})} = 1 .$$

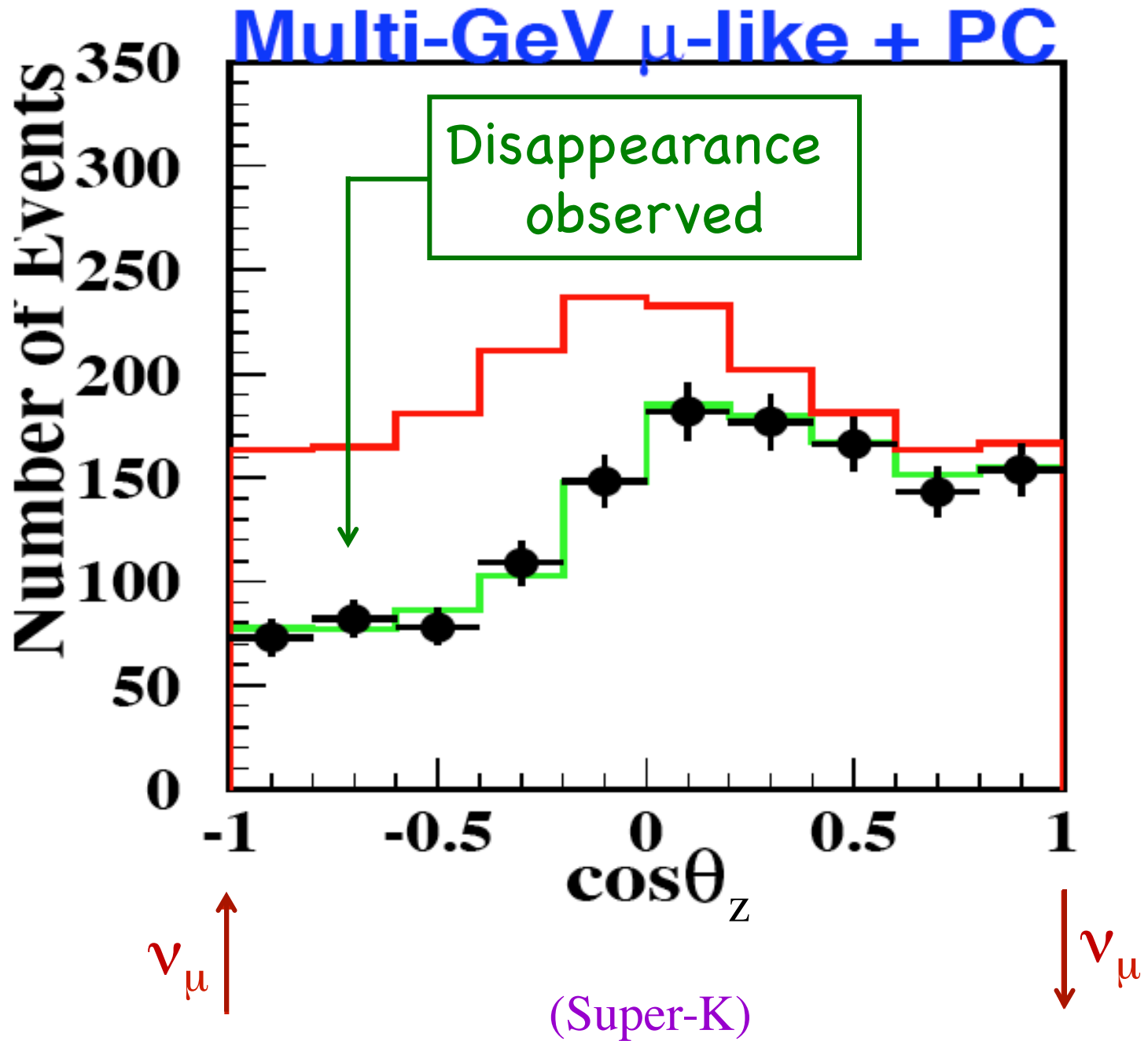
But Super-Kamiokande finds for  $E_\nu > 1.3$  GeV,  $\frac{\phi_{\nu_\mu}(\text{Up})}{\phi_{\nu_\mu}(\text{Down})} \cong 1/2 .$





***Super-Kamiokande:*** 50 ktons of water, surrounded by 11k phototubes that detect Cerenkov light from a  $\mu$  or  $e$





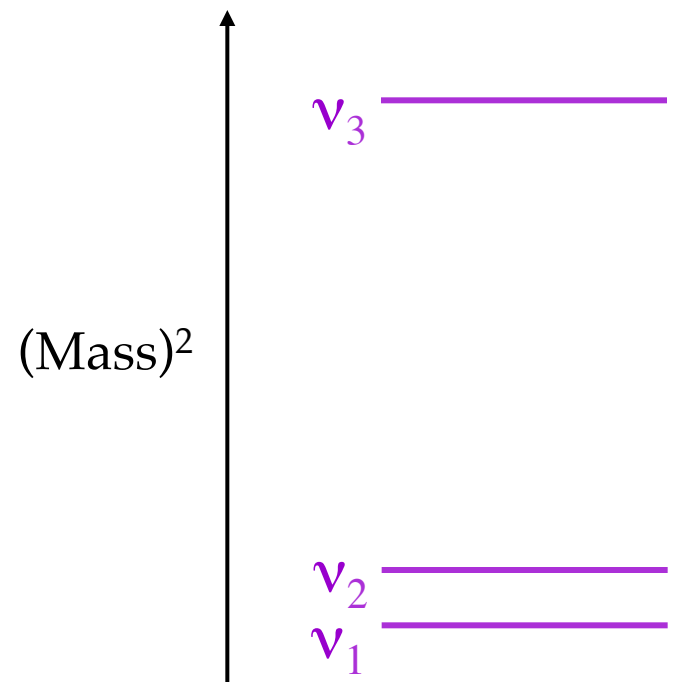


# The Physics of Neutrino Oscillation



# Flavor Change Requires *Neutrino Masses*

There must be some spectrum  
of neutrino mass eigenstates  $\nu_i$ :



$$\text{Mass}(\nu_i) \equiv m_i$$



# Flavor Change Requires *Leptonic Mixing*

The neutrinos  $\nu_{e,\mu,\tau}$  of definite flavor

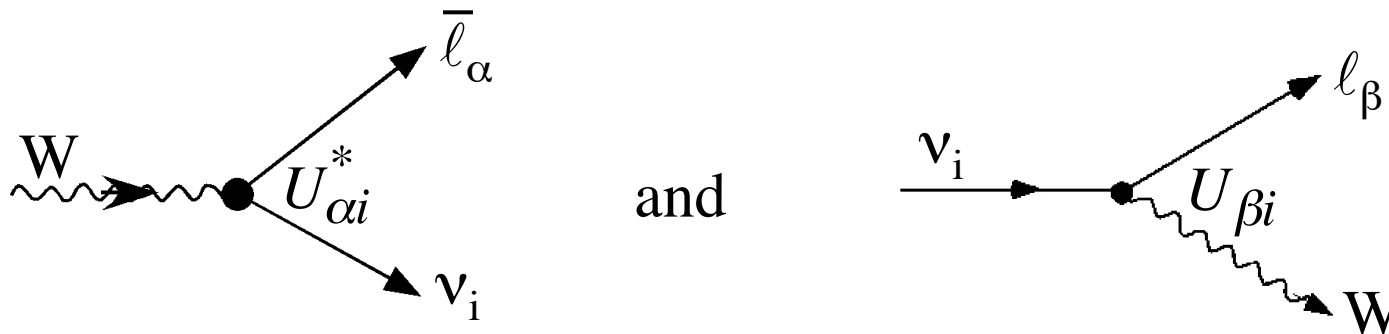
$$(W \rightarrow e\nu_e \text{ or } \mu\nu_\mu \text{ or } \tau\nu_\tau)$$

are **superpositions** of the neutrinos of definite mass:

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle .$$

Neutrino of flavor  
 $\alpha = e, \mu, \text{ or } \tau$

Neutrino of definite mass  $m_i$   
Unitary Leptonic Mixing Matrix

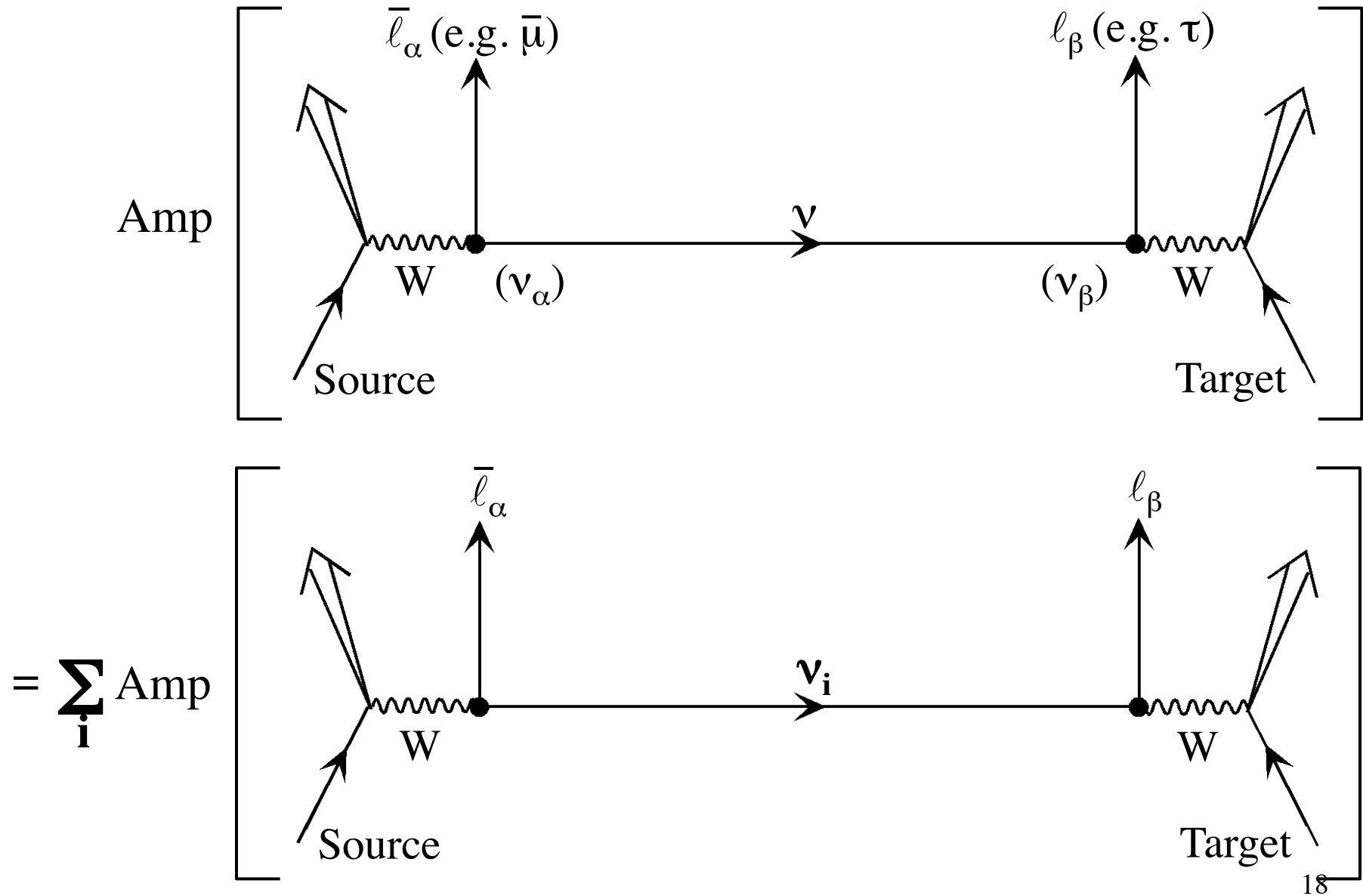


$\ell_\alpha$  is a charged lepton ( $\ell_e \equiv e, \ell_\mu \equiv \mu, \ell_\tau \equiv \tau$ ).



# Neutrino Flavor Change (“Oscillation”)

(Approach of BK and L. Stodolsky)





$$= \sum_{\mathbf{i}} \text{Amp} \left[ \begin{array}{c} \text{Amp}(\nu_{\alpha} \rightarrow \nu_{\beta}) \\ \begin{array}{c} \bar{\ell}_{\alpha} \\ \uparrow \\ U_{\alpha i}^* \\ \bullet \\ \text{Source} \end{array} \xrightarrow[\nu_i]{\begin{array}{c} \text{---} \frac{i}{\hbar c} \left( m_i^2 c^4 \right)^2 \frac{L}{2E} \text{---} \\ \text{---} \frac{i}{\hbar c} \left( m_i^2 c^4 \right)^2 \frac{L}{2E} \text{---} \\ \text{---} \frac{i}{\hbar c} \left( m_i^2 c^4 \right)^2 \frac{L}{2E} \text{---} \end{array}} \begin{array}{c} \ell_{\beta} \\ \uparrow \\ U_{\beta i} \\ \bullet \\ \text{Target} \end{array} \end{array} \right]$$

$$= \sum_i U_{\alpha i}^* e^{-\frac{i}{\hbar c} \left( m_i^2 c^4 \right)^2 \frac{L}{2E}} U_{\beta i}$$

Rest energy
Energy
Proper time
Distance

$$\text{Probability} \uparrow P(\nu_{\alpha} \rightarrow \nu_{\beta}) = \left| \text{Amp}(\nu_{\alpha} \rightarrow \nu_{\beta}) \right|^2$$



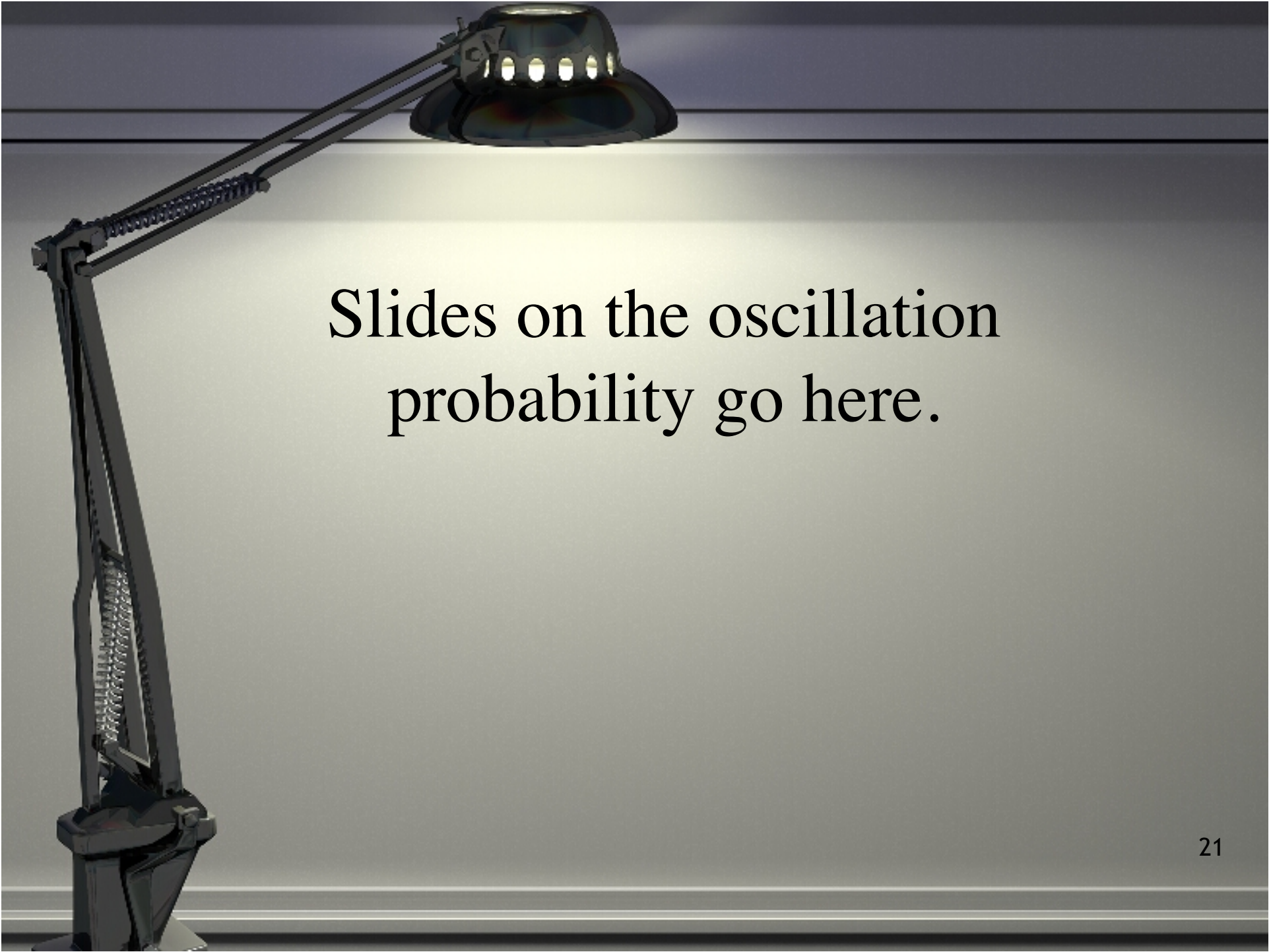
Why does  $e^{-\frac{i}{\hbar}(m_i c^2)\tau_i}$  describe neutrino propagation?

If, in the lab. frame, a neutrino  $\nu$  of mass  $m$ ,  
with momentum  $p$  and energy  $E$ ,  
travels a distance  $L$  in time  $t$ ,  
its wave function picks up a factor —

$$\exp\left[\frac{i}{\hbar}(pL - Et)\right] = \exp\left[-\frac{i}{\hbar}(mc^2)\tau\right]$$

By the Lorentz  
transformation



A 3D rendered desk lamp with a black adjustable arm and a silver-colored lamp head. The lamp is positioned on the left side of the frame, casting a warm, yellowish glow onto a light gray presentation slide. The slide is mounted on a dark gray wall with horizontal lines. The text on the slide is in a black serif font.

Slides on the oscillation  
probability go here.



# Neutrino Flavor Change In Matter



Coherent forward scattering via this  
W-exchange interaction leads to  
an extra interaction potential energy —

$$V_W = \begin{cases} +\sqrt{2}G_F N_e, & \nu_e \\ -\sqrt{2}G_F N_e, & \bar{\nu}_e \end{cases}$$

Fermi constant

Electron density

This raises the effective mass of  $\nu_e$ , and lowers that of  $\bar{\nu}_e$ .



The fractional importance of matter effects on an oscillation involving a vacuum splitting  $\Delta m^2$  is —

$$\begin{array}{cc} \text{Interaction} & \text{Vacuum} \\ \text{energy} & \text{energy} \\ \underbrace{\hspace{1.5cm}} & \underbrace{\hspace{1.5cm}} \\ [\sqrt{2}G_F N_e] & / \ [\Delta m^2/2E] \equiv x . \end{array}$$

The matter effect —

- Grows with neutrino energy  $E$
- Is sensitive to  $\text{Sign}(\Delta m^2)$
- Reverses when  $\nu$  is replaced by  $\bar{\nu}$

This last is a “fake CP violation”, but the matter effect is negligible when  $x \ll 1$ .



# *Evidence For Flavor Change*

## Neutrinos

## Evidence of Flavor Change

Solar  
Reactor  
(Long-Baseline)

Compelling  
Compelling


Atmospheric  
Accelerator  
(Long-Baseline)

Compelling  
Compelling

Accelerator & Reactor  
(Short-Baseline)

“Interesting”





# Further Highlights of the History



# Solar Neutrinos

History –

Nuclear reactions in the core of the sun  
produce  $\nu_e$ . Only  $\nu_e$ .





Theorists, especially **John Bahcall**, calculated the produced  $\nu_e$  flux vs. energy  $E$ .





Ray Davis' Homestake experiment measured the higher-E part of the  $\nu_e$  flux  $\phi_{\nu_e}$  that arrives at earth.



The Homestake experiment could detect only  $\nu_e$ . It found:

$$\frac{\phi_{\nu_e}(\text{Homestake})}{\phi_{\nu_e}(\text{Theory})} = 0.34 \pm 0.06$$

## The Possibilities:

The theory was wrong.

The experiment was wrong.

Both were wrong.

Neither was wrong. Two thirds of the  $\nu_e$  flux changes into a flavor or flavors that the Homestake experiment could not see.



## The Resolution —

Sudbury Neutrino Observatory (SNO) measures, for the high-energy part of the solar neutrino flux:

$$\nu_{\text{sol}} d \rightarrow e p p \Rightarrow \phi_{\nu_e}$$

$$\nu_{\text{sol}} d \rightarrow \nu n p \Rightarrow \phi_{\nu_e} + \phi_{\nu_\mu} + \phi_{\nu_\tau} \quad (\nu \text{ remains a } \nu)$$

---

From the two reactions,

$$\frac{\phi_{\nu_e}}{\phi_{\nu_e} + \phi_{\nu_\mu} + \phi_{\nu_\tau}} = 0.301 \pm 0.033$$

Clearly,  $\phi_{\nu_\mu} + \phi_{\nu_\tau} \neq 0$ . Neutrinos change flavor.

For solar neutrinos,  $P(\nu_e \rightarrow \nu_e) = 0.3$ .



Change of flavor does not change the total number of neutrinos.

The total flux,  $\phi_{\nu_e} + \phi_{\nu_\mu} + \phi_{\nu_\tau}$ , should agree with Bahcall's prediction.

SNO:  $\phi_{\nu_e} + \phi_{\nu_\mu} + \phi_{\nu_\tau} = (5.54 \pm 0.32 \pm 0.35) \times 10^6/\text{cm}^2\text{sec}$

Theory\*:  $\phi_{\text{total}} = (5.69 \pm 0.91) \times 10^6/\text{cm}^2\text{sec}$

\*Bahcall, Basu, Serenelli

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John Bahcall and Ray Davis both stuck to  
their results for several decades,  
and both were *right* all along.



A 3D-rendered desk lamp with a black adjustable arm and a silver-colored lamp head is positioned on the left side of the frame. The lamp is turned on, casting a warm, yellowish glow onto a light gray surface that serves as the background for the text. The lamp's arm is extended upwards and to the right, with a coiled spring visible. The lamp head has a circular base with several small, glowing light sources.

# KamLAND Evidence for Oscillatory Behavior



The **KamLAND** detector studied  $\bar{\nu}_e$  produced by Japanese nuclear power reactors  $\sim 180$  km away.

For **KamLAND**,  $x_{\text{Matter}} < 10^{-2}$ . Matter effects were negligible.

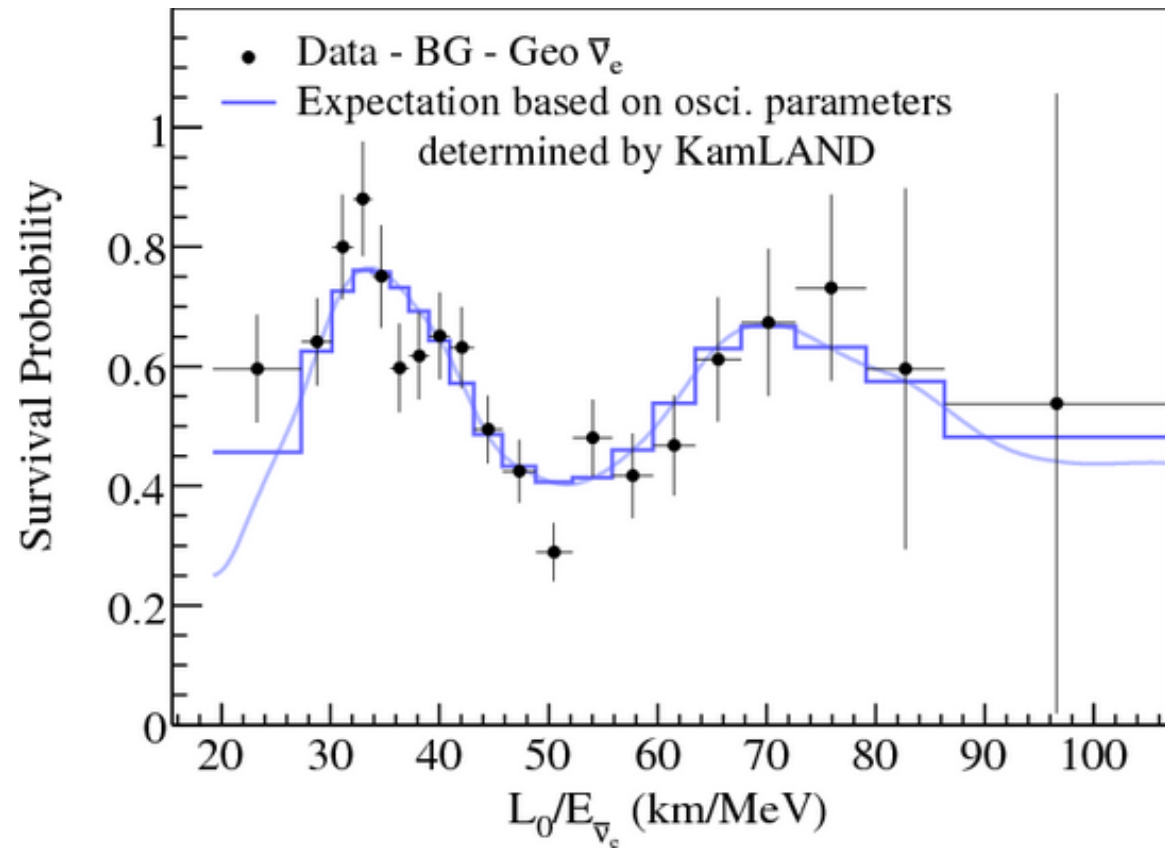
The  $\bar{\nu}_e$  survival probability,  $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$ , should **oscillate** as a function of  $L/E$  following the vacuum oscillation formula.

In the two-neutrino approximation, we expect —

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2 2\theta \sin^2 \left[ 1.27 \Delta m^2 (eV^2) \frac{L(km)}{E(GeV)} \right].$$



Survival  
probability  
 $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$   
of reactor  $\bar{\nu}_e$



$L_0 = 180$  km is a flux-weighted average travel distance.

*$P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$  actually oscillates!*



# The End

## — Part 1